

Hidden dark sector and dark matter



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arXiv:0904.1745 [hep-ph]

arXiv:0907.2918 [hep-ph]

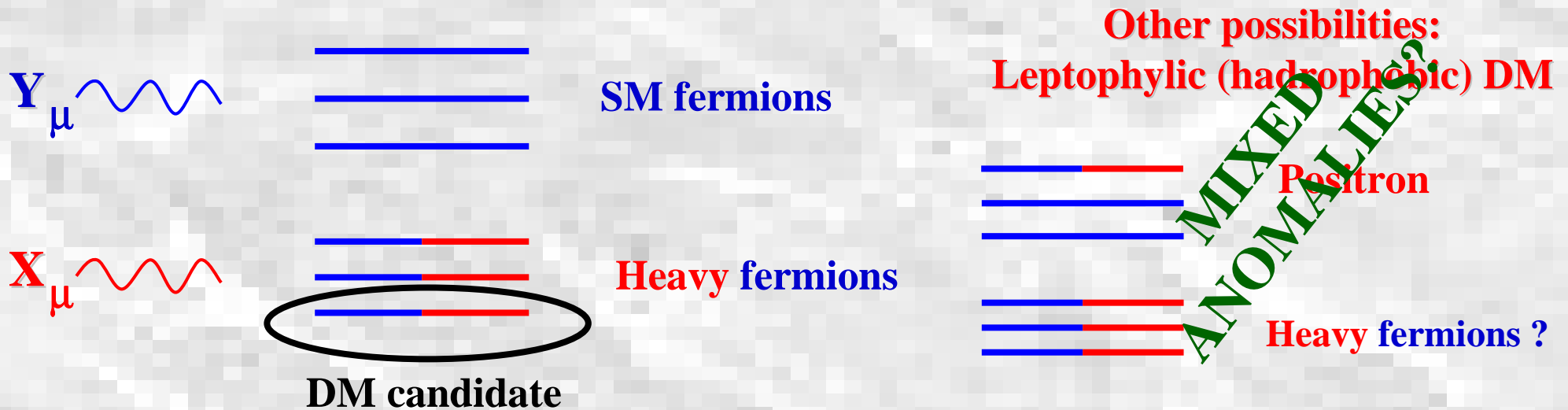
Madrid DM Workshop, September 16th 2009

Extra U(1) models

Question : is it possible to see an invisible gauge boson X?

What is an invisible X?

A boson that does not couple with SM particles



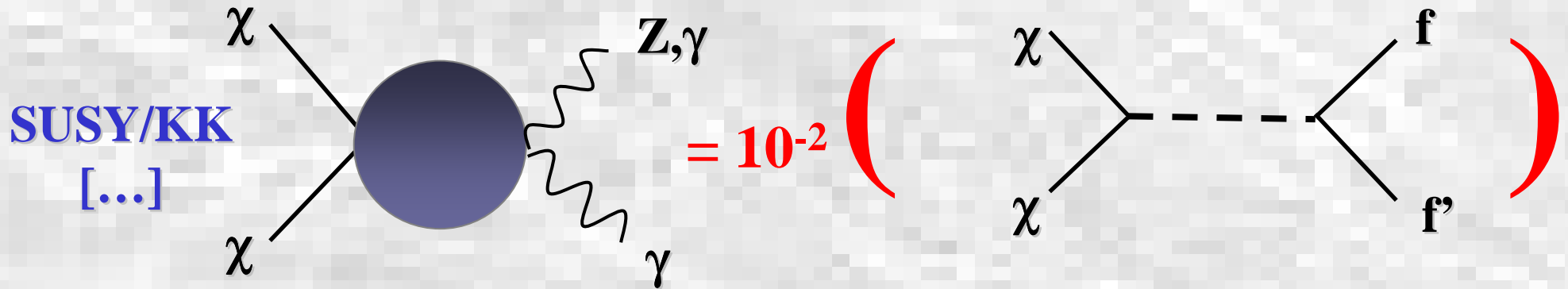
A Y-X kinetic mixing generated at loop level

$$\mathcal{L}_{mix} = \delta F^{Y\mu\nu} F^X_{\mu\nu}$$

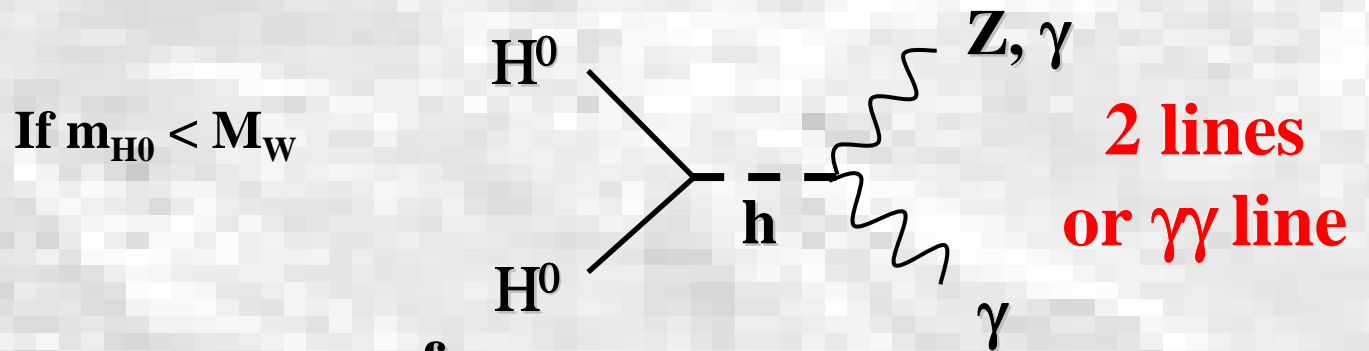
[Weiner, 08]



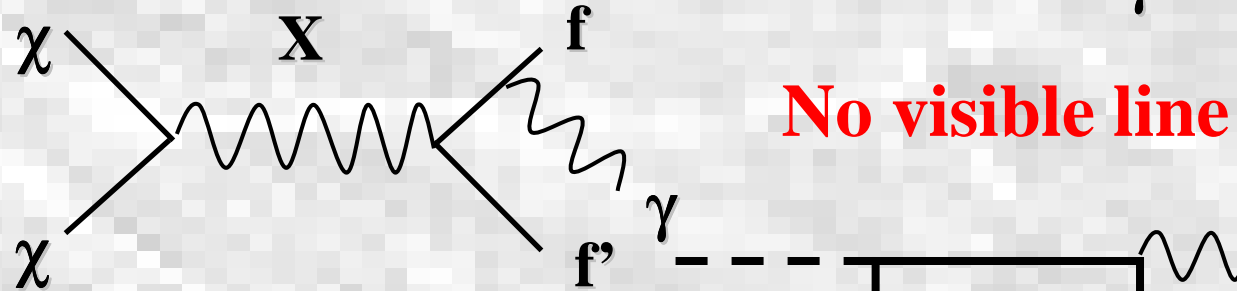
γ ray lines



Inert Higgs Doublet
[Gustafsson et al. 07]

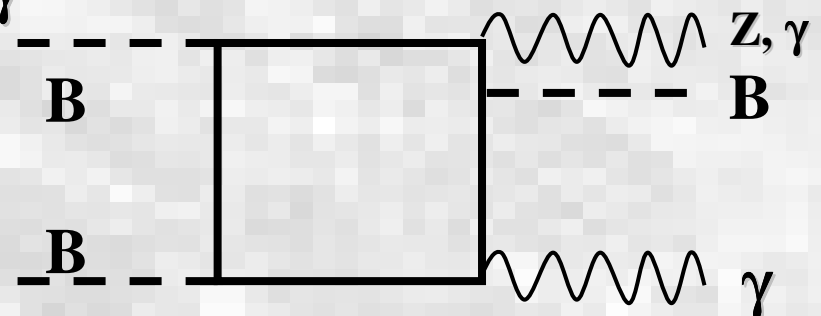


Visible X
[Cheung et al. 07]
[Baek & Ko, 08]

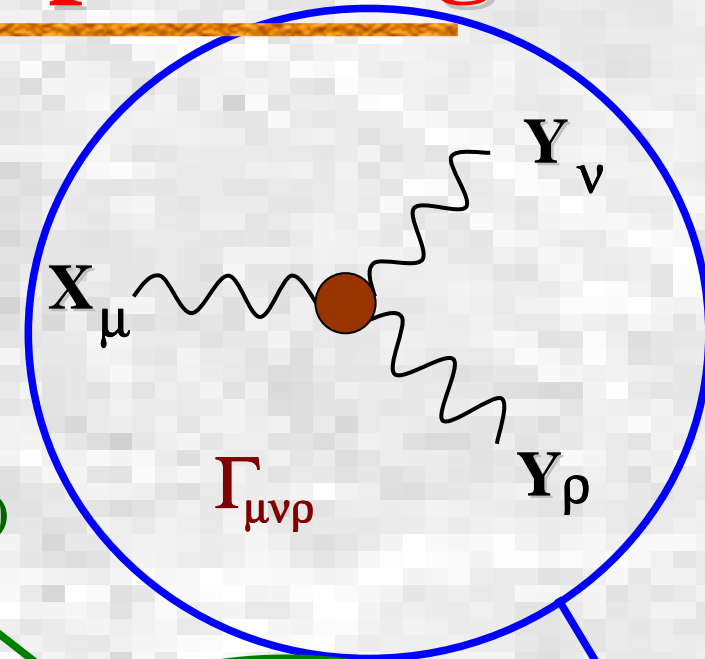
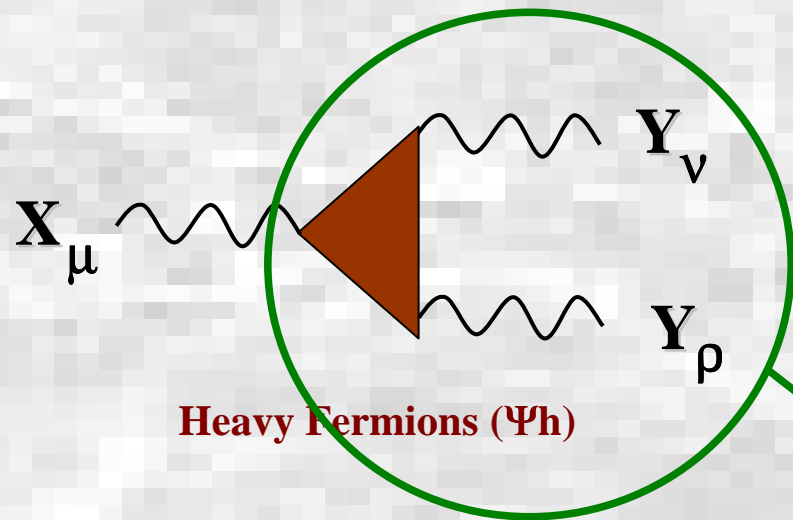


Chiral Square
[Bertone et al. 09]

3 visible lines



Anomalies and loops triangle



$U'(1)$

| Ψ | $U(1)$ | $U'(1)$ |
|-------------|----------|---------|
| Ψ_{SM} | X_{SM} | 0 |
| Ψ_h | X_h | X'_h |

$$\mathcal{L} \xrightarrow{U'(1)} \mathcal{L} + \lambda \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^Y F_{\rho\sigma}^Y + \mathcal{L}_3$$

$$X_\mu \longrightarrow X_\mu + d_\mu \lambda$$

$$a \longrightarrow a - M_X \lambda$$

$$X_\mu + M_X d_\mu a \longrightarrow X_\mu + M_X d_\mu a$$

(Stuckelberg Lagrangian)

The Lagrangian :

Green-Schwarz mechanism

(Intersecting Brane Models)

$$\mathcal{L}_{inv} = F^{Y\mu\nu} F^Y_{\mu\nu} - (d_\mu a - M_X X_\mu)^2 - i \bar{\Psi}_h \gamma^\mu D_\mu \Psi_h$$

$$\mathcal{L}_{var} = \underbrace{B a \varepsilon^{\mu\nu\rho\sigma} F^Y_{\mu\nu} F^Y_{\rho\sigma}}_{\text{Peccei-Quinn terms}} + \underbrace{C \varepsilon^{\mu\nu\rho\sigma} X_\mu Y_\nu F^Y_{\rho\sigma}}_{\text{Chern-Simons terms}}$$

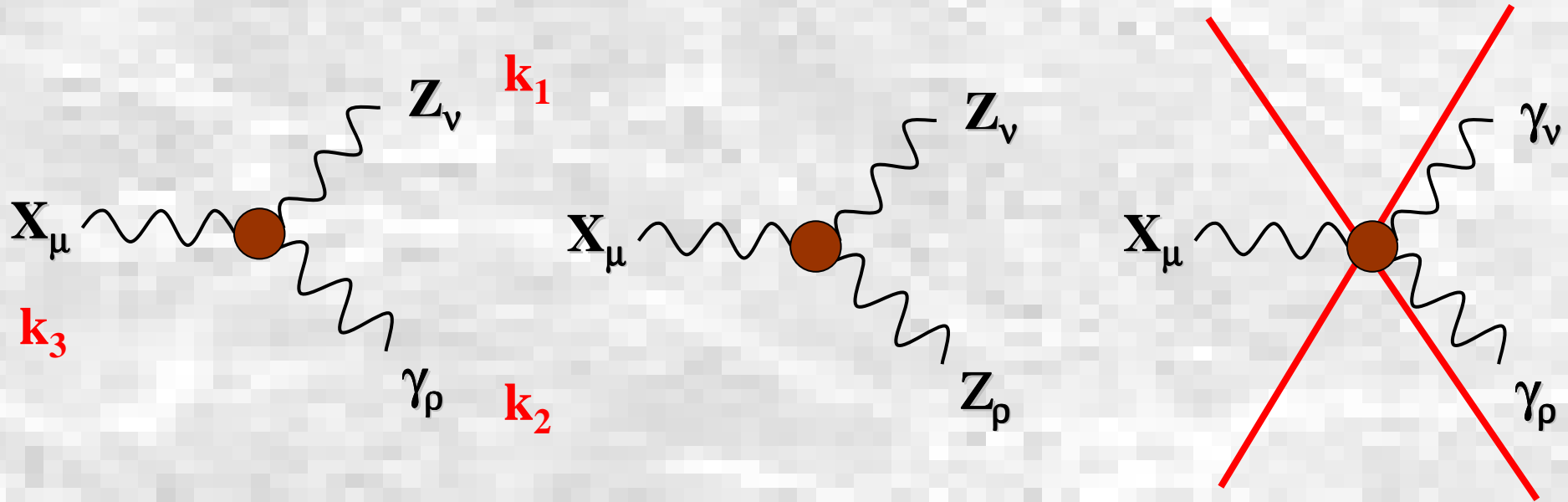
Peccei-Quinn terms

Chern-Simons terms

$$\delta \mathcal{L}_{var} = - \delta \left(\begin{array}{c} X_\mu \text{ wavy line} \\ \text{triangle} \\ Y_\nu \text{ wavy line} \\ Y_\rho \text{ wavy line} \end{array} \right)$$

Heavy Fermions (Ψ_h)

Effective couplings : $\mathcal{L}_{eff} = \mathcal{L}_{loops} + \mathcal{L}_{var}$



$$\Gamma_{\mu\nu\rho}^{\alpha} = t^{\alpha} \{ A_1 \epsilon^{\mu\nu\rho\sigma} k_{2\sigma} - A_2 \epsilon^{\mu\nu\rho\sigma} k_{1\sigma} + B_1 k_{2\nu} \epsilon^{\mu\rho\sigma\tau} k_{2\sigma} k_{1\tau} + B_2 k_{1\nu} \epsilon^{\mu\rho\sigma\tau} k_{2\sigma} k_{1\tau} + B_3 k_{2\rho} \epsilon^{\mu\nu\sigma\tau} k_{2\sigma} k_{1\tau} + B_4 \epsilon^{\mu\nu\sigma\tau} k_{2\sigma} k_{1\tau} + C k_{3\mu}/k_3^2 \epsilon^{\nu\rho\sigma\tau} k_{2\sigma} k_{1\tau} + D \epsilon^{\mu\nu\rho\sigma} (k_{2\sigma} - k_{1\sigma}) \}$$

Peccei-Quinn

Chern-Simons

[Dudas, Kiritsis 09]

[YM, 09]

[Kumar, Wells 08]

$\delta \mathcal{L}_{eff} = 0$ 3 Ward identities + $(\mathbf{k}_1; \mathbf{k}_2)$ symmetries
 \rightarrow the vertex can be express as function of $|B_2 - B_1| = 1/\Lambda_X^2$

With $B_1, B_2 =$ computable loops integrals

Cc : only 3 parameters : Λ_X [$\langle S \rangle$] ; M_X [g_X] ; $M\chi$ [Y_{heavy}]

Interpretation as higher dimensional operators

[Antoniadis 09; Dudas, YM, Romagnoni 09]

$$\mathcal{L}_1 = \frac{1}{M^2} * \{ b \text{Tr}[F^X F^Y F^Y] + c \varepsilon^{\mu\nu\rho\sigma} (\mathcal{D}_\mu \mathbf{a}) (\mathbf{D}_\nu \mathbf{H})^+ F^Y_{\rho\sigma} \mathbf{H} \},$$

with

$$\mathcal{D}_\mu \mathbf{a} = \mathbf{d}_\mu \mathbf{a} - g' \mathbf{X}_\mu ; \quad \mathbf{D}_\nu = \mathbf{d}_\nu - i g \mathbf{Y}_\nu - i g' \mathbf{X}_\nu$$

Masses suppression coming from the fermions
which decouple after $U'(1)$ breaking

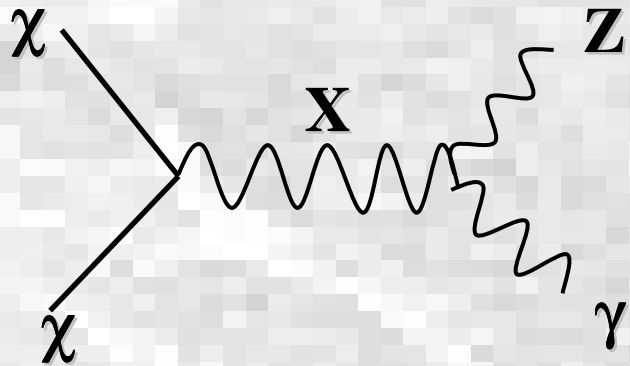
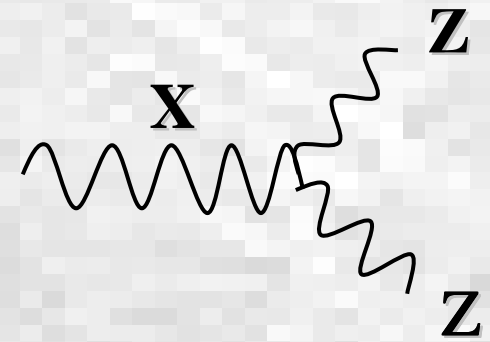
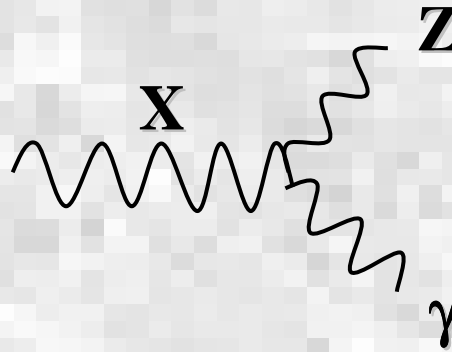
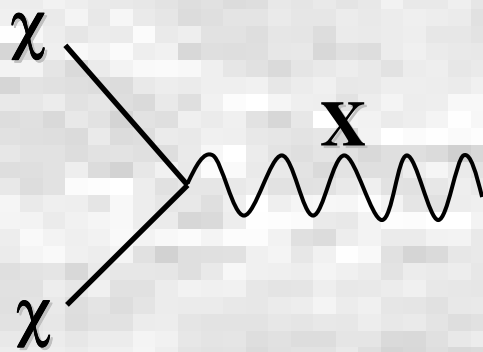
Equivalent to the D'Hoker-Farhi term

$$\{ 1/(\mathbf{H}^+ \mathbf{H}) \varepsilon^{\mu\nu\rho\sigma} (\mathcal{D}_\mu \mathbf{a}) (\mathbf{D}_\nu \mathbf{H})^+ F^Y_{\rho\sigma} \mathbf{H} \} \text{ for SM}$$

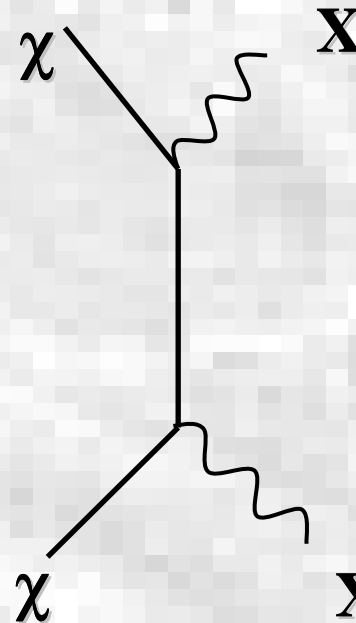
Remark :if two Z' are present, we can build an unsuppressed operator,

$$\varepsilon^{\mu\nu\rho\sigma} (\mathcal{D}_\mu \mathbf{a}_1) (\mathcal{D}_\nu \mathbf{a}_2) F^Y_{\rho\sigma}$$

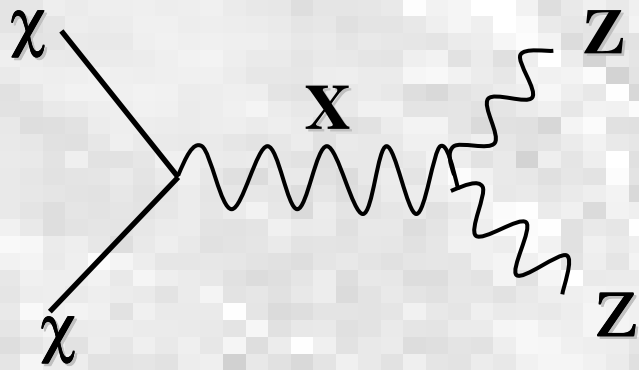
Dark matter: Annihilation channels



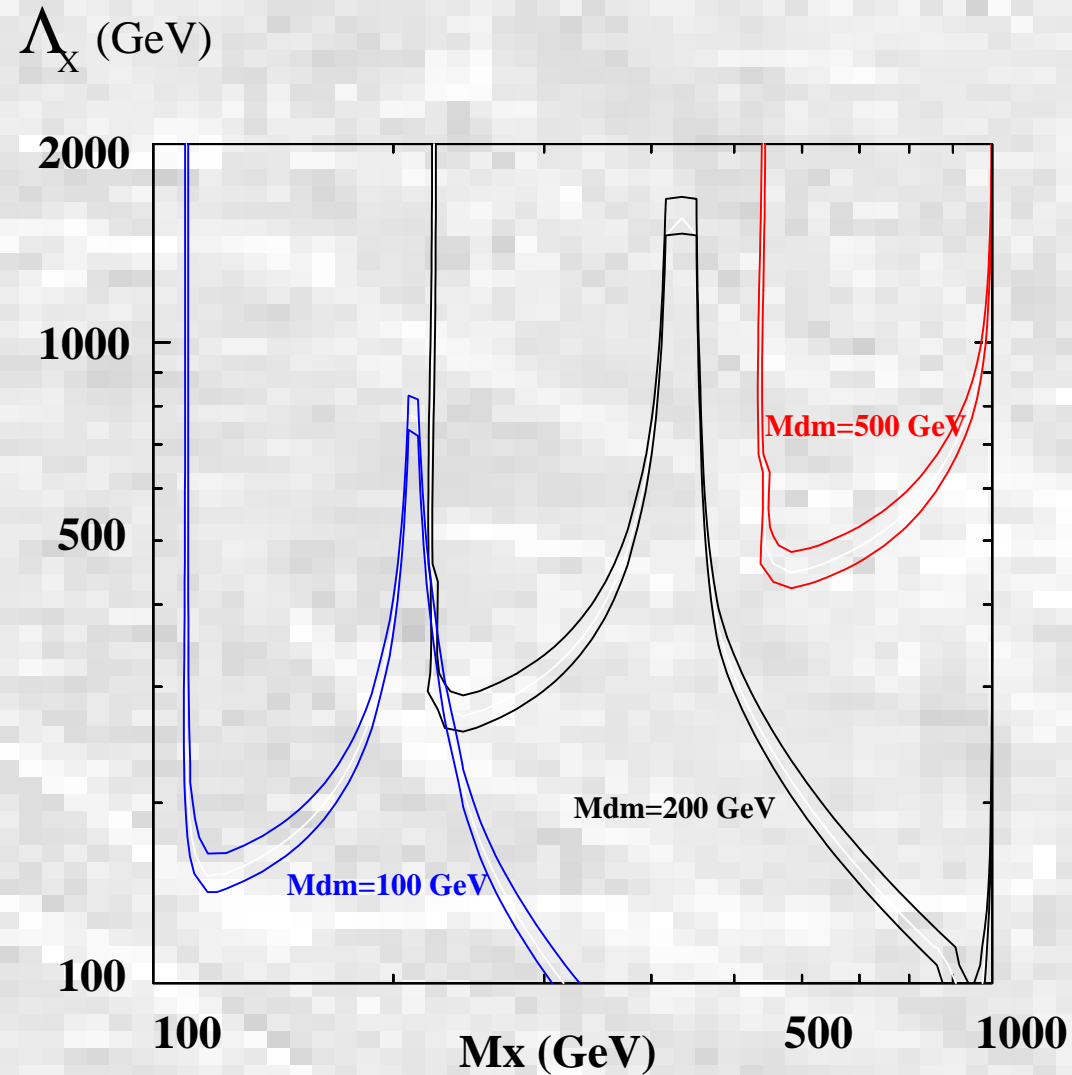
$M_X > M_\chi$ (natural)



$M_X < M_\chi$ (unnatural)

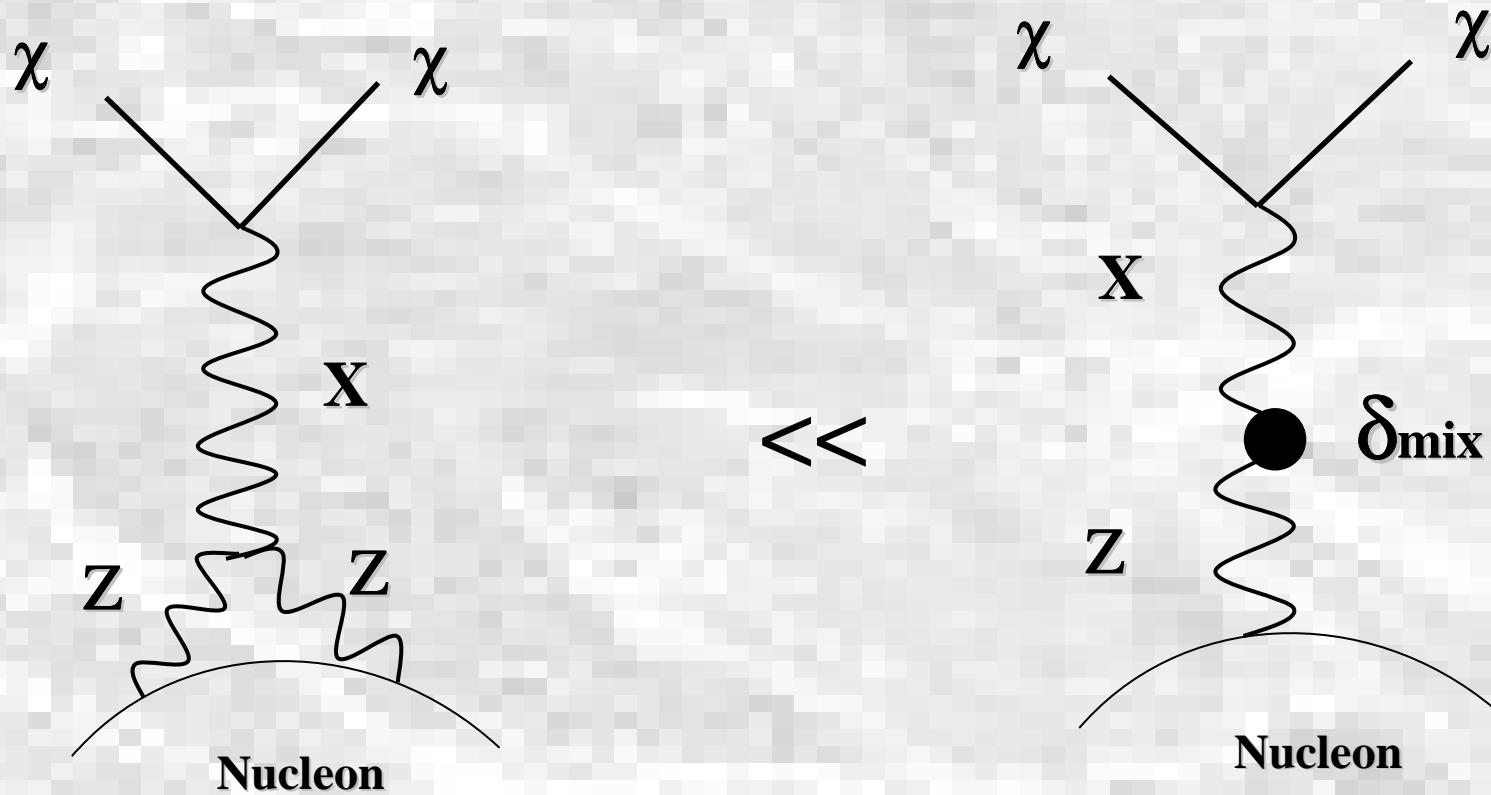


The relic density



$$\Lambda_X^2 = |\mathbf{B}_2 - \mathbf{B}_1| = g_h * g^2 / (8\pi^2) * \text{Tr}[\mathbf{X}' \mathbf{X}^2 / M_{heavy}^2] * \text{Integral}$$

Direct detection

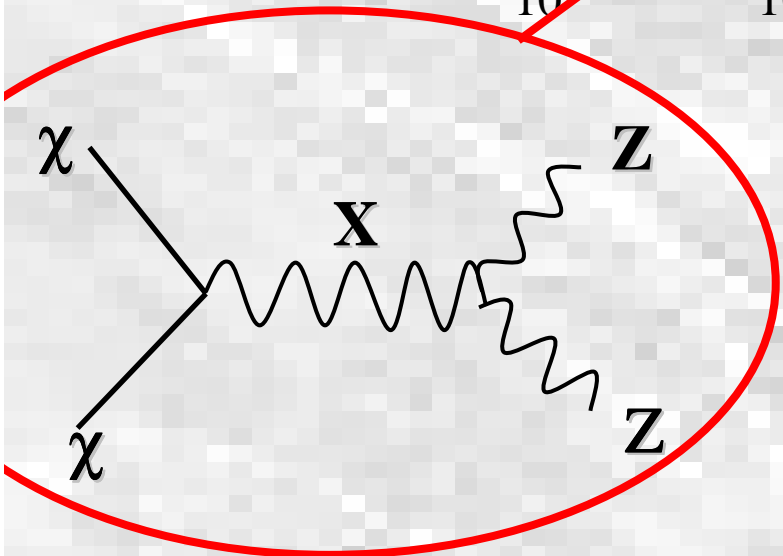
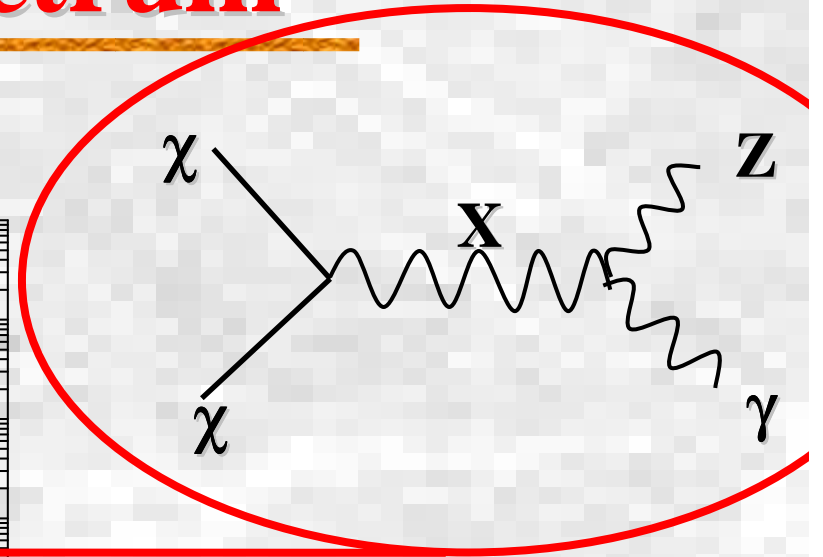
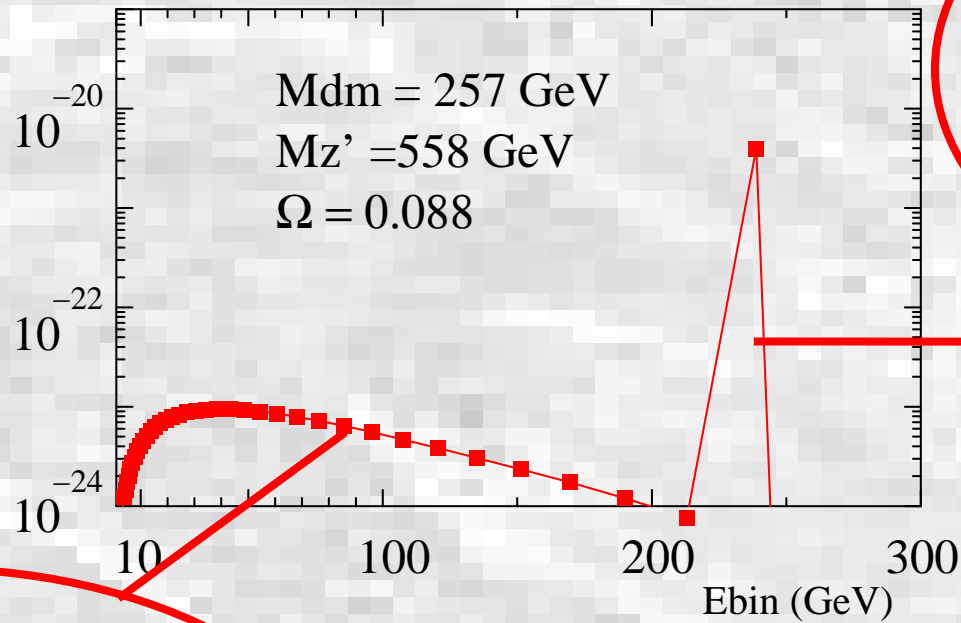


$$\mathcal{L}_{\text{mix}} = \delta_{\text{mix}} \mathbf{F}_{\mu\nu} \mathbf{F}'_{\mu\nu}$$

Millicharged DM,
XENON $\Rightarrow \delta_{\text{mix}} < 0.01$

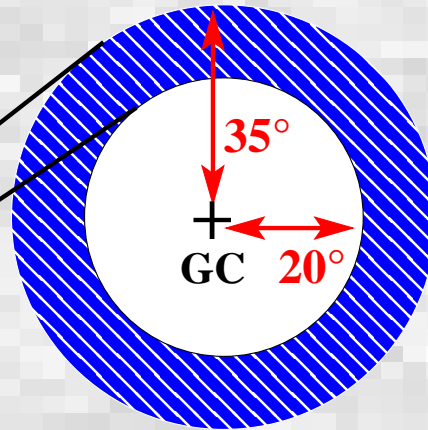
Indirect detection: examples of spectrum

$E^2 * dN/dE$



$$E_{\gamma} = M_{\chi} [1 - (M_Z/2M_{\chi})^2]$$

Indirect detection astro-parameters

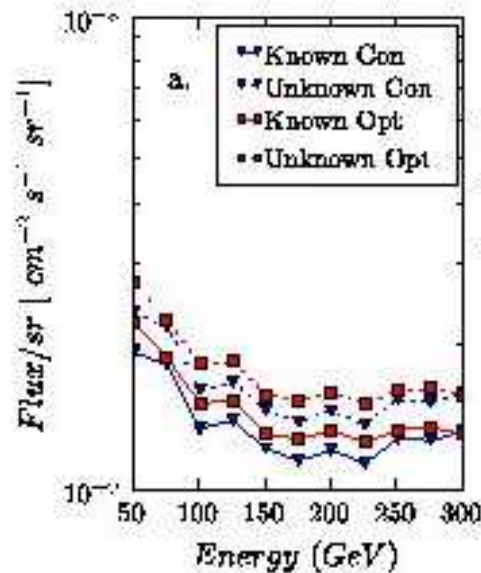


Signal to noise ratio
12 time greater than
GC

Galactic Centered Annulus
(Stoehr et al 2003, GLAST col. 2008)

Independant of the Galactic profile

$$J \Delta\Omega \sim 10$$



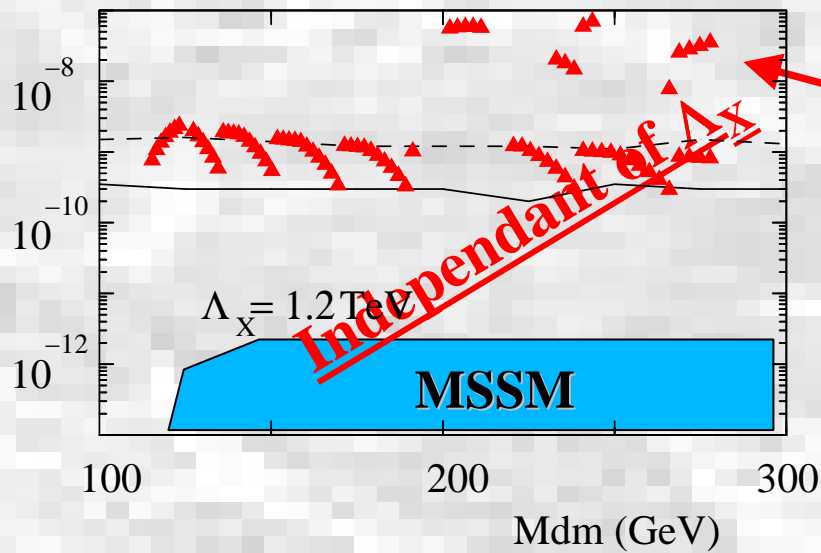
Galprop conventional model
for the background

5 years of data, signals at 5σ and
95% CL

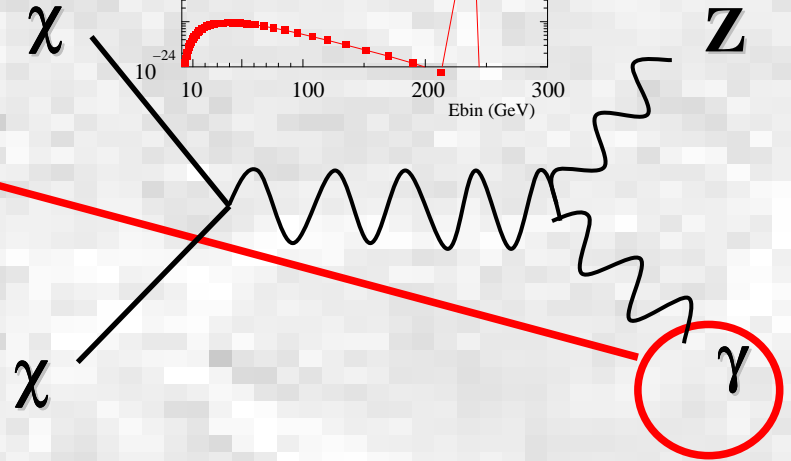
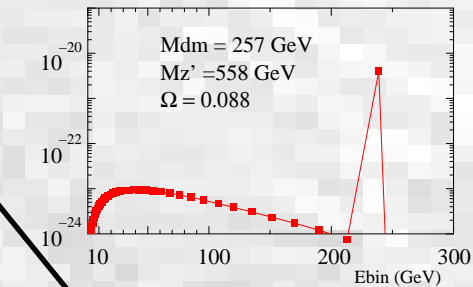
[FERMI estimates, Morselli et al. 08]

Observability

Flux (/sr/cm²/s)

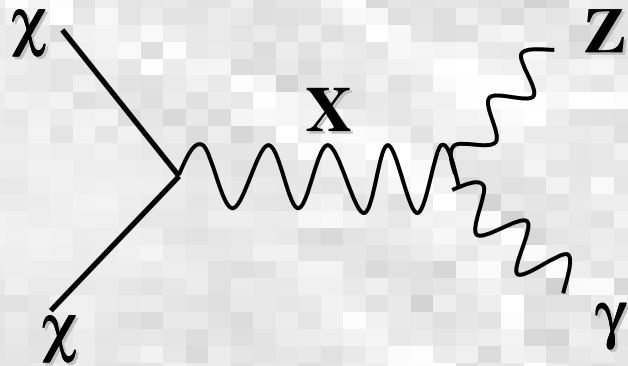
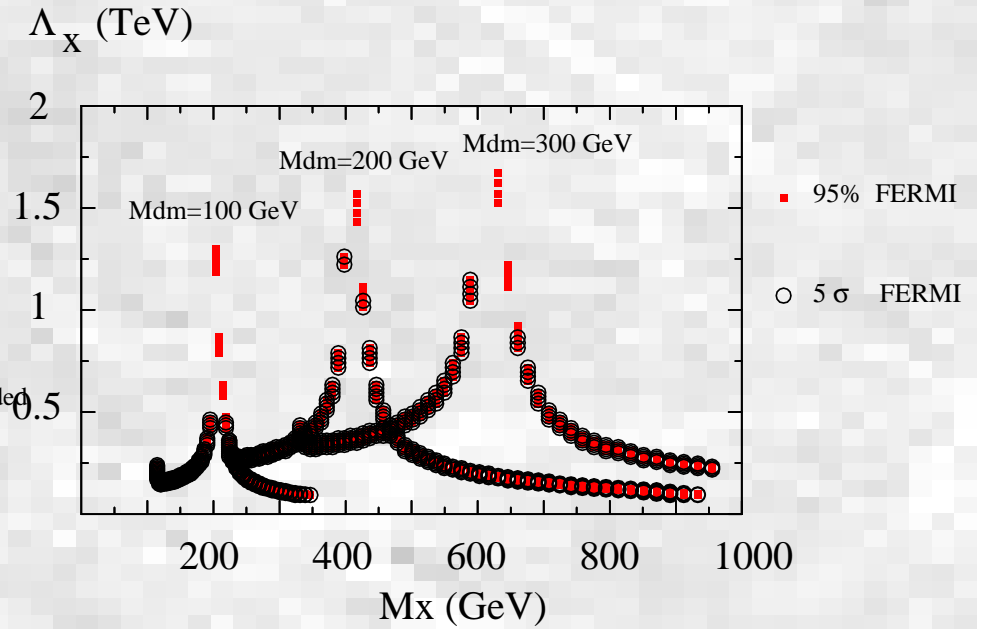
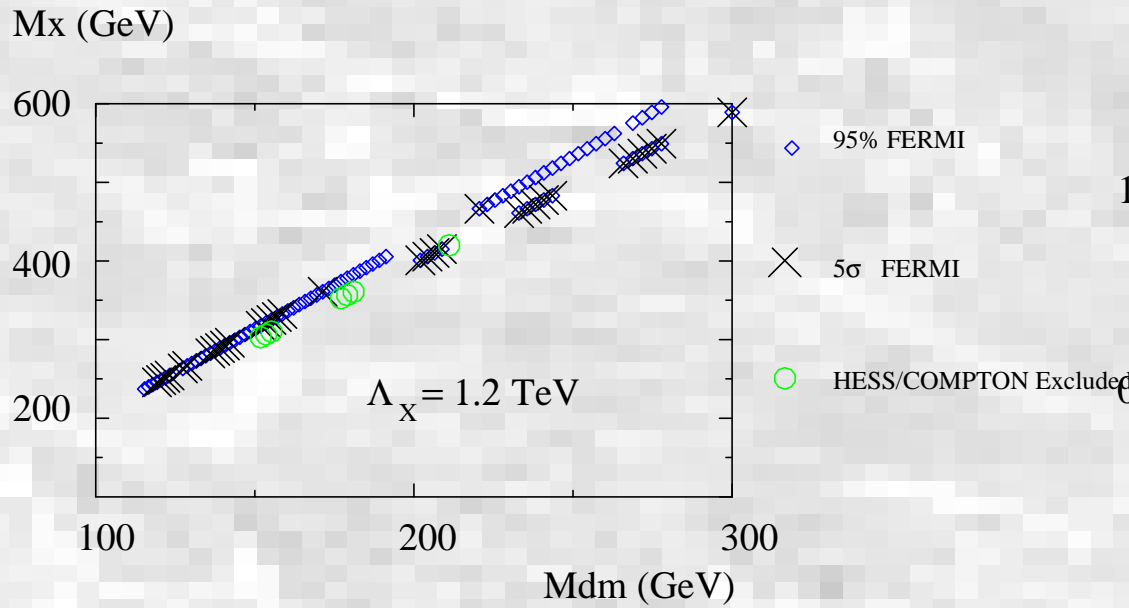


$E^2 \cdot dN/dE$



No excess with current constraints (EGRET, HESS.. [Jacques, Bell 08])

Consequences on M_x and M_{chi}



Comparison with other models

| Masse | Direct detection | Indirect detection | LHC |
|-------------------------|-------------------------|---------------------------|------------|
| SUSY/KK | Yes | No line | Yes |
| Chiral Square | Yes | 3 lines | ? |
| Inert HiggsModel | Yes | 2 lines | Yes |
| Milli-charged | Yes | No line | Yes |
| (In)visible X | No | 1 line | No |

Conclusions

An (In)visible Z' can be quite visible

Indirect detection would be THE ONLY WAY to observe it

1 γ ray line is a smoking gun signal distinguishing it clearly from other constructions

Possibility to test up to 1TeV BSM scale