

Vortex Solution in Holographic Two-Band Superfluid/Superconductor

Hai-Qing Zhang
Utrecht University

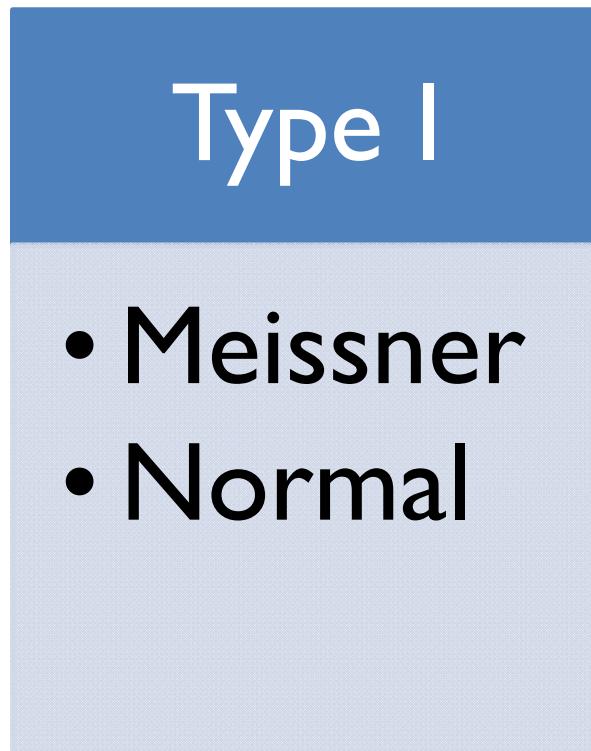
Instituto de Física Teórica
Universidad Autónoma de Madrid
10/11/2015

arXiv:1511.01325, M.-S. Wu, Shang-Yu Wu and HQZ

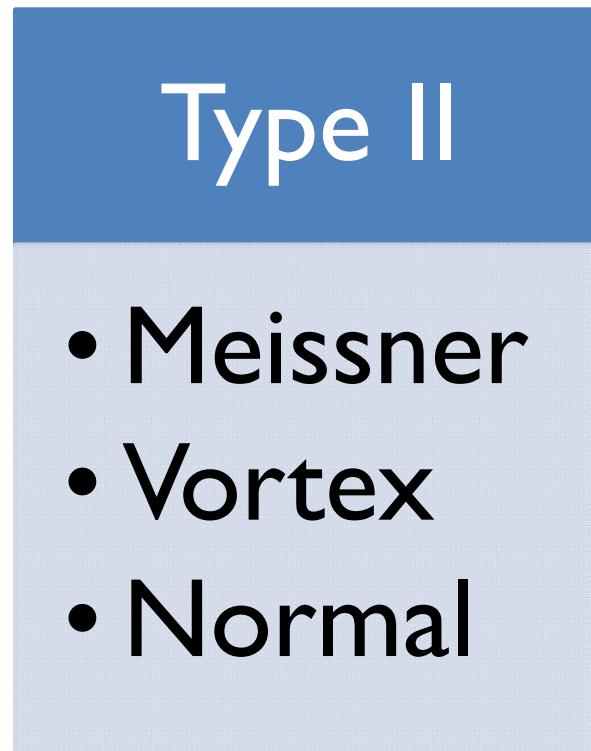


SUPERCONDUCTING VORTEX IN CONDENSED MATTER PHYSICS

* Ginzburg-Landau (GL) parameter κ



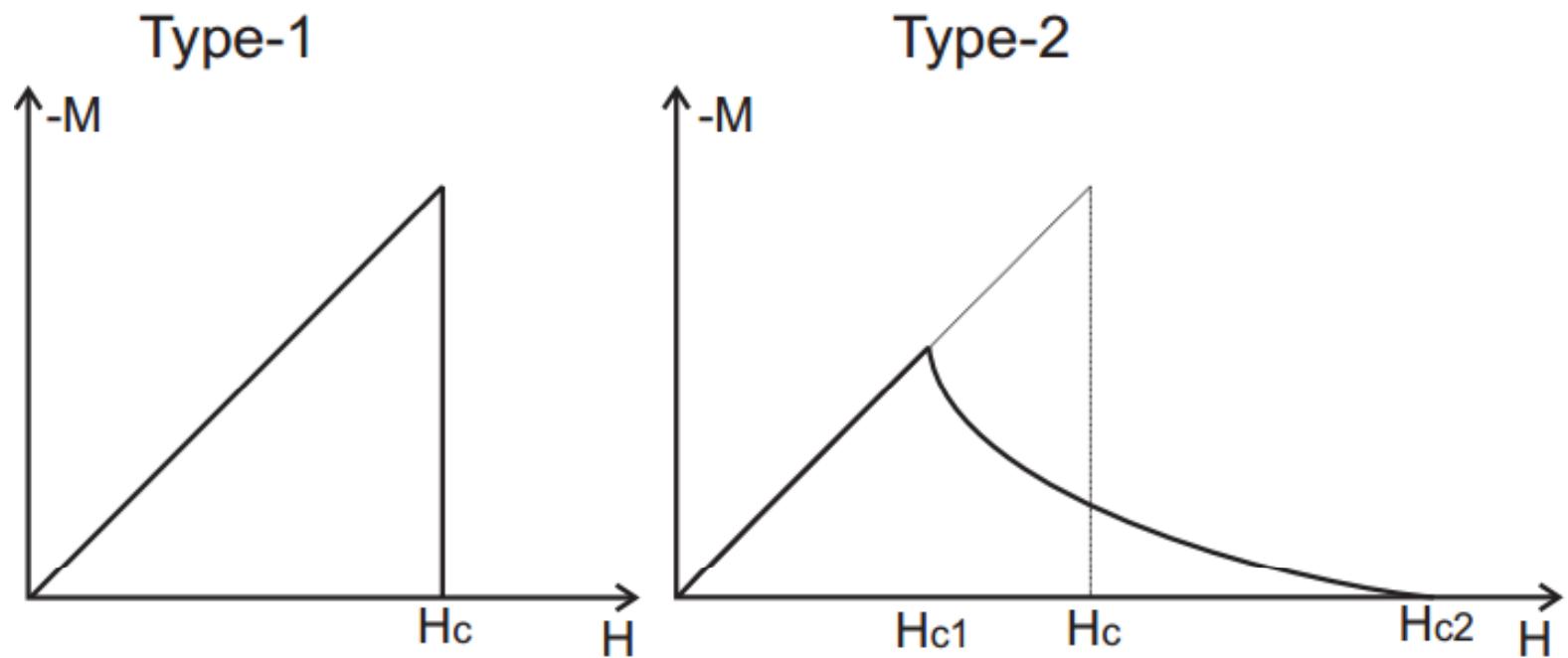
$$\kappa < 1/\sqrt{2}$$



$$\kappa > 1/\sqrt{2}$$



$\kappa = \lambda/\xi$, λ -penetration length, ξ -coherence length



M- Magnetization

H- External Magnetic Field

* Supercurrent velocity from GL theory

$$m^* v_s = \hbar \nabla \varphi - \frac{e^* A}{c}$$

* Coherence length ξ and penetration length λ near T_c

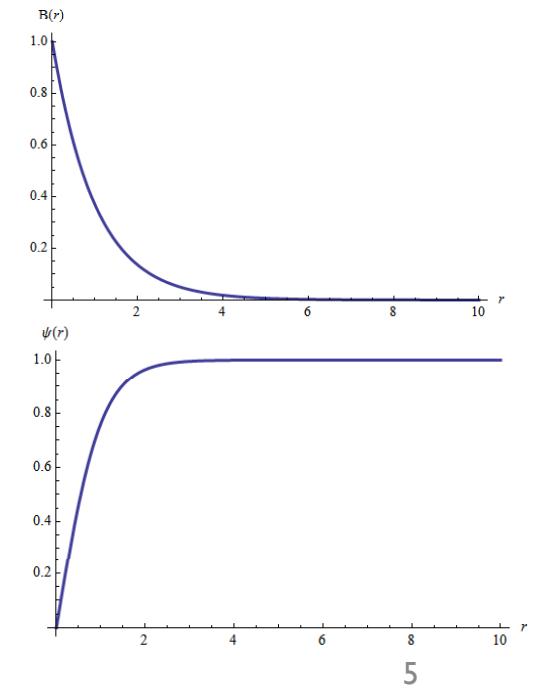
$$\xi(T) \propto \frac{1}{\sqrt{1-T/T_c}} \quad \lambda(T) \propto \frac{1}{\sqrt{1-T/T_c}}$$

* Magnetic field and penetration length λ

$$B(r) \approx \sqrt{\frac{\lambda}{r}} e^{-r/\lambda}$$

* Condensate and coherence length

$$\psi(r) = \psi_0 \tanh\left(\frac{r}{\sqrt{2}\xi}\right)$$



PRB, 69, 054508 (2004)

* Two component Ginzburg-Landau (TCGL)

$$F_{GL} = \int dx \left[\alpha_1 |\Delta_1|^2 + \alpha_2 |\Delta_2|^2 - \gamma (\Delta_1^* \Delta_2 + \Delta_2^* \Delta_1) \right. \\ \left. + \frac{1}{2} \beta_1 |\Delta_1|^4 + \frac{1}{2} \beta_2 |\Delta_2|^4 + K_{1i} |\nabla_i \Delta_1|^2 + K_{2i} |\nabla_i \Delta_2|^2 \right],$$

$$\nabla_i = \partial_i + i \frac{2\pi}{\Phi_0} A_i$$

γ — Josephson
interband coupling

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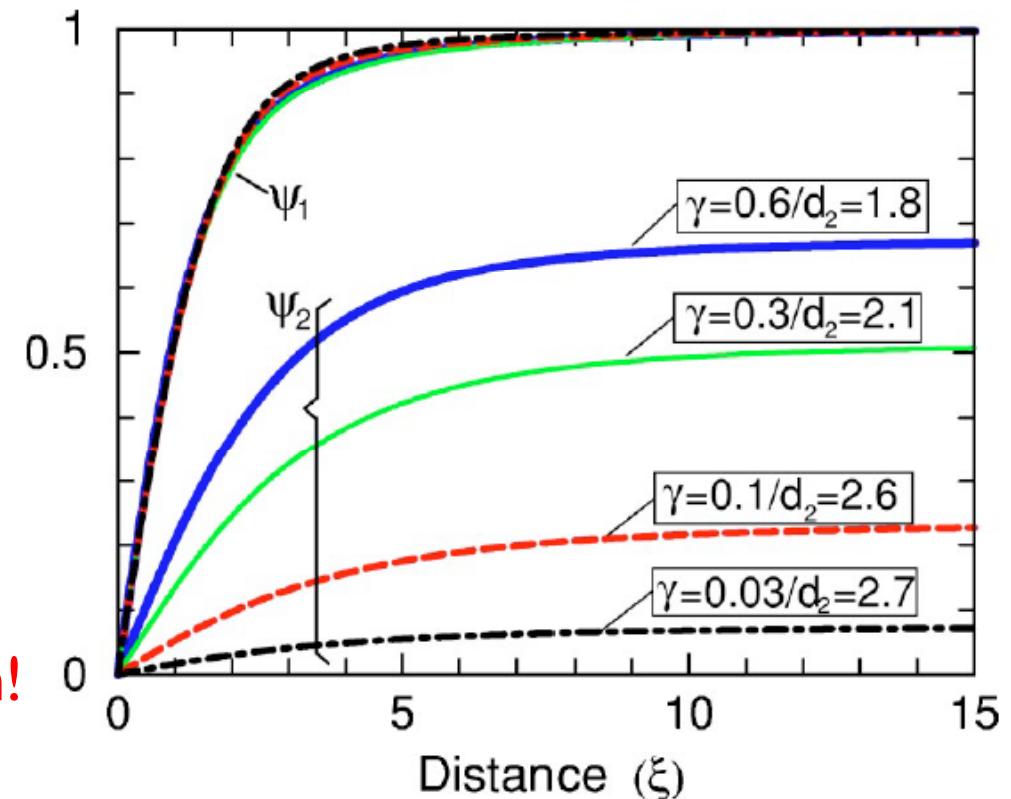
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$$\nabla_i = \partial_i + i \frac{2\pi}{\Phi_0} A_i$$

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Note: one single vortex solution!



HOLOGRAPHIC STUDY

* Action of the two-band superconductor

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} [R + \frac{6}{L^2} - \frac{1}{4}F^2 - |\partial\psi_1 - iqA\psi_1|^2 - |\partial\psi_2 - iqA\psi_2|^2 - V(\psi_1, \psi_2)],$$

$$V(\psi_1, \psi_2) = m_1^2|\psi_1|^2 + m_2^2|\psi_2|^2 + \epsilon(\psi_1\psi_2^* + \psi_1^*\psi_2) + \eta|\psi_1|^2|\psi_2|^2$$

ϵ is the intercomponent Josephson coupling

η is a density-density coupling

* 4d AdS-Schwarzschild black hole

$$ds^2 = \frac{1}{z^2} \left(-f(z)dt^2 + \frac{dz^2}{f(z)} + d\rho^2 + \rho^2 d\phi^2 \right), \quad f(r) = 1 - \left(\frac{z}{z_h} \right)^3$$

$z \rightarrow 0$ is the infinite boundary

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Disregard this term later

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* Ansatz for vortex

$$\psi_1 = \varphi_1 e^{i\theta_1}, \psi_2 = \varphi_2 e^{i\theta_2}$$

$$\varphi_1(\rho, z), \varphi_2(\rho, z), A_t(\rho, z), A_\phi(\rho, z), \theta_1 = n_1\phi, \theta_2 = n_2\phi$$

* Equation of Motions

$$0 = f\partial_z^2 A_t + \frac{\partial_\rho A_t}{\rho} + \partial_\rho^2 A_t - \frac{2q^2 A_t}{z^2} (\varphi_1^2 + \varphi_2^2) ,$$

$$0 = \partial_z f \partial_z A_\phi + f \partial_z^2 A_\phi - \frac{\partial_\rho A_\phi}{\rho} + \partial_\rho^2 A_\phi + \frac{2q}{z^2} \varphi_1^2 (n_1 - qA_\phi) + \frac{2q}{z^2} \varphi_2^2 (n_2 - qA_\phi)$$

$$\begin{aligned} 0 = & -\frac{\varphi_1}{\rho^2} (qA_\phi - n_1)^2 + \frac{q^2 A_t^2 \varphi_1}{f} - \frac{m_1^2 \varphi_1}{z^2} - \frac{\epsilon e^{i(n_2-n_1)\phi} \varphi_2}{z^2} - \frac{\eta \varphi_1 \varphi_2^2}{z^2} \\ & + \left(\partial_z f - \frac{2f}{z} \right) \partial_z \varphi_1 + f \partial_z^2 \varphi_1 + \frac{\partial_\rho \varphi_1}{\rho} + \partial_\rho^2 \varphi_1 , \end{aligned}$$

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$n_1 = n_2 = n$

- * Boundary condition (B.C.) near $z \sim 0$

$$\varphi_i(\rho, z)|_{z \rightarrow 0} = \varphi_i^{(1)}(\rho)z^{3-\Delta_i} + \varphi_i^{(2)}(\rho)z^{\Delta_i}, \quad i = 1, 2,$$

$$A_t(\rho, z)|_{z \rightarrow 0} = \mu(\rho) - \varrho(\rho)z,$$

$$A_\phi(\rho, z)|_{z \rightarrow 0} = a_\phi(\rho) + J_\phi(\rho)z.$$

- * B.C. near horizon

At=0, regular b.c. for other fields

- * Numerical parameters:

$$q = L = 1 \quad m_1^2 = -2, m_2^2 = -5/4 \quad \boxed{\eta = 0}$$

$$\Delta_1 = 2, \Delta_2 = 5/2$$

Holographic Superfluid Vortex

* B.C. at vortex core $\rho=0$

$$\varphi|_{\rho \rightarrow 0} = 0, \quad \partial_\rho A_t|_{\rho=0} = 0, \quad A_\phi|_{\rho=0} = 0, \quad \text{for } n \neq 0$$

$$\partial_\rho \varphi|_{\rho=0} = 0 \quad \dots \quad \dots \quad \text{for } n = 0$$

* B.C. at $z=0$

$$\varphi_i^{(1)} \equiv 0$$

* Dirichlet B.C. at z=0

$$A_\phi|_{z=0} = a_\phi(\rho) = \frac{1}{2}\rho^2 B$$

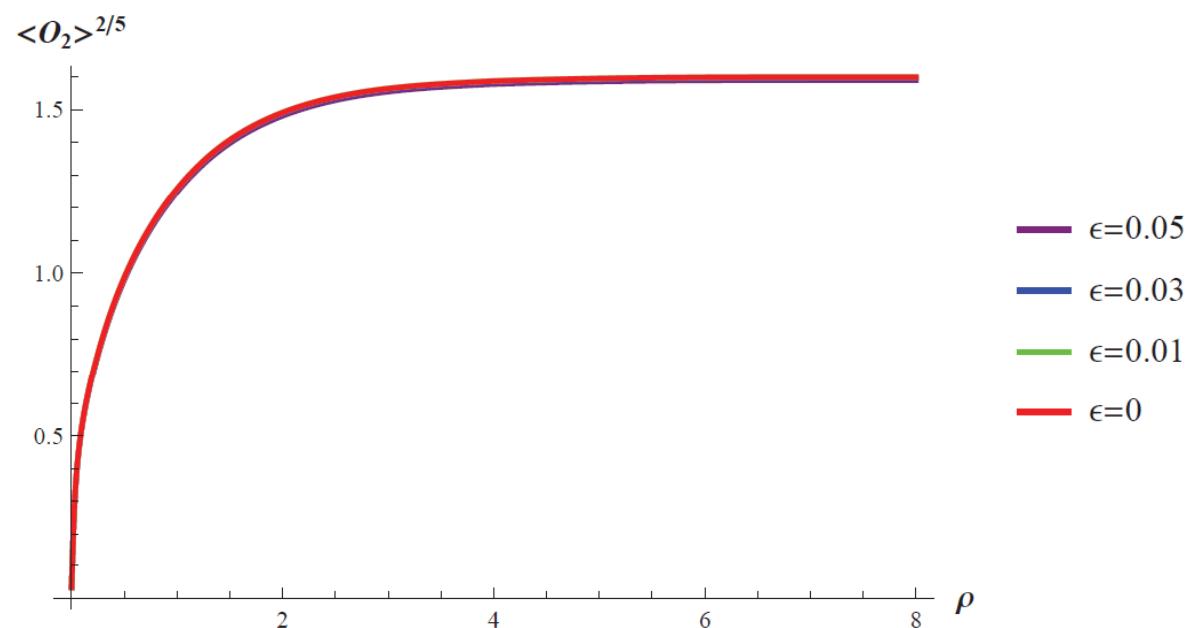
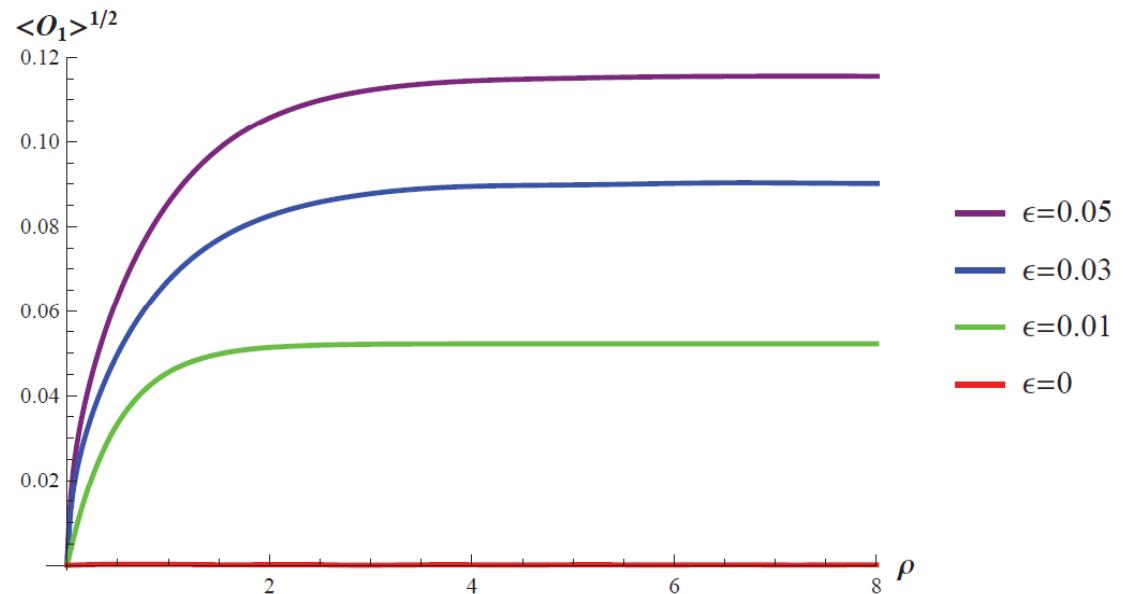
* B.C. at vortex border $\rho=R$

$$A_\phi|_{\rho=R} = \frac{1}{2}BR^2$$

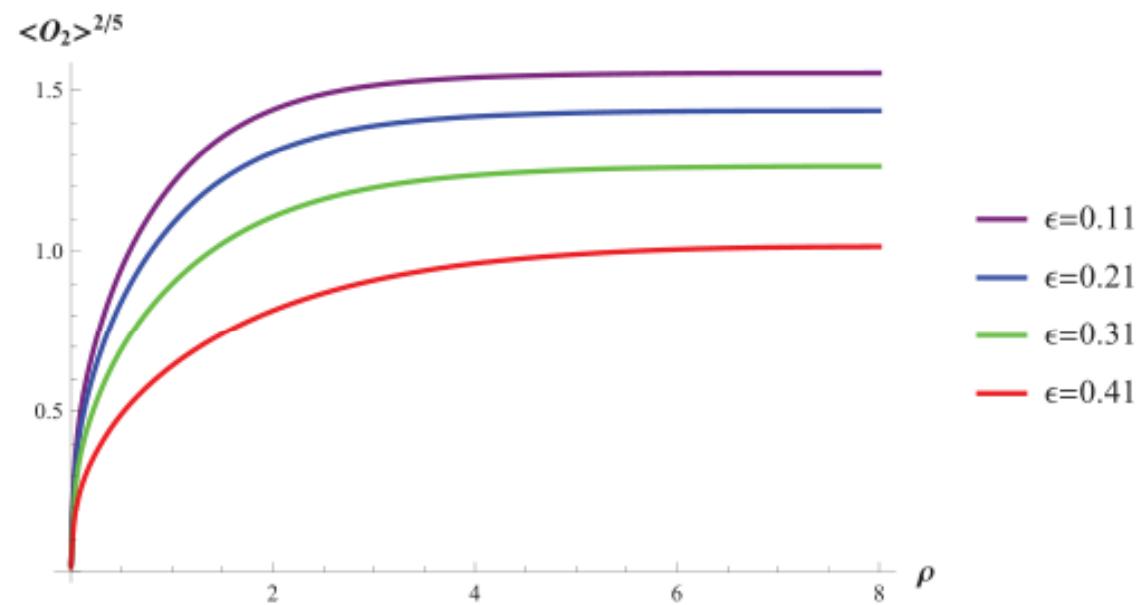
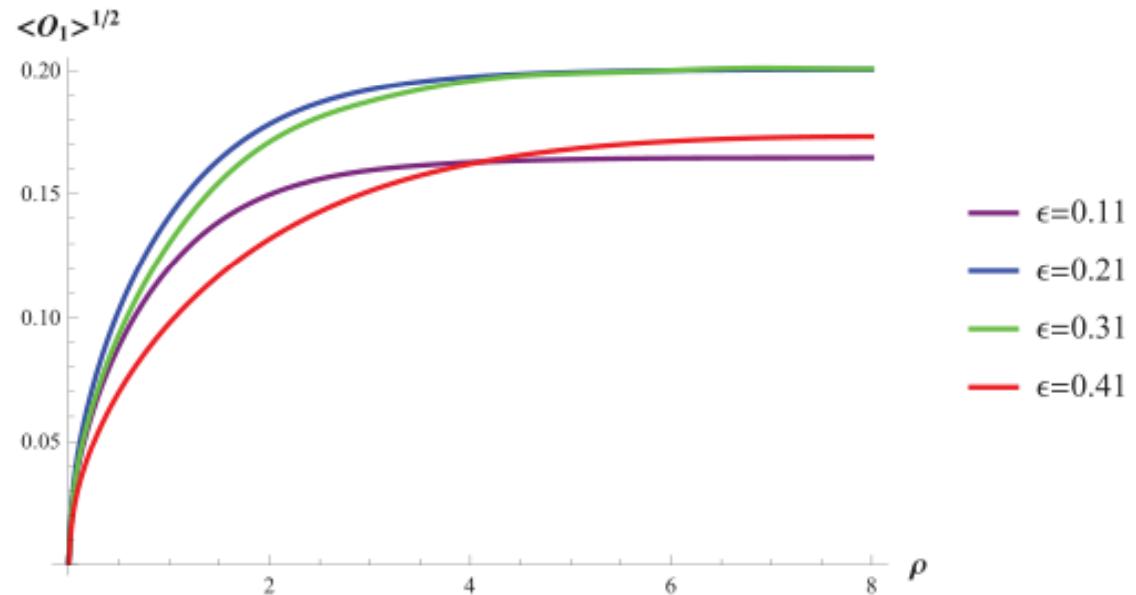
← Angular velocity

Homogeneous b.c. for A_ϕ and φ_i

For small ϵ



For $\epsilon \geq 0.1$



Free Energy

* On-shell action

$$\begin{aligned} S_{os} = & \frac{1}{2\kappa^2} \left(-\frac{1}{2} \right) \int d^4x \partial_a \left[\sqrt{-g} \left(A_b F^{ab} + \psi_1^* \partial^a \psi_1 + \psi_1 \partial^a \psi_1^* + \psi_2^* \partial^a \psi_2 + \psi_2 \partial^a \psi_2^* \right) \right] \\ & + \frac{1}{2\kappa^2} \frac{iq}{2} \int d^4x \sqrt{-g} A_b \left[\psi_1^* \left(\partial^b - iqA^b \right) \psi_1 - \psi_1 \left(\partial^b + iqA^b \right) \psi_1^* \right. \\ & \quad \left. + \psi_2^* \left(\partial^b - iqA^b \right) \psi_2 - \psi_2 \left(\partial^b + iqA^b \right) \psi_2^* \right] \\ & + \frac{\eta}{2\kappa^2} \int d^4x \sqrt{-g} |\psi_1|^2 |\psi_2|^2. \end{aligned}$$

* Counter term

$$S_{ct} = \frac{-1}{2\kappa^2} \int d^3x \sqrt{-\gamma} (\psi_1 \psi_1^*) + \frac{-1/2}{2\kappa^2} \int d^3x \sqrt{-\gamma} (\psi_2 \psi_2^*)$$

Free Energy

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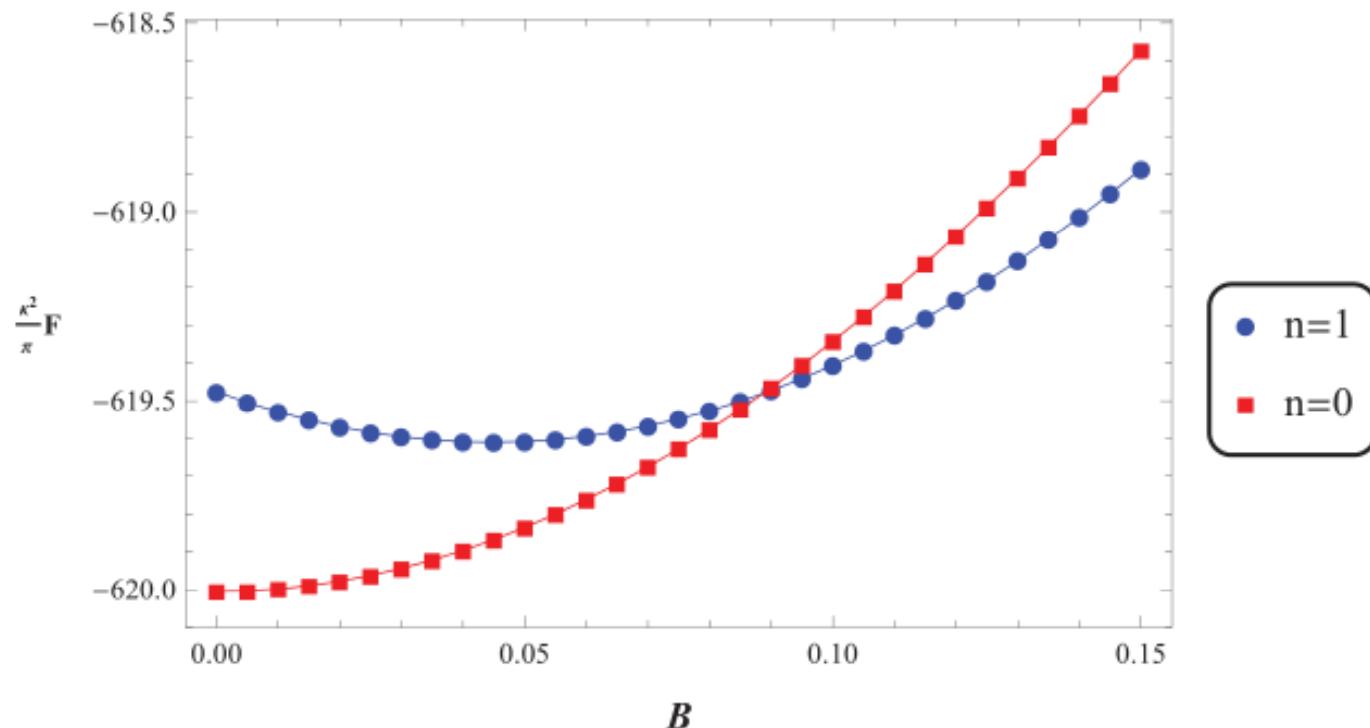
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$$S_{ct} = \frac{-1}{2\kappa^2} \int d^3x \sqrt{-\gamma} (\psi_1 \psi_1^*) + \frac{-1/2}{2\kappa^2} \int d^3x \sqrt{-\gamma} (\psi_2 \psi_2^*)$$

* Regularized free energy

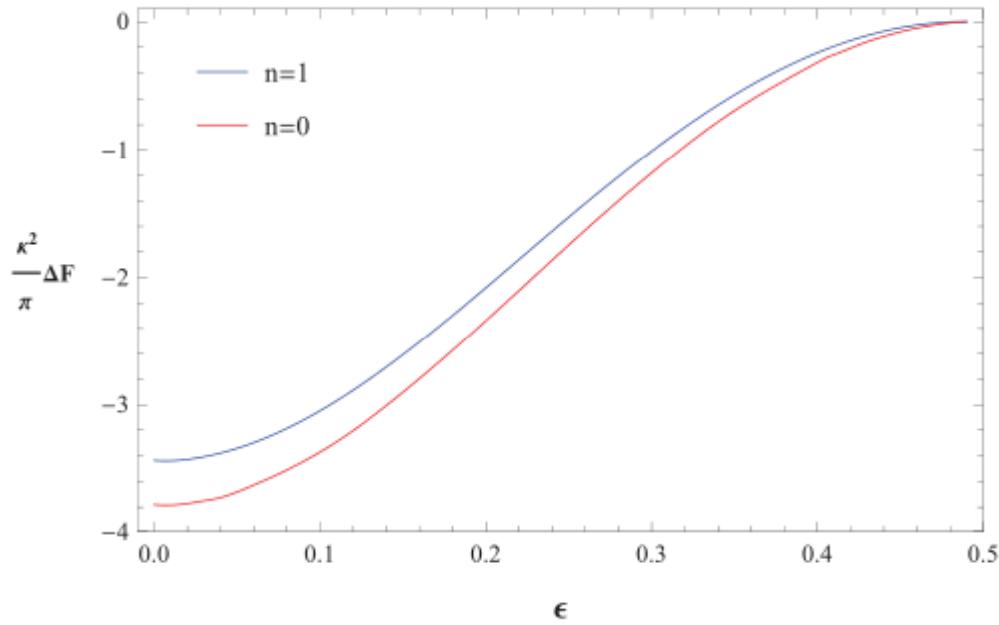
$$\begin{aligned}
F &= -TS_{reg.} = -T(S_{os} + S_{ct}) \\
&= \frac{-T}{2\kappa^2} \int dt d\phi \left[\int d\rho \rho \left(\frac{\varrho\mu}{2} + \frac{BJ_\phi}{4} \right) \Big|_{z=0} - \int dz \frac{A_\phi \partial_\rho A_\phi}{2\rho} \Big|_{\rho=R} \right. \\
&\quad \left. + \int dz d\rho \frac{\rho}{z^2} \left(-\frac{q^2 A_t^2 (\varphi_1^2 + \varphi_2^2)}{f} - \frac{q A_\phi}{\rho^2} (\varphi_1^2 (n_1 - q A_\phi) + \varphi_2^2 (n_2 - q A_\phi)) + \frac{\eta \varphi_1^2 \varphi_2^2}{z^2} \right) \right]
\end{aligned}$$

* Critical angular velocity B_{cl}

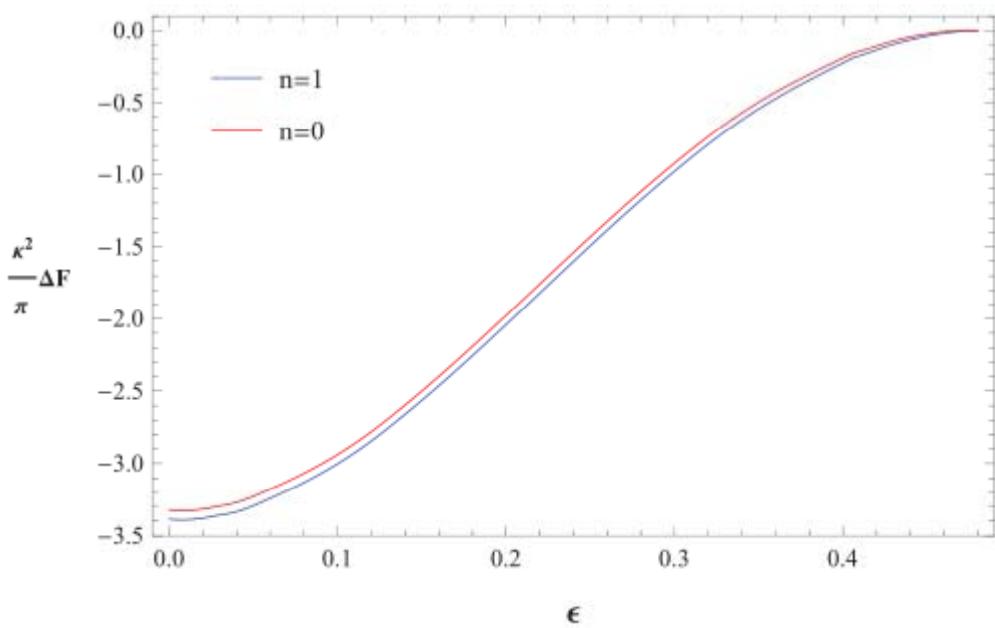


* For n=1

$$\Delta F = F(\varphi_i \neq 0) - F(\varphi_i = 0)$$



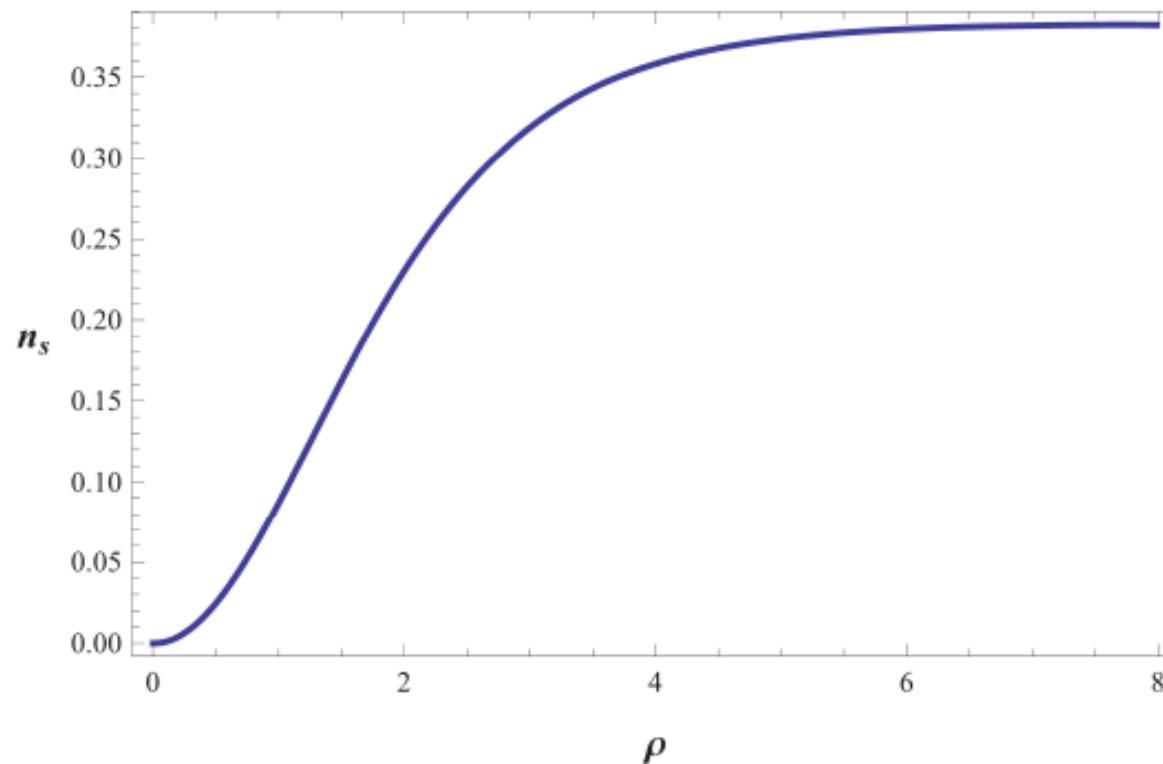
$$B = 0.03 < B_{c1}$$



$$B = 0.1 > B_{c1}$$

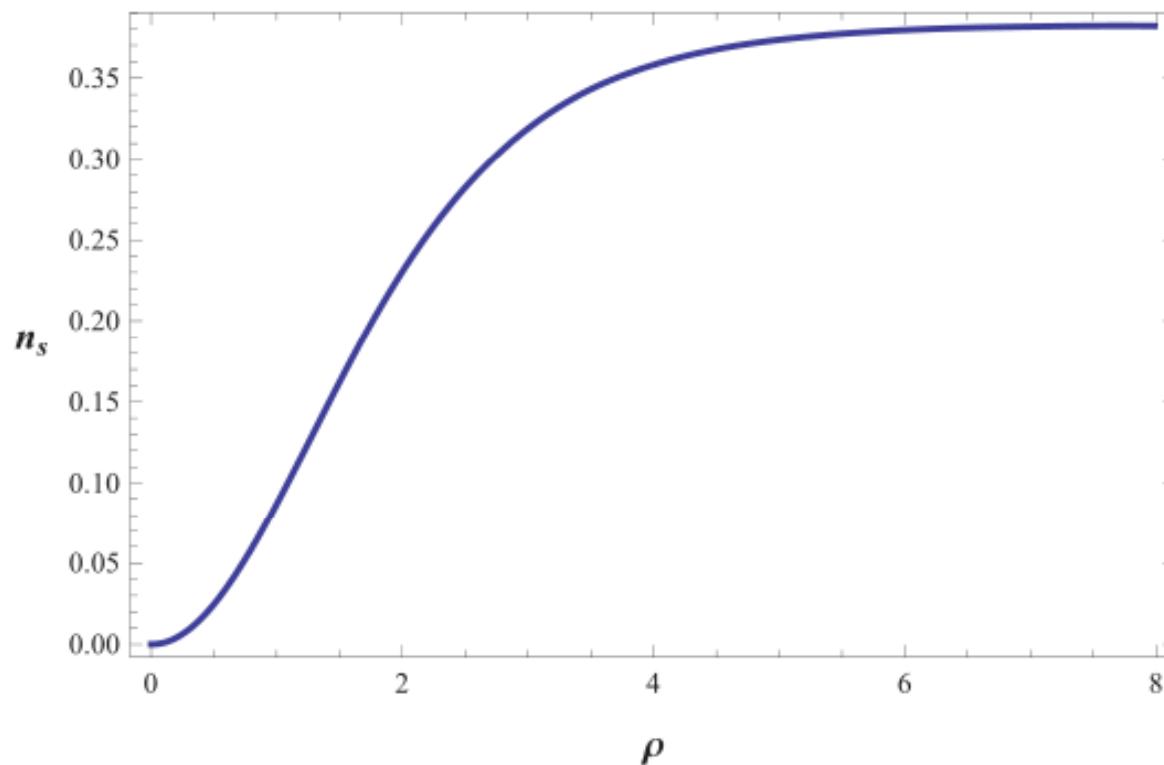
* Superfluid density and current

$$n_s = \frac{J_\phi}{n - a_\phi}$$



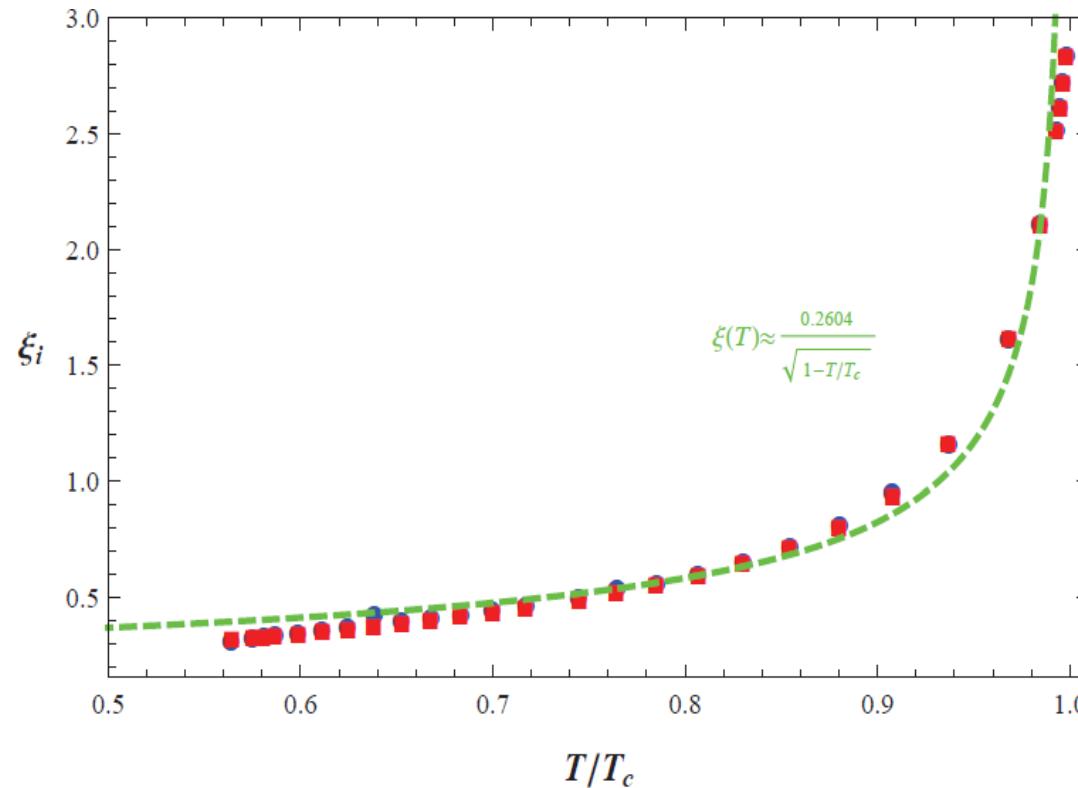
* Superfluid density and current

$$n_s = \frac{J_\phi}{n - a_\phi} \quad \xleftarrow{\text{red arrow}} \quad v_\phi = \nabla_\phi \varphi_i - A_\phi$$



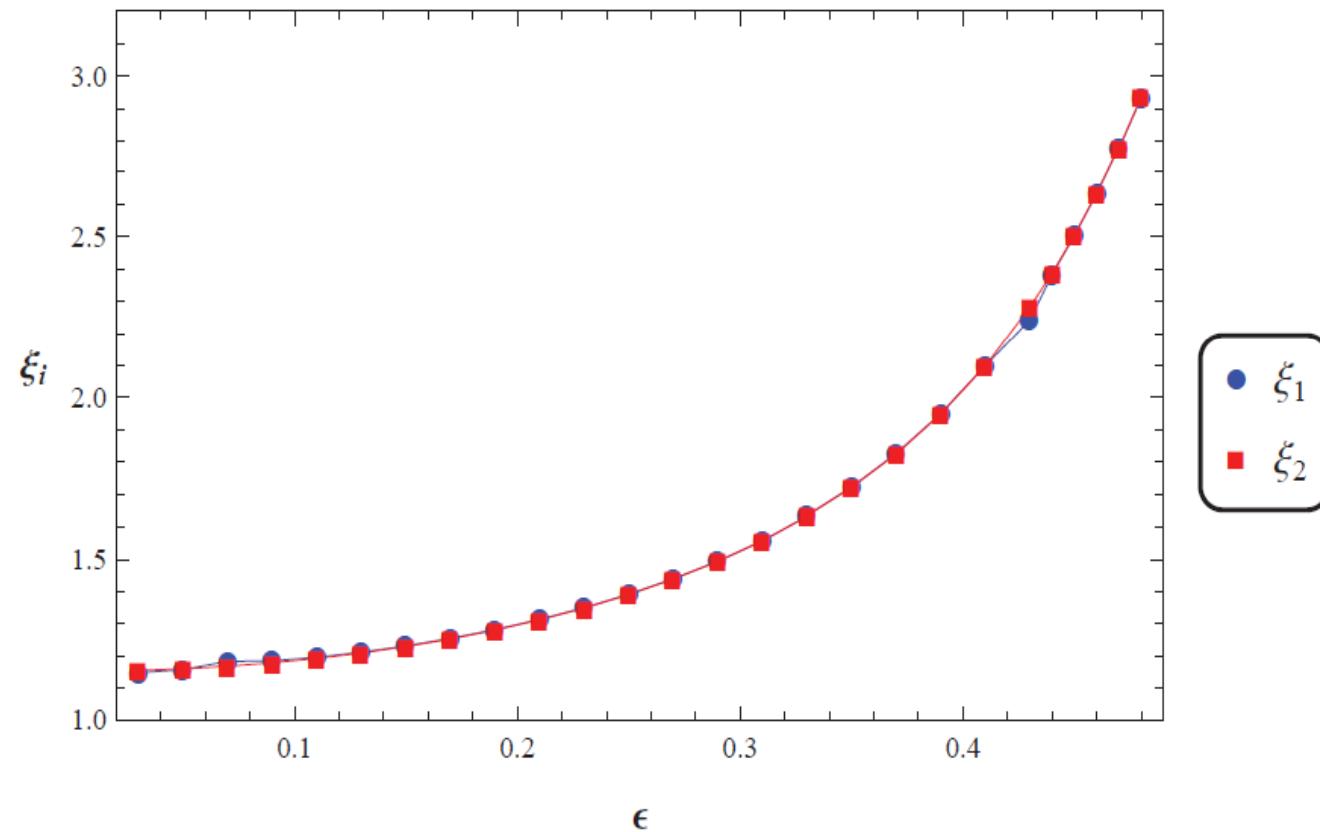
* Coherence length vs.T

$$\langle O_i(\rho) \rangle = O_i(\infty) \tanh \left(\frac{\rho}{\sqrt{2}\xi_i} \right)$$

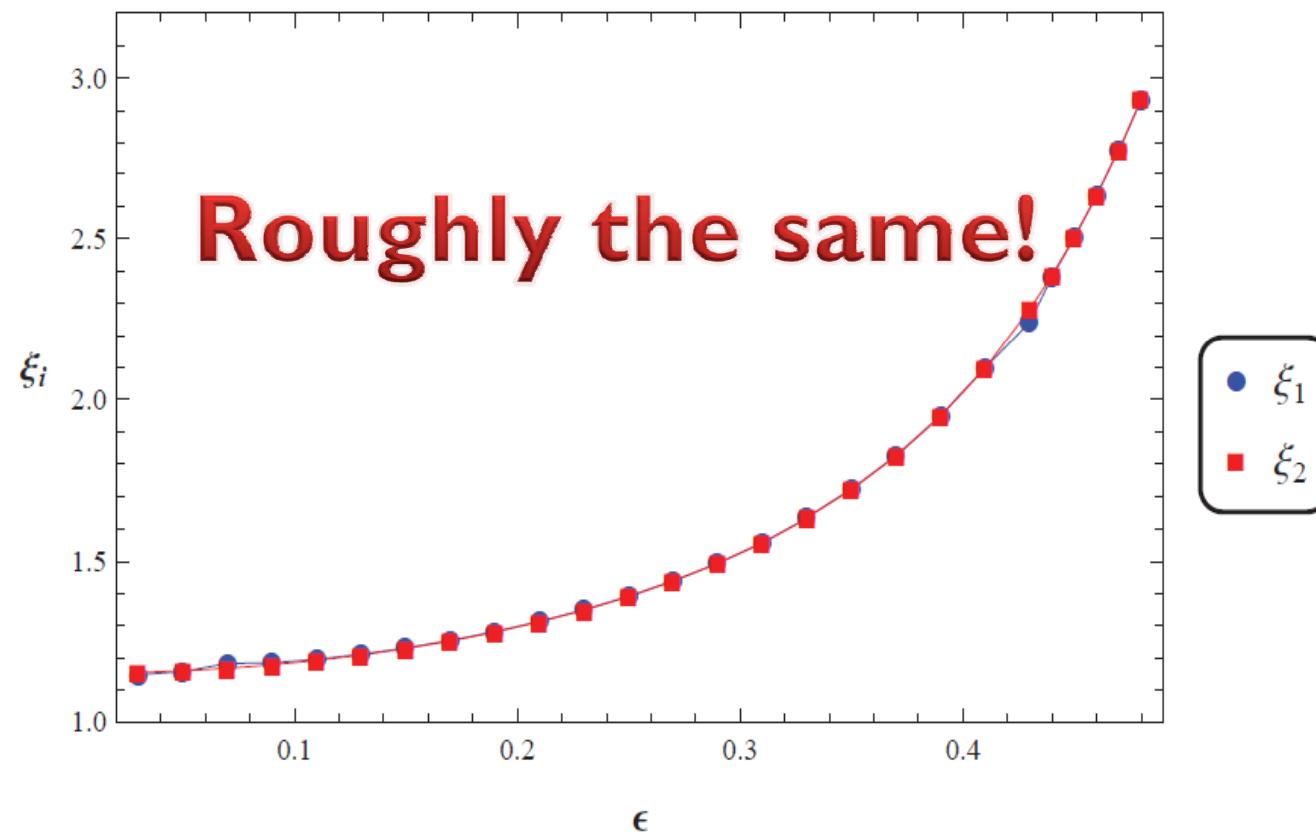


$$\xi_i(T) = 0.2604(1 - T/T_c)^{-1/2}$$

* Dependence on Josephson coupling



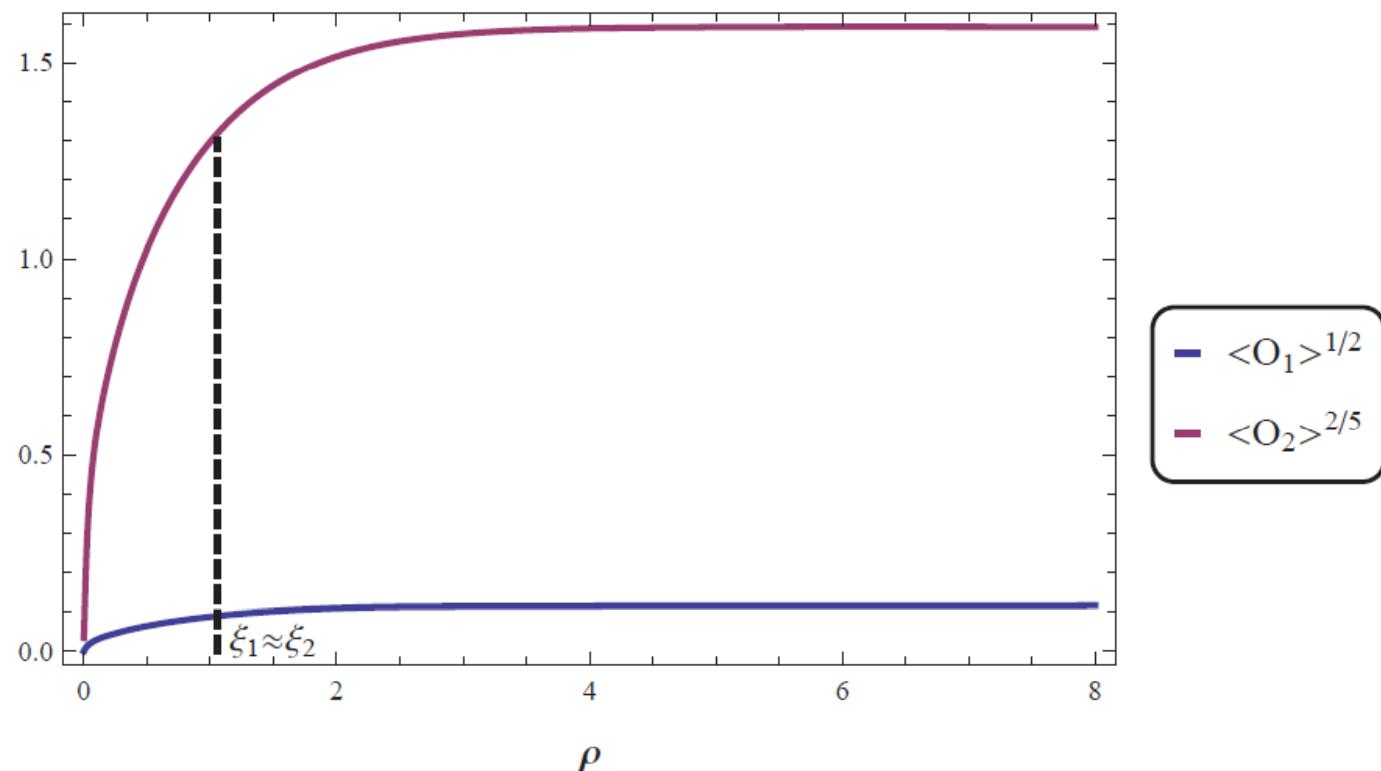
* Dependence on Josephson coupling



Holographic Superconductor Vortex

* Neumann B.C. at $z=0$

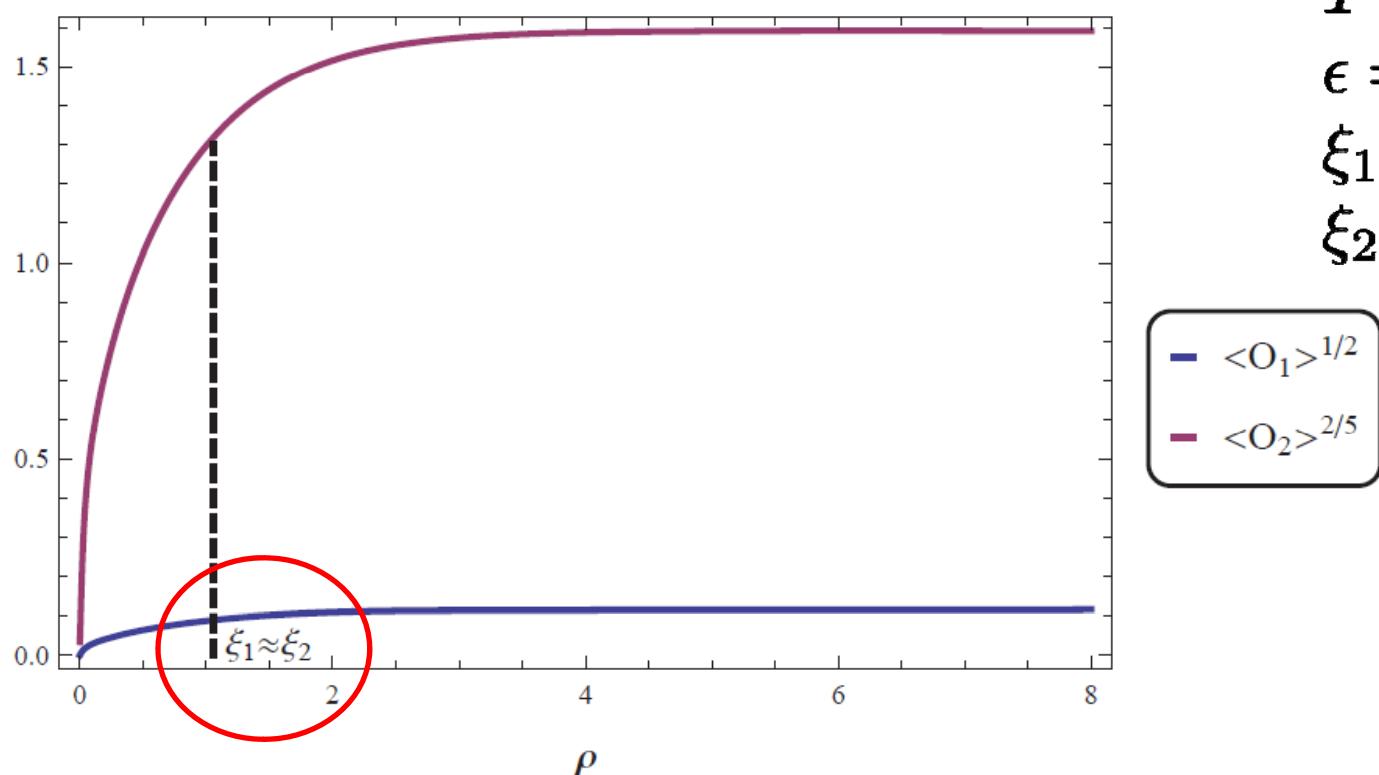
$$\partial_z A_\phi(\rho, z) = J_\phi(\rho) = 0$$



Holographic Superconductor Vortex

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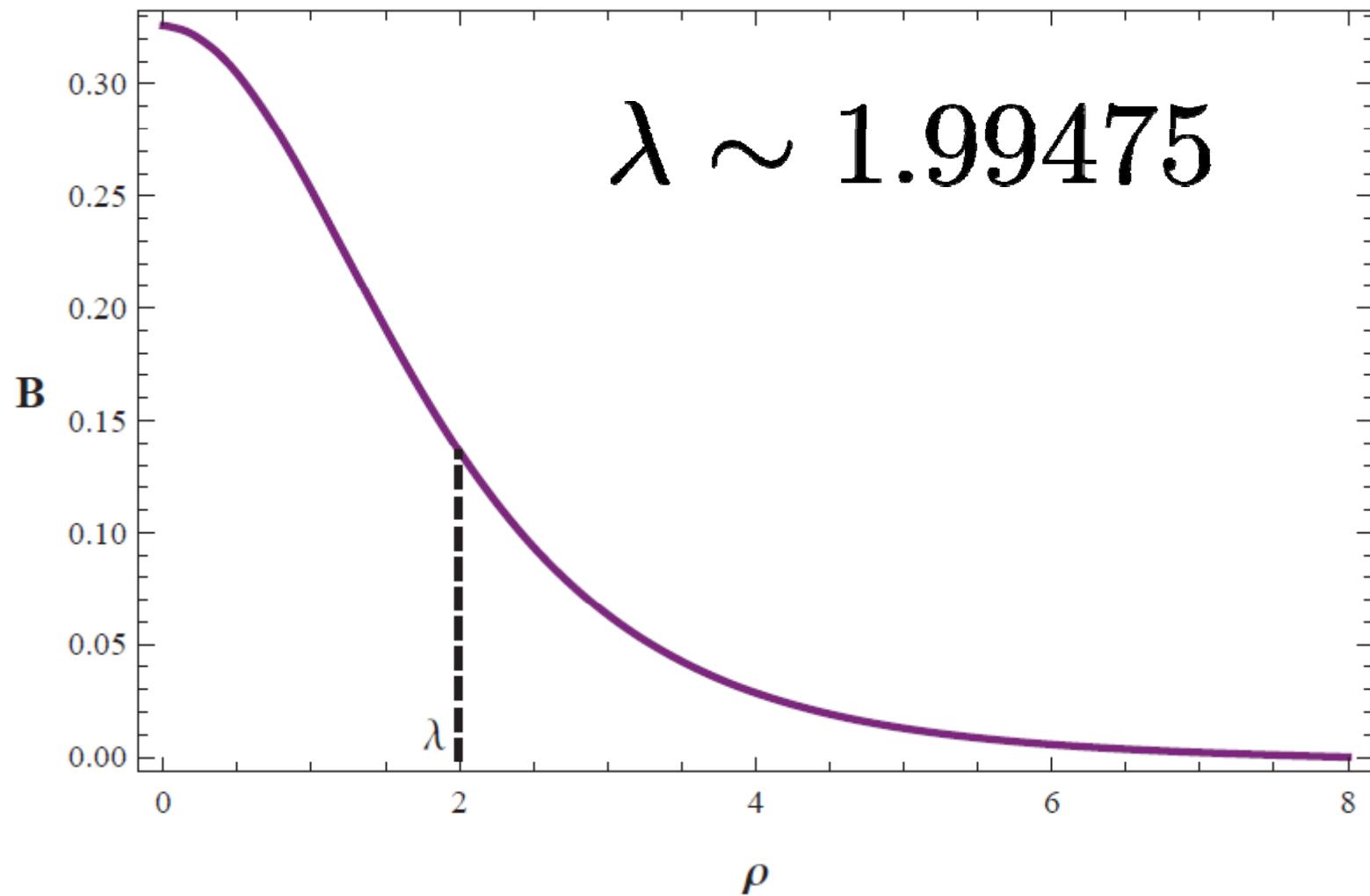
$$\partial_z A_\phi(\rho, z) = J_\phi(\rho) = 0$$



$$\begin{aligned} T &= 0.937T_c \\ \epsilon &= 0.05 \\ \xi_1 &\sim 1.05687 \\ \xi_2 &\sim 1.01591 \end{aligned}$$

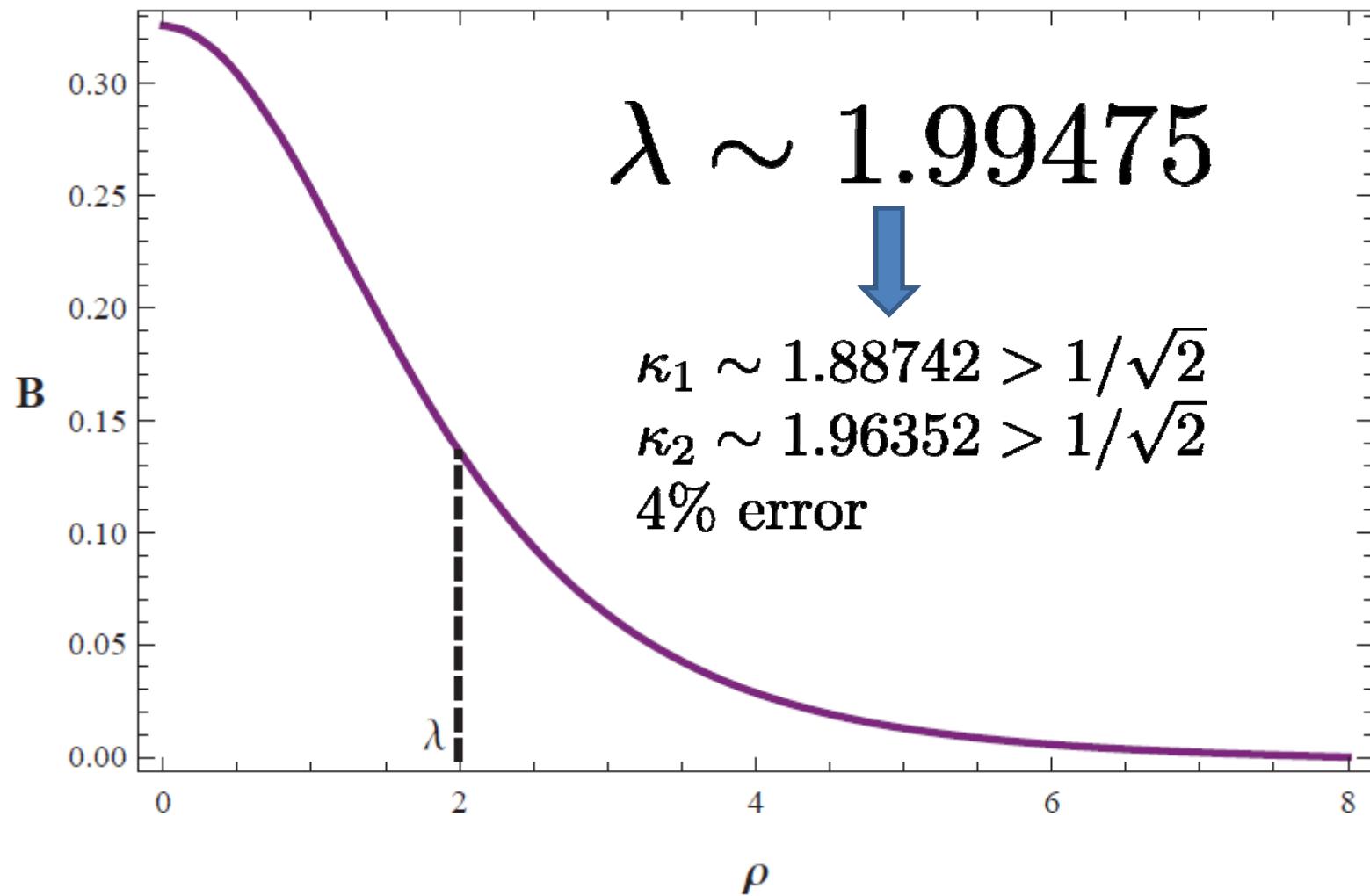
* Magnetic field B

$$B = b e^{-\rho/\lambda}$$



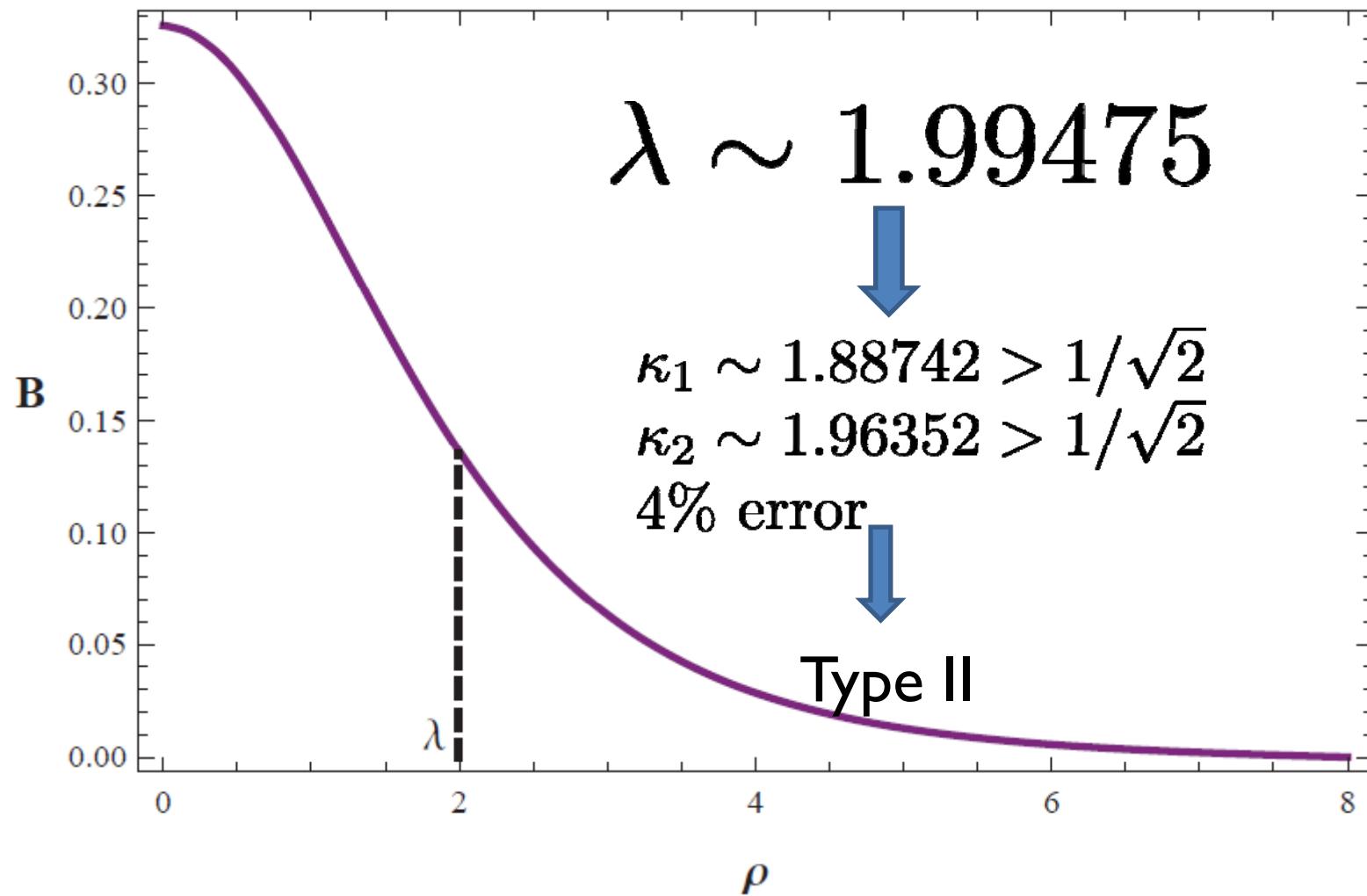
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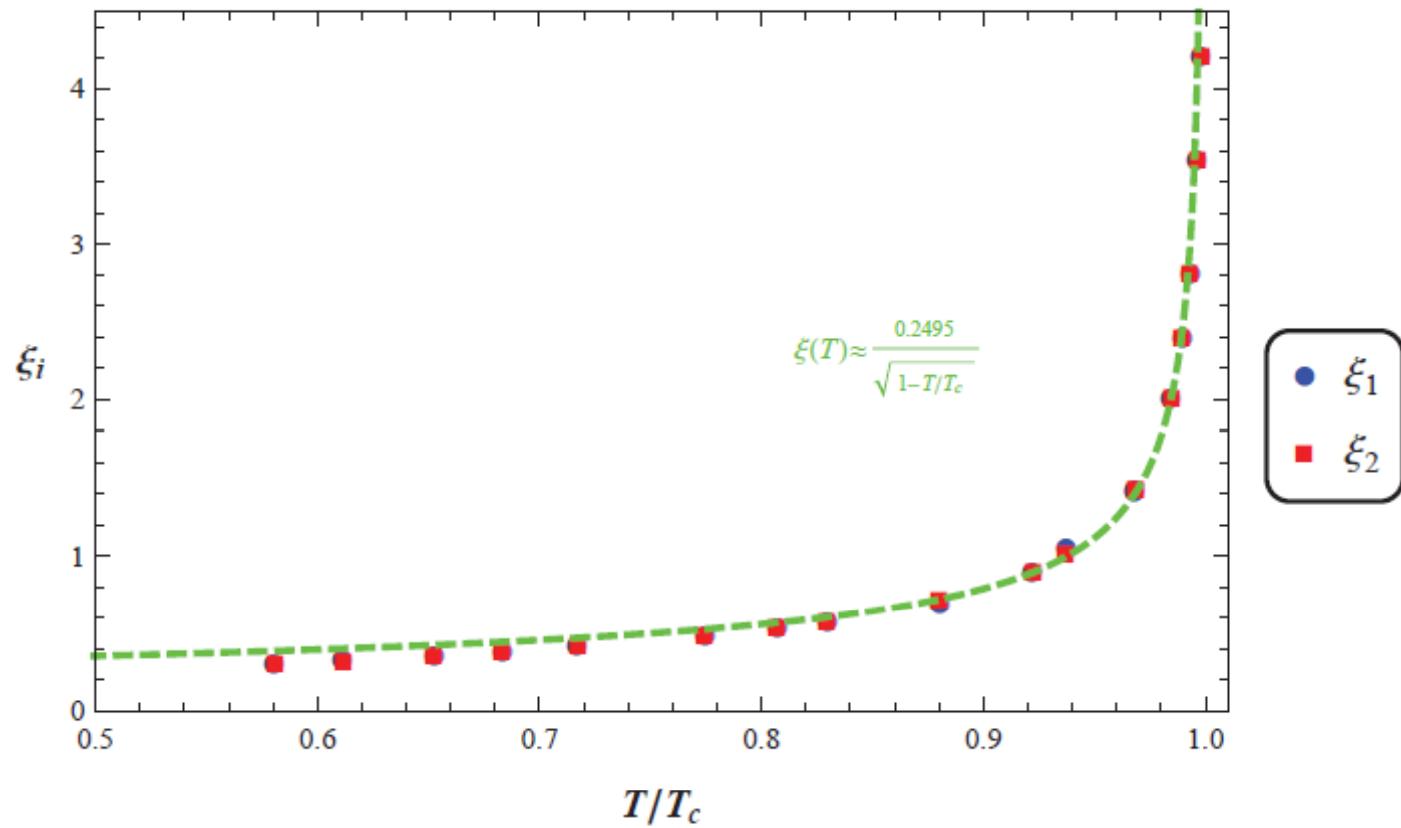


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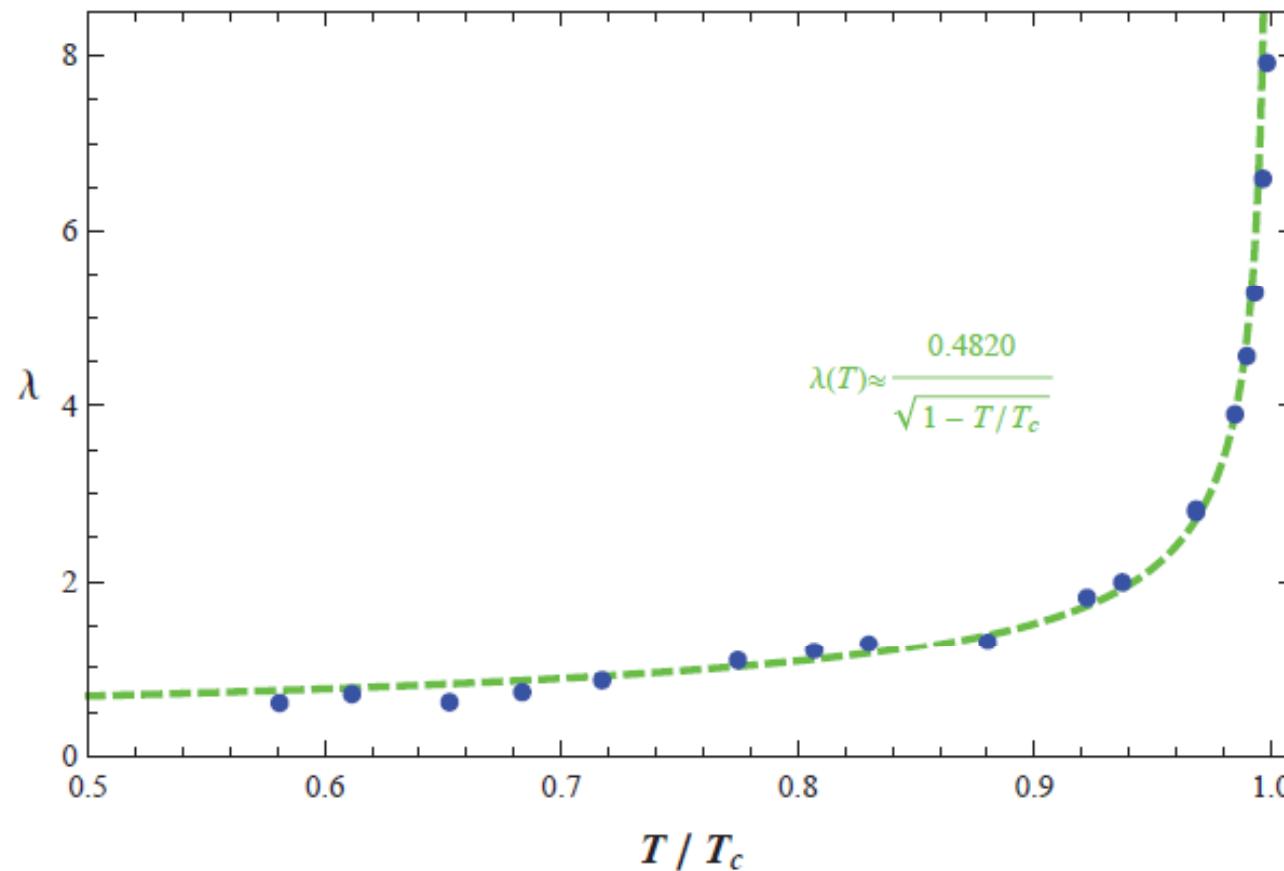
$$B = b e^{-\rho/\lambda}$$



* Coherence length vs. T



* Penetration length vs.T



$$\kappa_{1,2} > 1/\sqrt{2}$$

SUMMARIES & OUTLOOK

Summaries

- Studied the angular velocity response of superfluid vortex;
- The magnetic response of superconductor vortex;
- Effects of Josephson coupling and temperature on vortex;
- By calculating the coherence length and penetration length, we get the type II superconductor vortex.

Outlook

Type I.5

- Meissner
- Semi-Meissner
- Vortex
- Normal



$$\kappa_1 < 1/\sqrt{2}, \quad H$$
$$\kappa_2 > 1/\sqrt{2}$$

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- Normal

$$\kappa_1 < 1/\sqrt{2},$$

$$\kappa_2 > 1/\sqrt{2}$$



PRL 102, 117001 (2009)

Selected for a Viewpoint in Physics
PHYSICAL REVIEW LETTERS

week ending
20 MARCH 2009



Type-1.5 Superconductivity

Victor Moshchalkov,^{1,*} Mariela Menghini,¹ T. Nishio,¹ Q. H. Chen,¹ A. V. Silhanek,¹ V. H. Dao,¹ L. F. Chibotaru,¹ N. D. Zhigadlo,² and J. Karpinski²

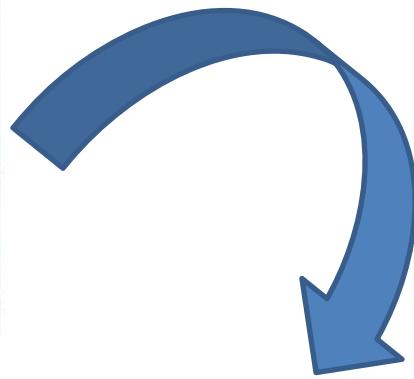
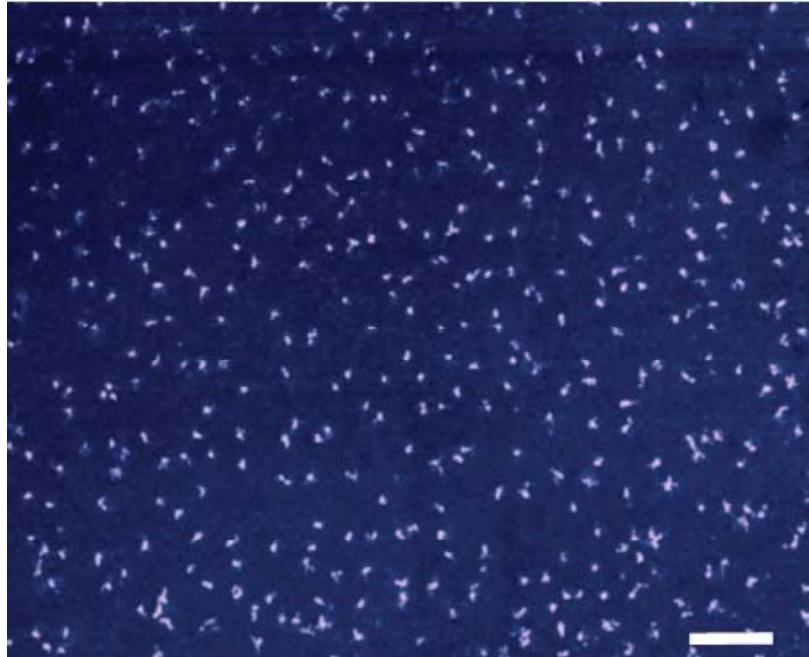
¹INPAC-Institute for Nanoscale Physics and Chemistry, Katholieke Universiteit Leuven,
Celestijnenlaan 200 D, B-3001 Leuven, Belgium

²Laboratory for Solid State Physics, ETH Zürich, 8093-Zurich, Switzerland
(Received 22 December 2008; published 16 March 2009)

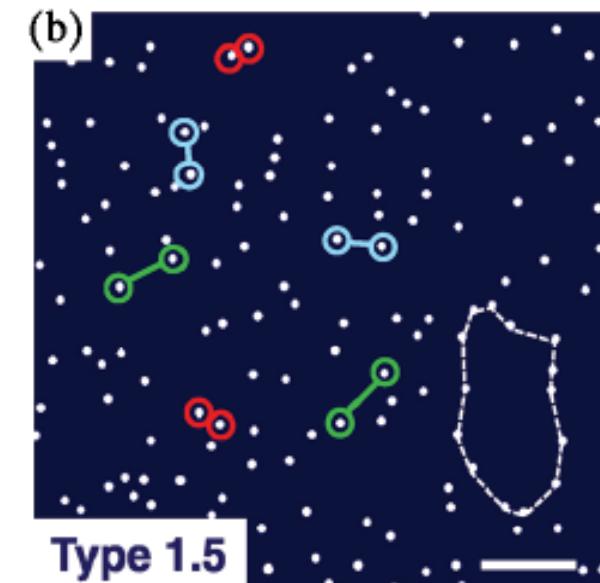
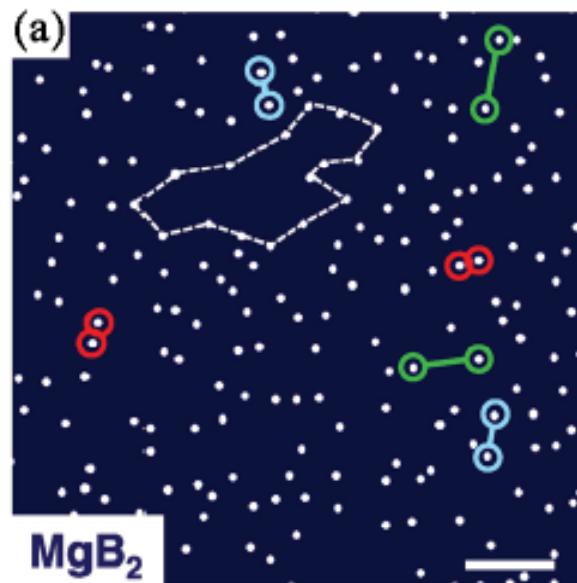
We demonstrate the existence of a novel superconducting state in high quality two-component MgB₂ single crystalline superconductors where a unique combination of both type-I ($\lambda_1/\xi_1 < 1/\sqrt{2}$) and type-II ($\lambda_2/\xi_2 > 1/\sqrt{2}$) superconductor conditions is realized for the two components of the order parameter. This condition leads to a vortex-vortex interaction attractive at long distances and repulsive at short distances, which stabilizes unconventional stripe- and gossamerlike vortex patterns that we have visualized in this type-1.5 superconductor using Bitter decoration and also reproduced in numerical simulations.

$$\kappa_\pi = 0.66 < 1/\sqrt{2} \quad (\text{type I})$$
$$\kappa_\sigma = 3.68 > 1/\sqrt{2} \quad (\text{type II})$$

H



Vortex clusters



Features of type I.5

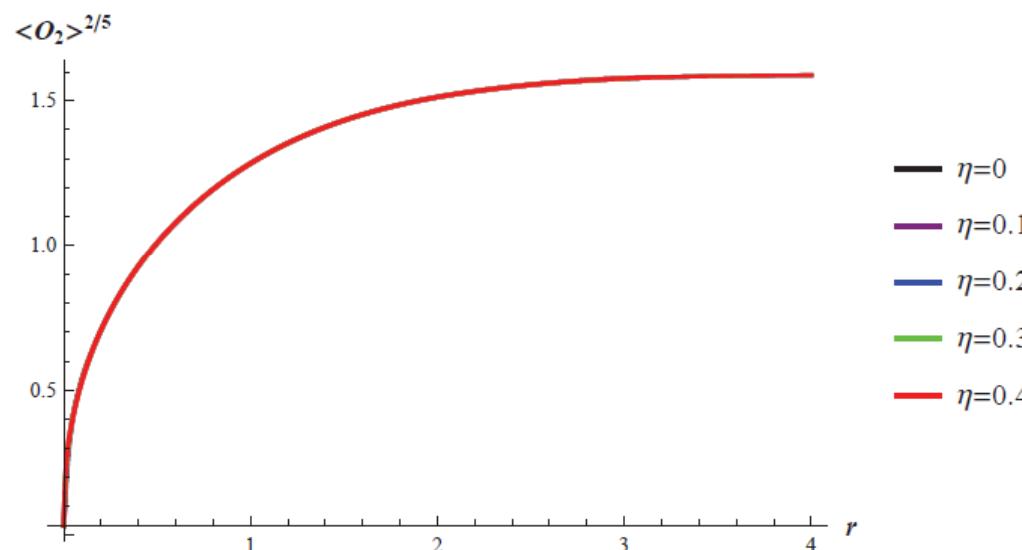
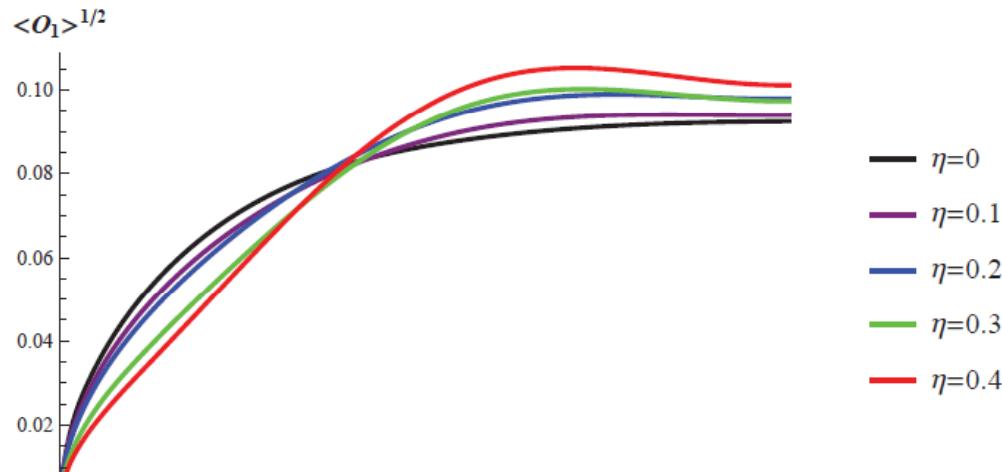
Long-range attractive , short-range repulsive

Semi-Meissner: mixtures of Meissner and vortex

Non-axialsymmetric vortex

Vortex clusters

- The effects of η



Thank you!