

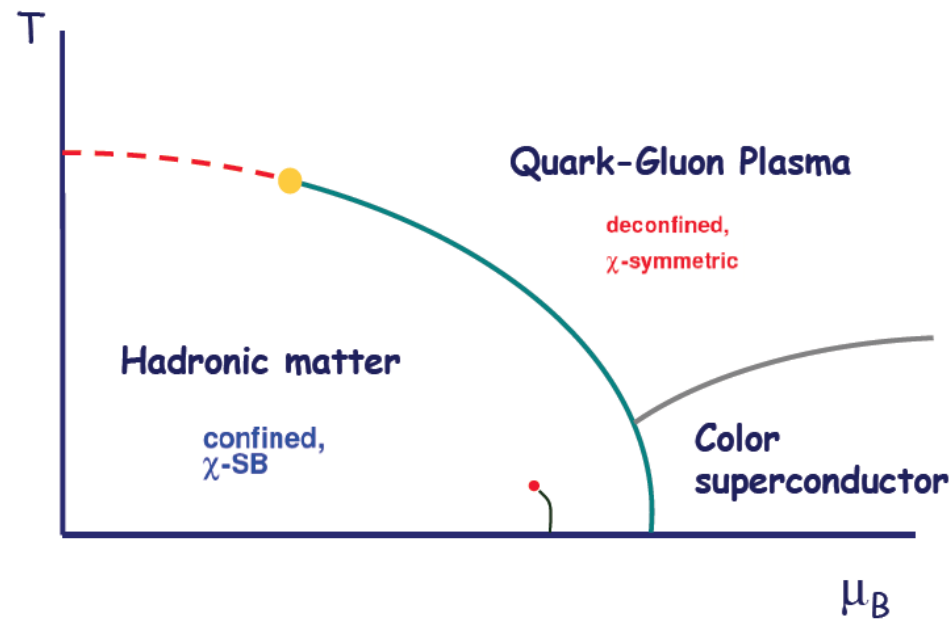
# The strongly coupled quark-gluon plasma and heavy ion collisions



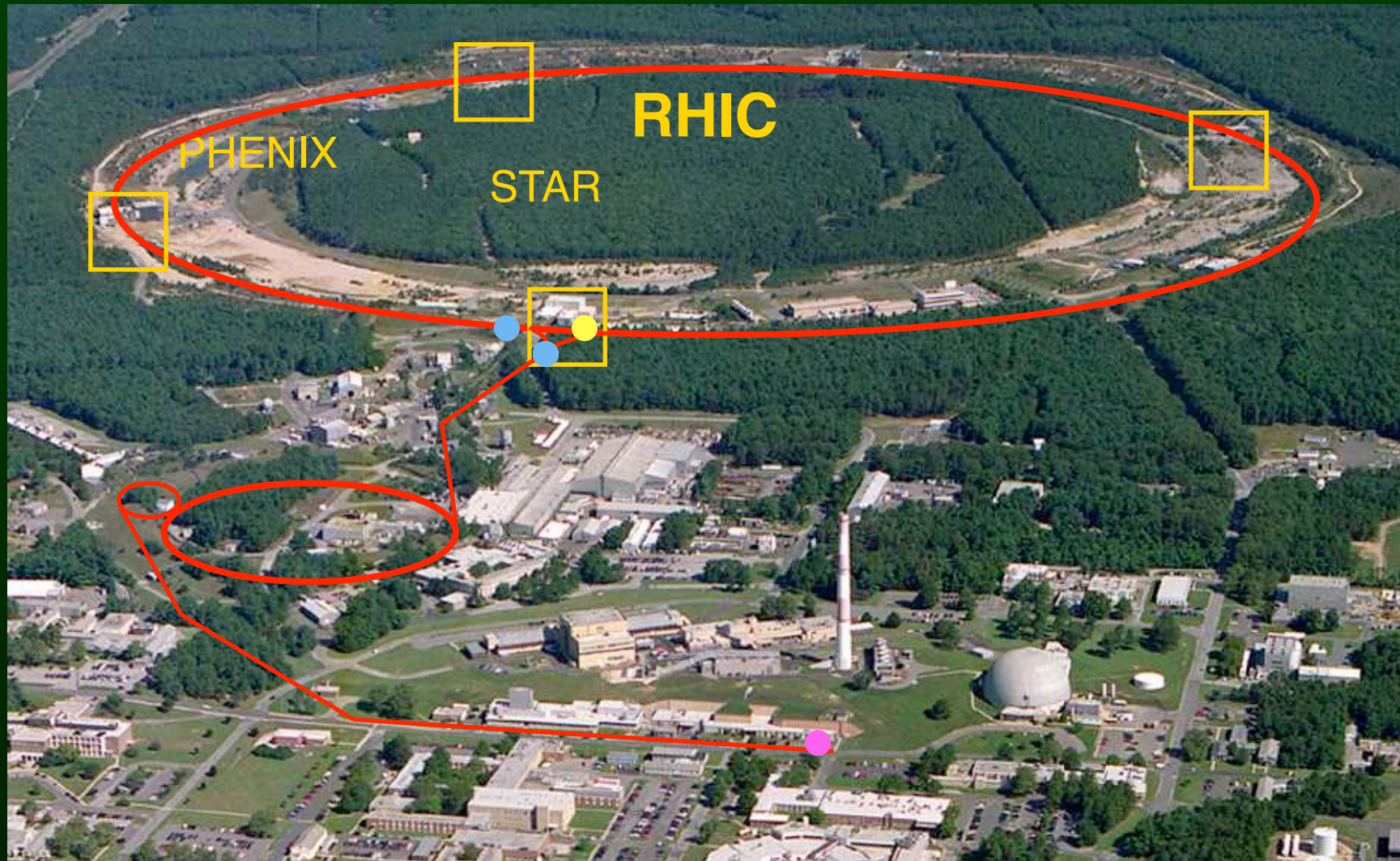
Jean-Paul Blaizot, IPHT- Saclay

# Early concept - Asymptotic freedom

$$\alpha_s = \frac{g^2}{4\pi} \approx \frac{2\pi}{b_0 \ln(\mu / \Lambda_{QCD})}$$



# Relativistic Heavy Ion Collider RHIC @ BNL



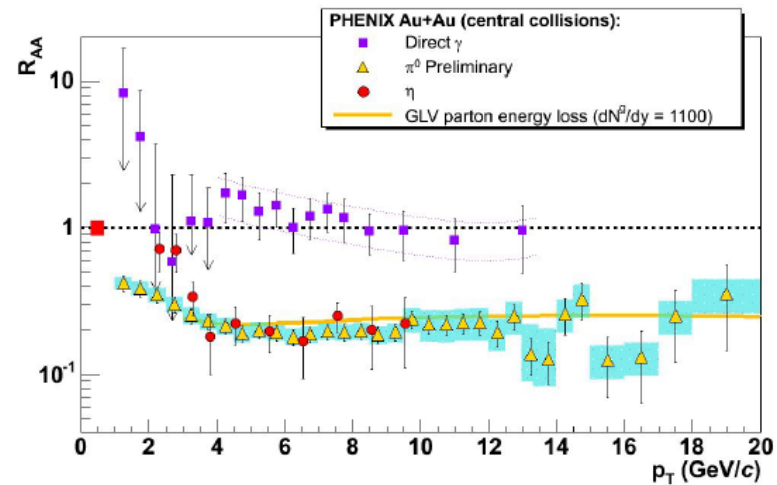
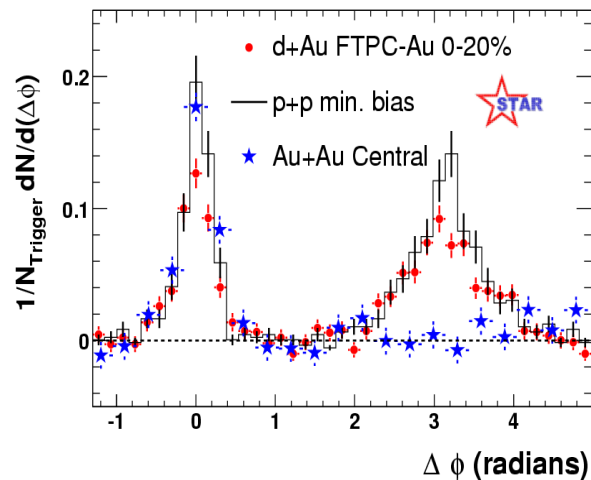
3 km ring

animation by Mike Lisa

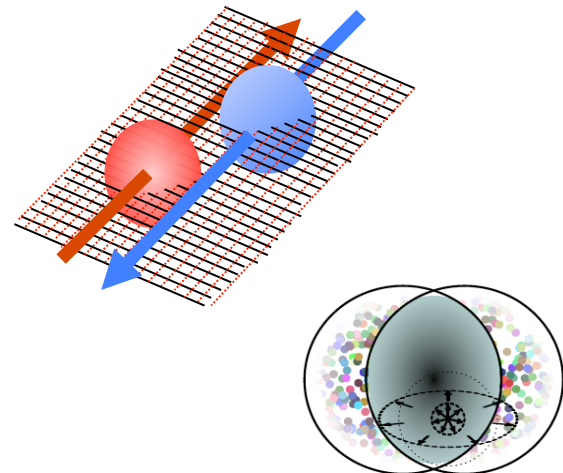
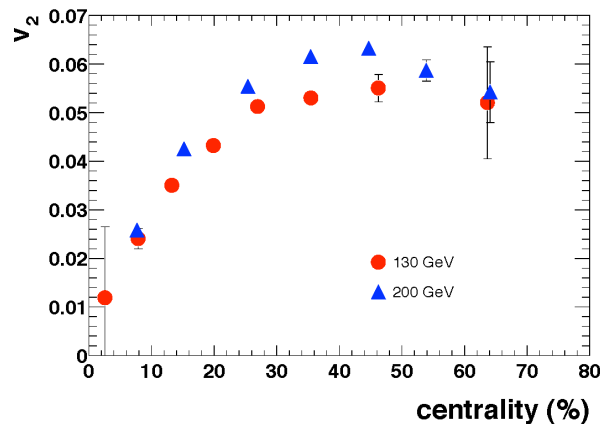
major international enterprise: thousands of scientists and engineers

# RHIC perfect liquid

Strong opacity of matter to propagation of hard partons



Collective (elliptic) flow of matter



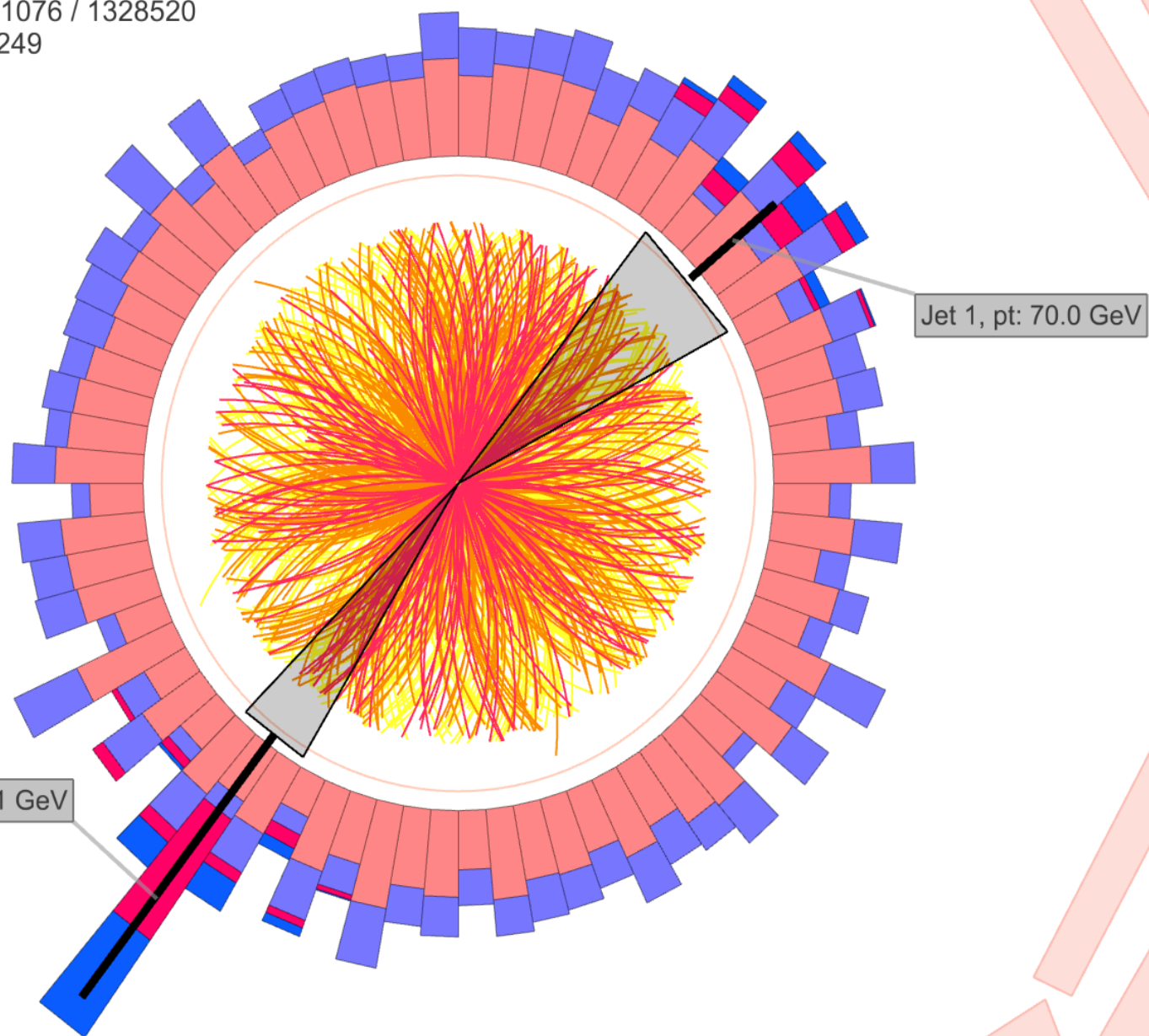


# News from the Large Hadron Collider



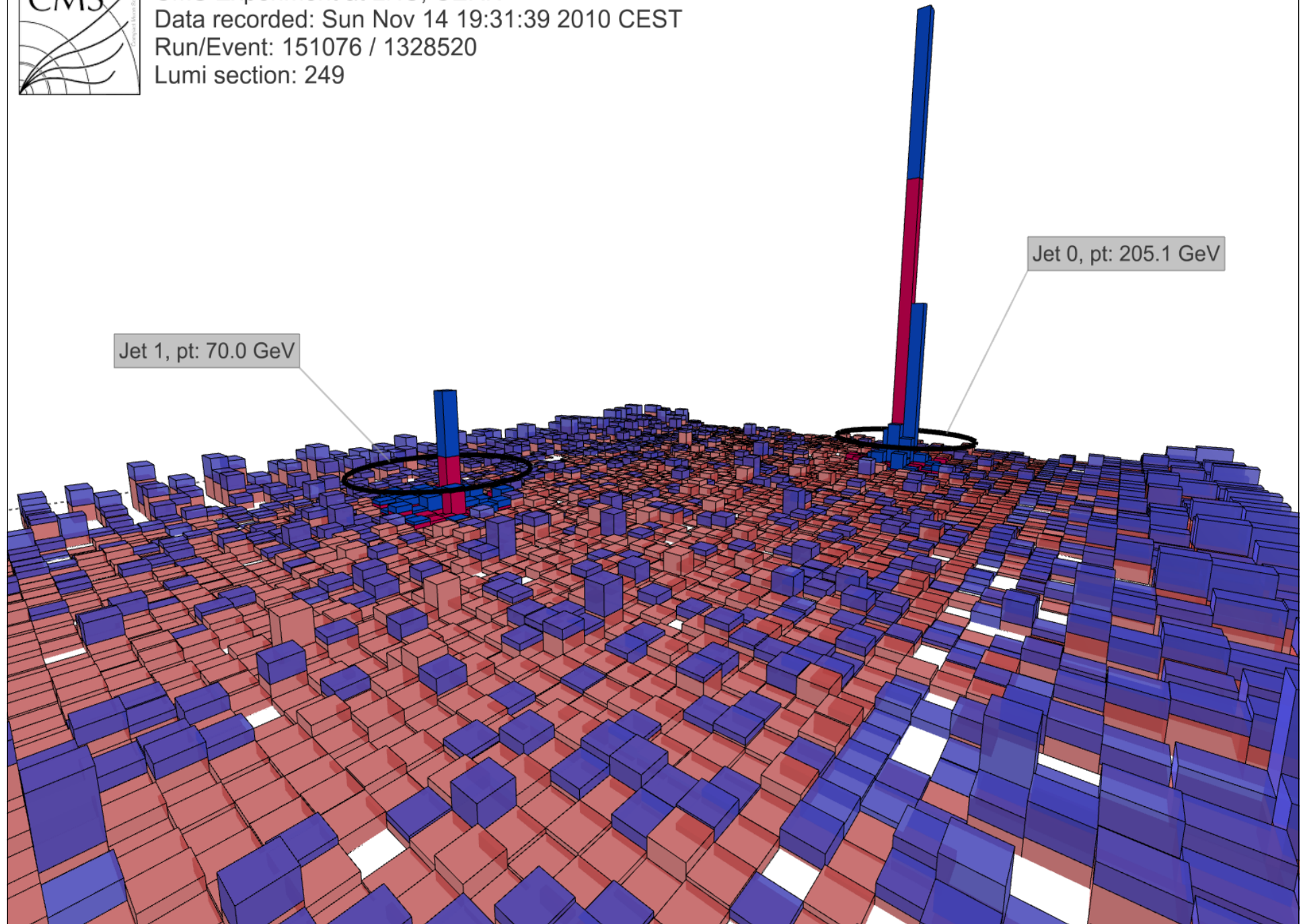


CMS Experiment at LHC, CERN  
Data recorded: Sun Nov 14 19:31:39 2010 CEST  
Run/Event: 151076 / 1328520  
Lumi section: 249

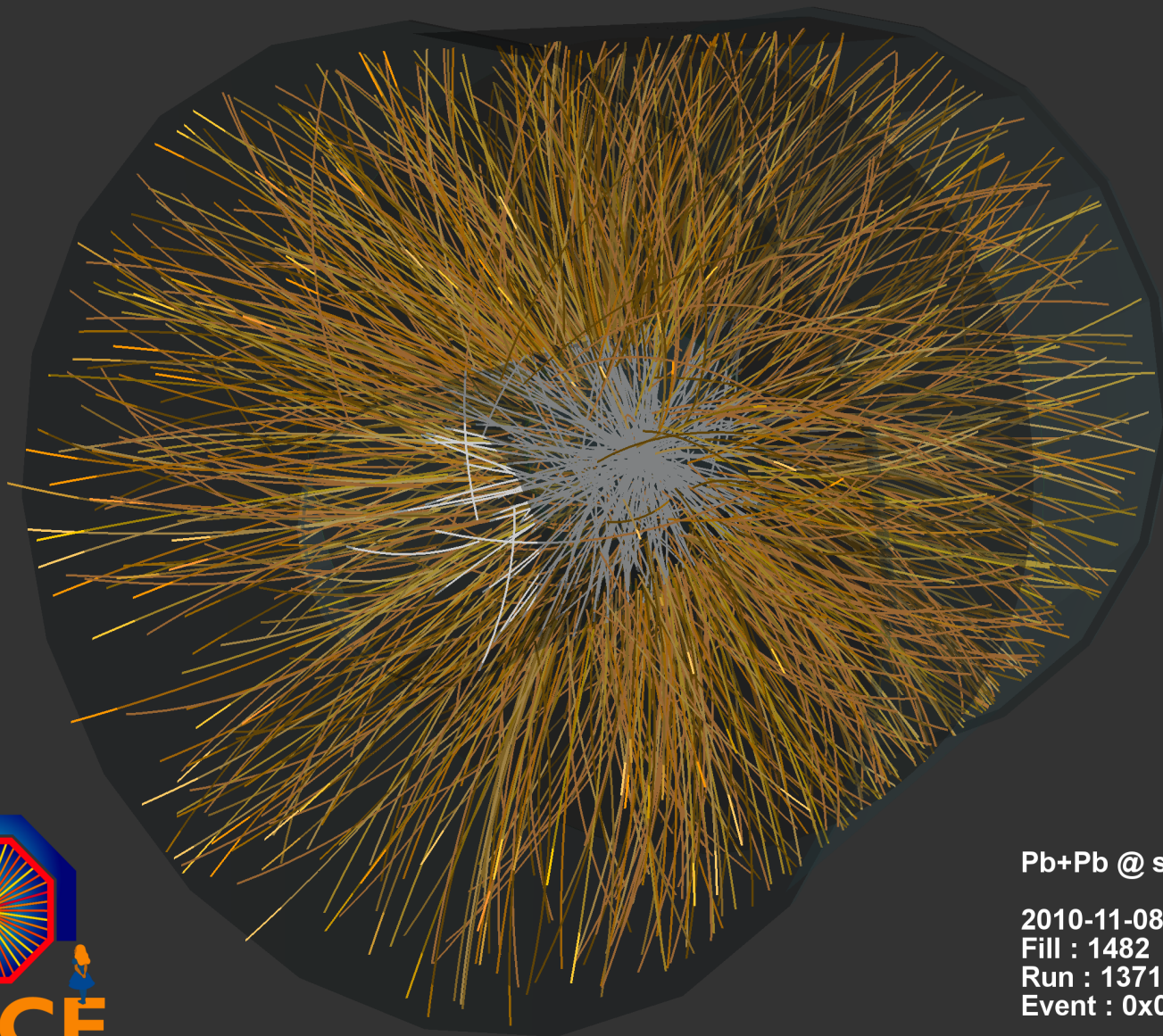




CMS Experiment at LHC, CERN  
Data recorded: Sun Nov 14 19:31:39 2010 CEST  
Run/Event: 151076 / 1328520  
Lumi section: 249







Pb+Pb @  $\sqrt{s} = 2.76$  ATeV

2010-11-08 11:33:20

Fill : 1482

Run : 137124

Event : 0x00000007FE50693

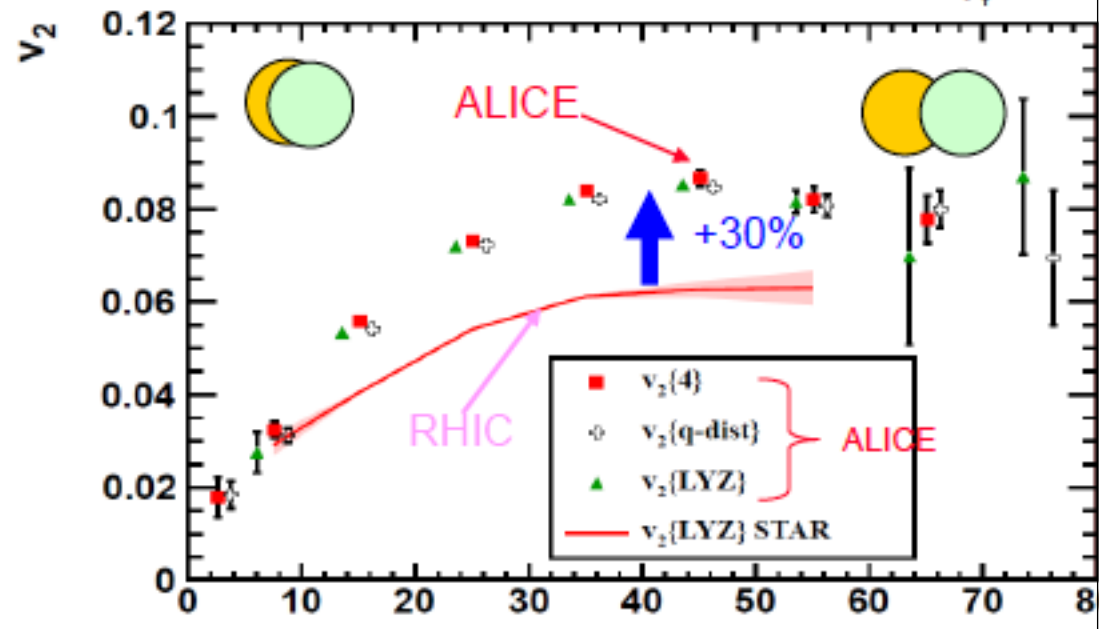
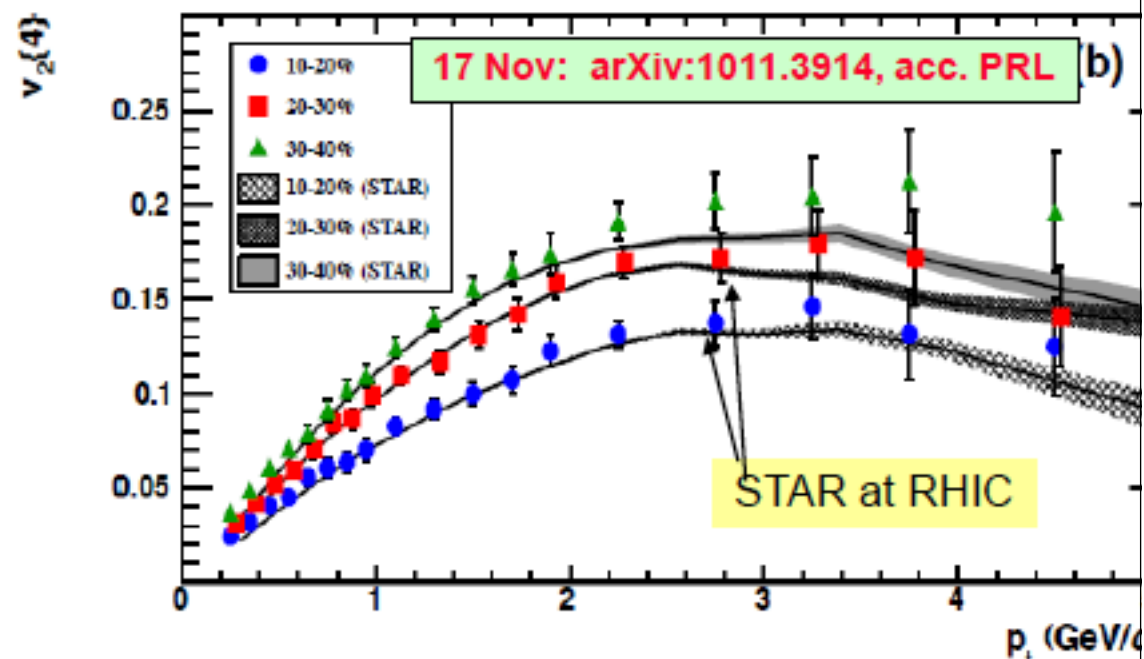




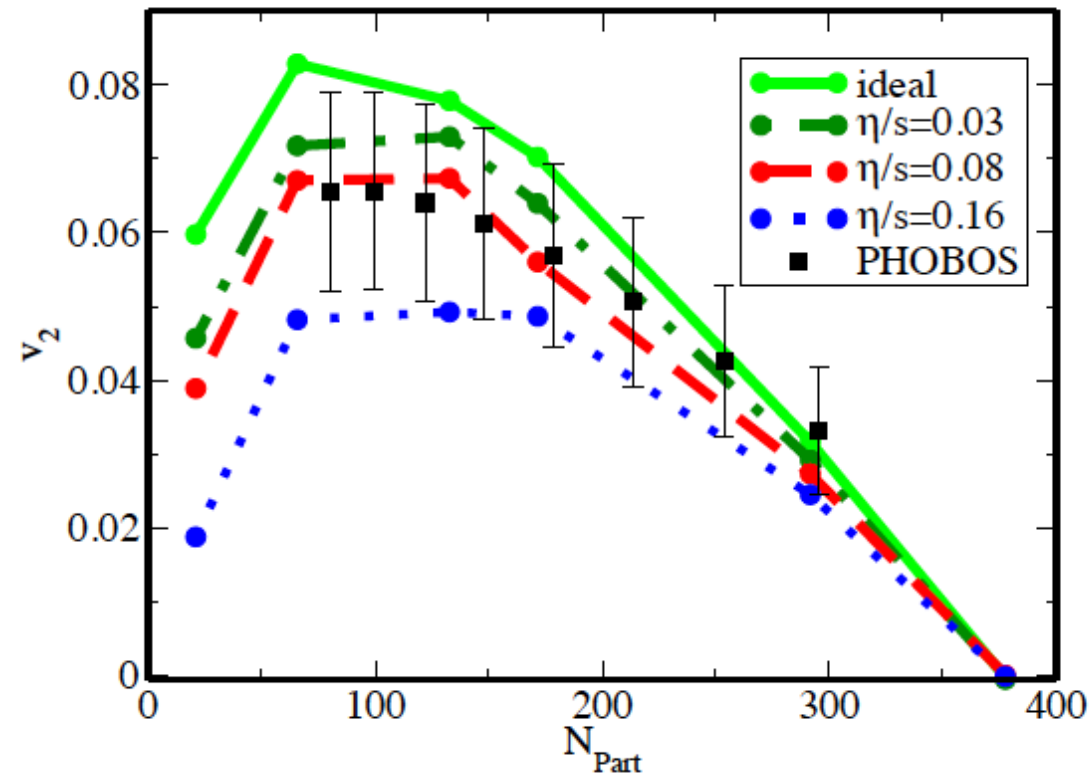
# First Elliptic Flow Measurement at LHC



- $v_2$  as function of  $p_t$ 
  - ⇒ practically no change with energy !
  - ✱ extends towards larger centrality/higher  $p_t$  ?
- $v_2$  integrated over  $p_t$ 
  - ⇒ 30% increase from RHIC
  - ⇒  $\langle p_t \rangle$  increases with  $\sqrt{s}$
  - ✱ pQCD powerlaw tail ?
  - ⇒ Hydro predicts increased 'radial flow'
  - ✱ very characteristic  $p_t$  and mass dependence; to be confirmed !



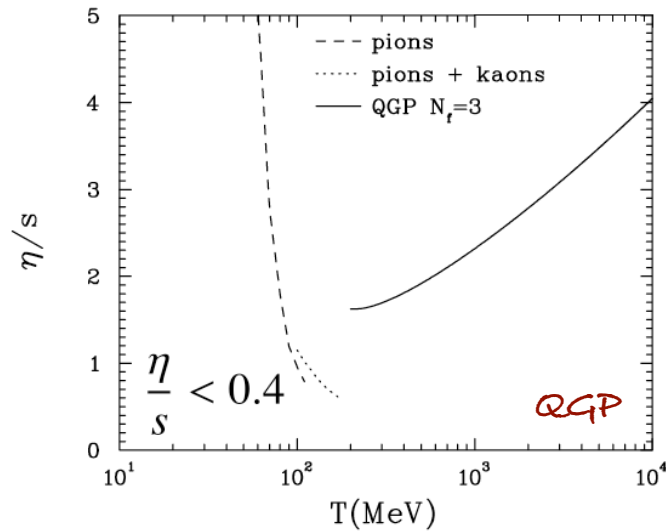
# The low viscosity of the quark-gluon plasma



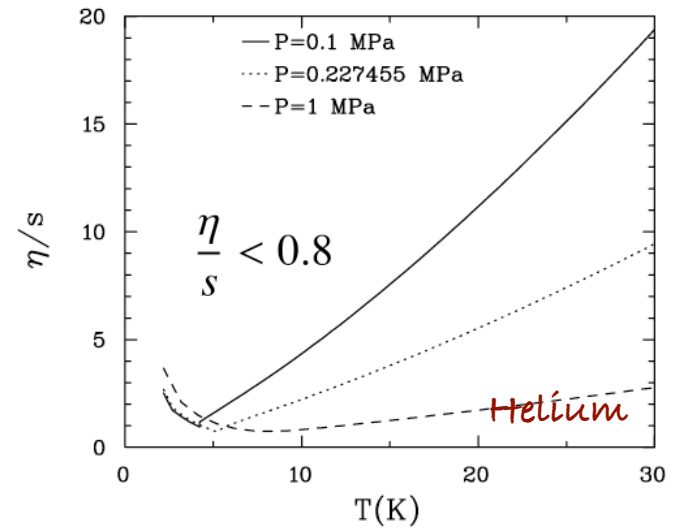
$$\frac{\eta}{s} < 5 \times \frac{1}{4\pi}$$

(Luzum, Romatschke, 2007)

# Low viscosity, phase transition and strong coupling



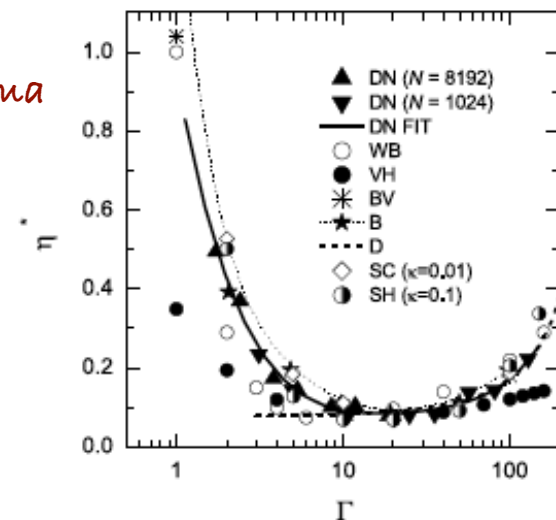
(Csernai et al, 2006)



Strongly correlated plasma

(Z. Donko et al,  
arXiv: 0710.5229)

Also, cold fermionic gas  
at unitarity  $\frac{\eta}{s} \sim 0.5$



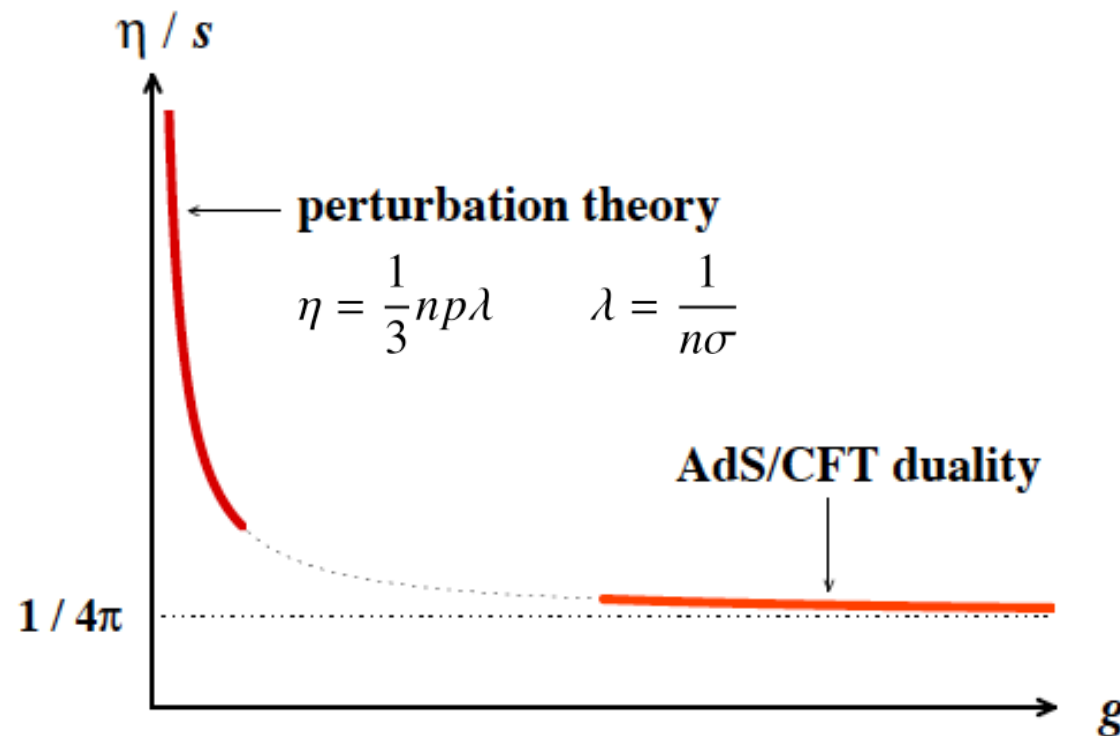


# AdS/CFT Duality

A 'natural' explanation for the small viscosity

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

(Polícastro et al, 2001)



A puzzling situation

Where is the apparent  
strongly coupled character  
of the QGP coming from?

Is initial concept wrong ?

No...

QCD asymptotic freedom works !



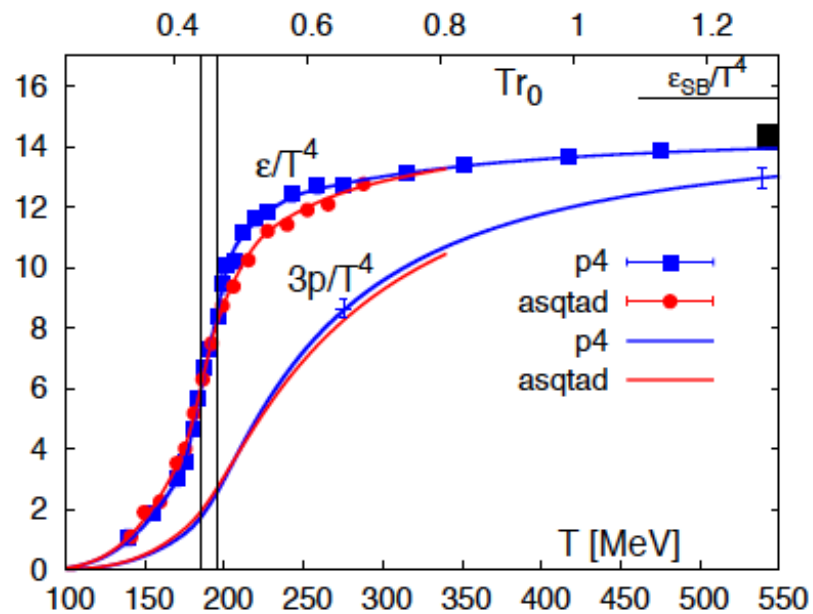
# Asymptotic freedom

$$\alpha_s = \frac{g^2}{4\pi} \approx \frac{2\pi}{b_0 \ln(\mu / \Lambda_{QCD})} \quad (\mu \approx 2\pi T)$$

Matter is « simple » at high temperature:

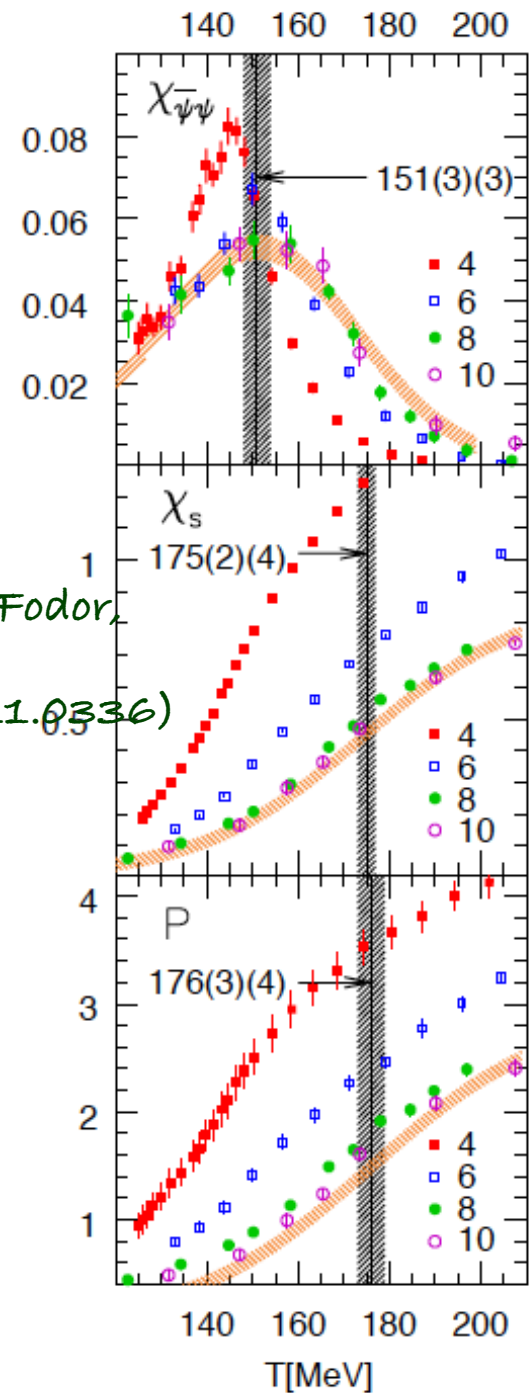
- an ideal gas of quarks and gluons
- the dominant effect of interactions is to turn (massless) quarks and gluons into weakly interacting (massive) quasiparticles.

# Crossover from hadrons to quarks and gluons

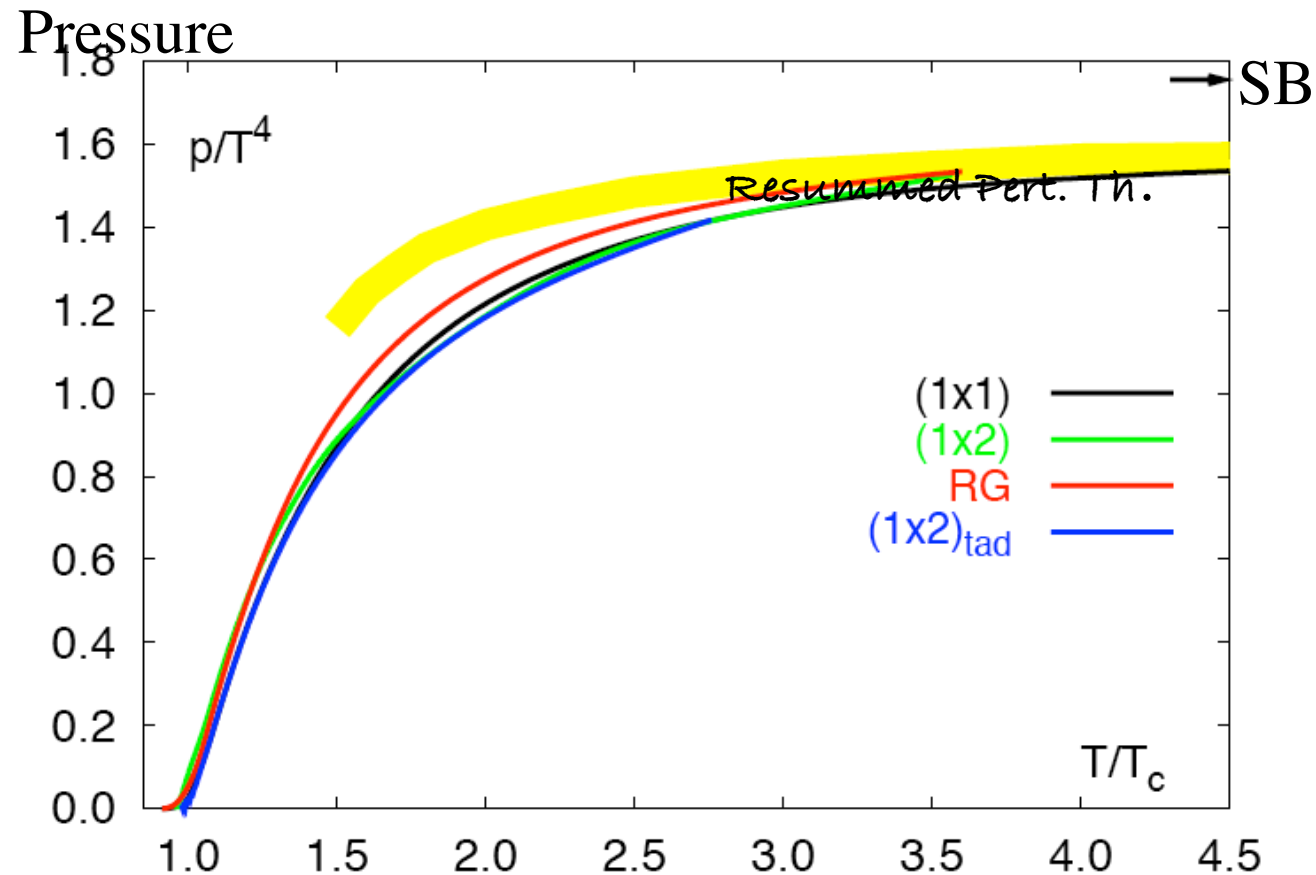


(from M. Bazavov et al, arXiv:0903.4379)

(from Z. Fodor,  
arXiv:0711.0336)



At  $T > 3T_c$  Resummed Pert. Theory  
accounts for lattice results

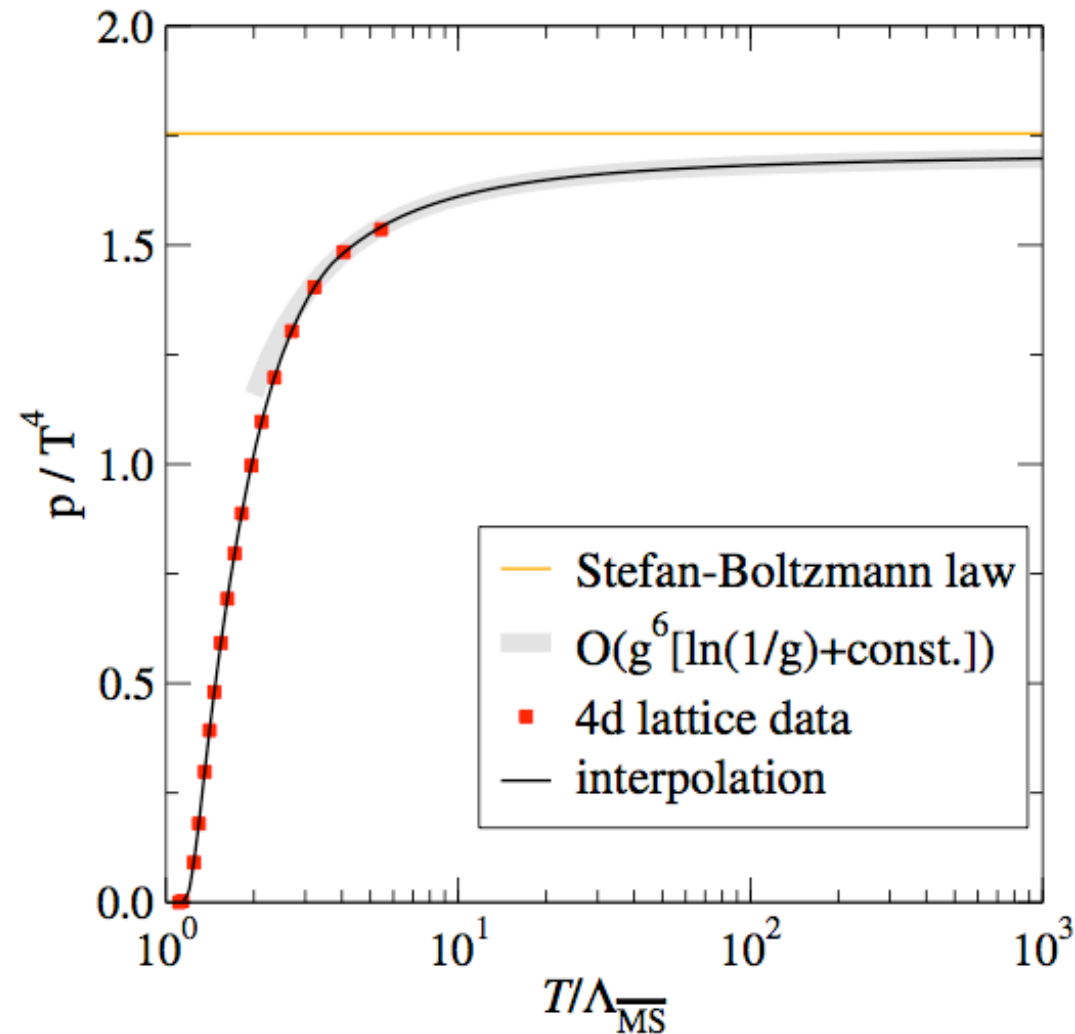


(SU(3) lattice gauge calculation from Karsch et al, hep-lat/0106019)

(resummed pert. th. from J.-P. B., E. Iancu, A. Rebhan: Nucl.Phys.A698:404-407,2002)

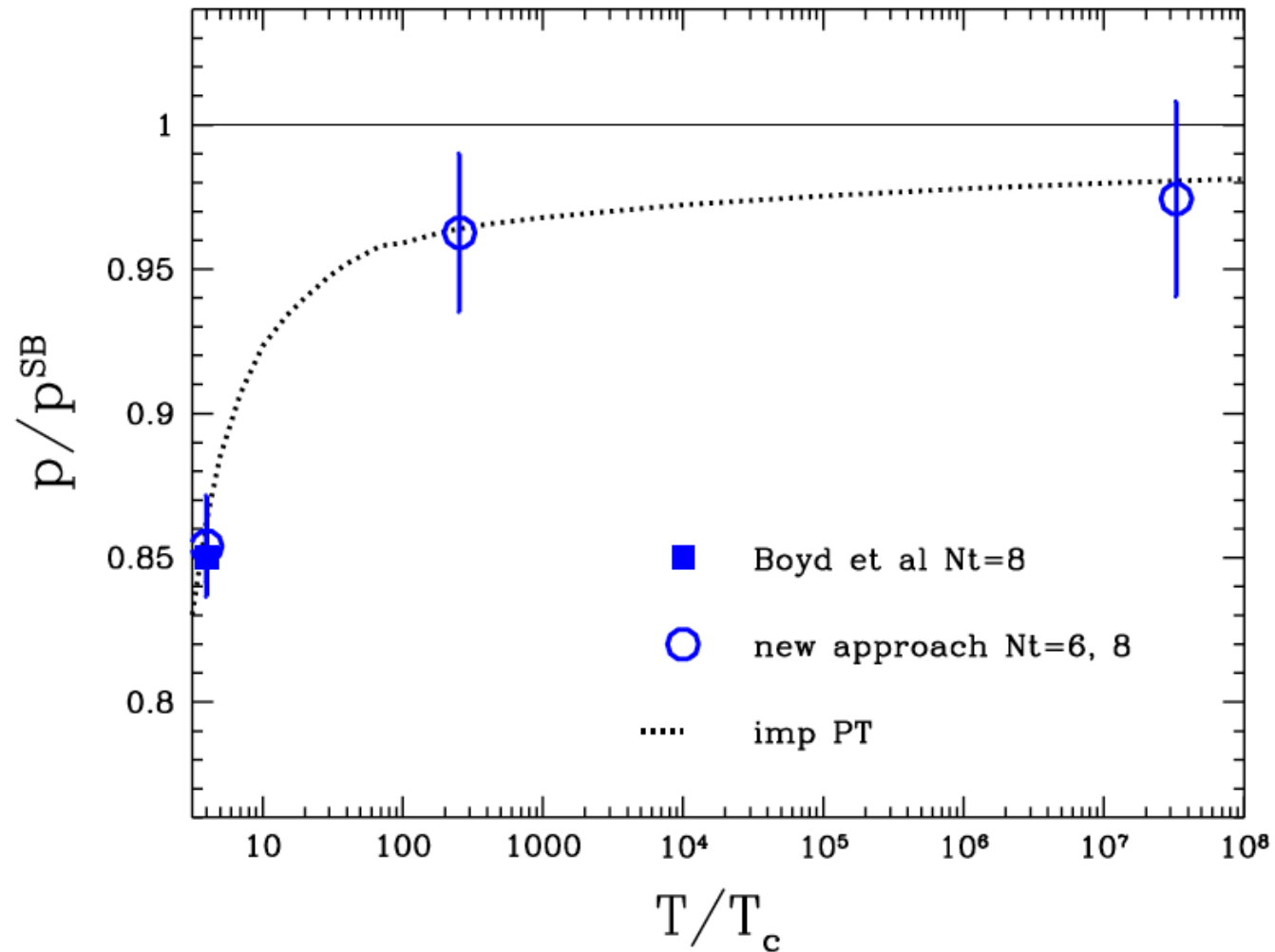


## State of the art in high order perturbative calculations



(from M. Laine, Y Schroeder, hep-ph/0603048)

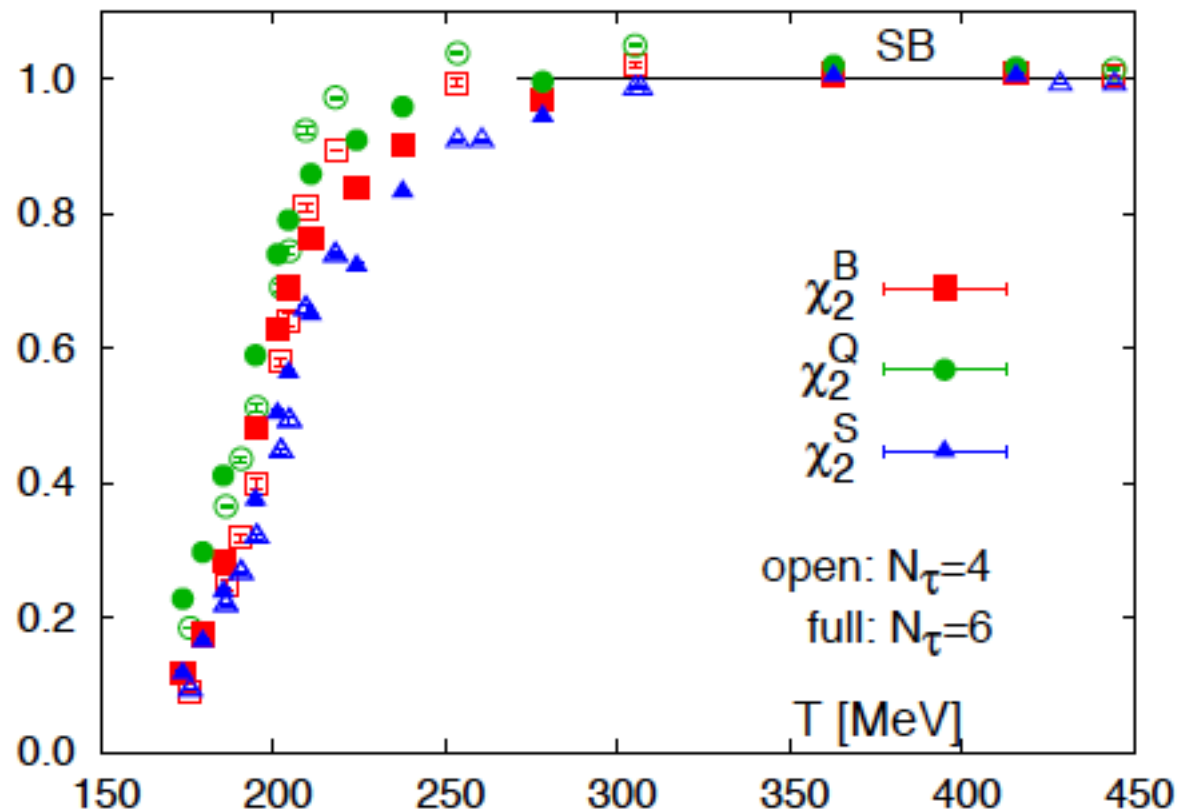
## Pressure for $SU(3)$ YM theory at (very) high temperature



(from G. Endrodi et al, arXiv: 0710.4197)

# Conserved charge susceptibilities

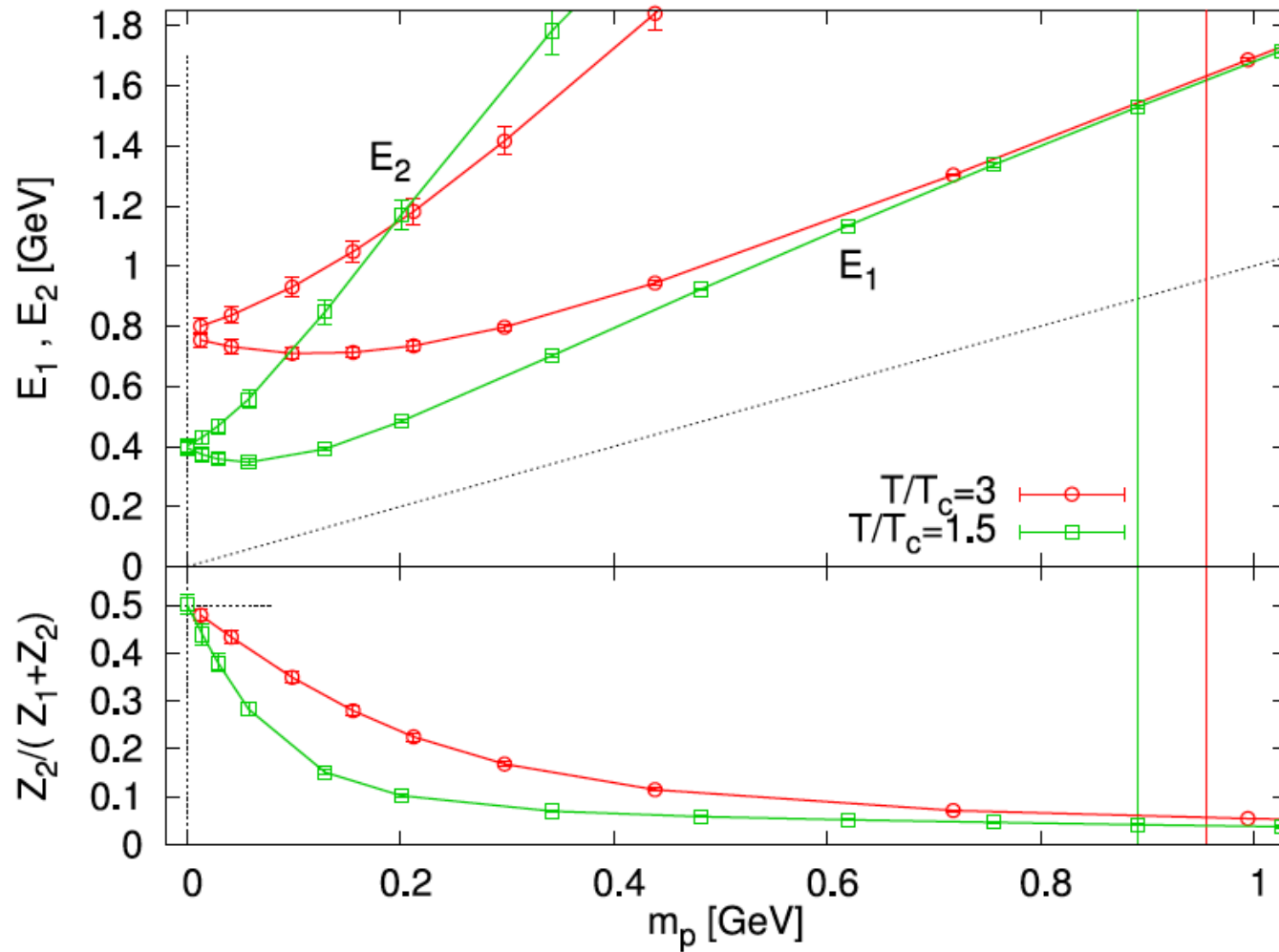
$$\chi_C \sim \langle C^2 \rangle \quad C = B, Q, S$$



(from M. Cheng et al, arXiv: 0811.1006)



## Quark quasiparticles seen on the lattice

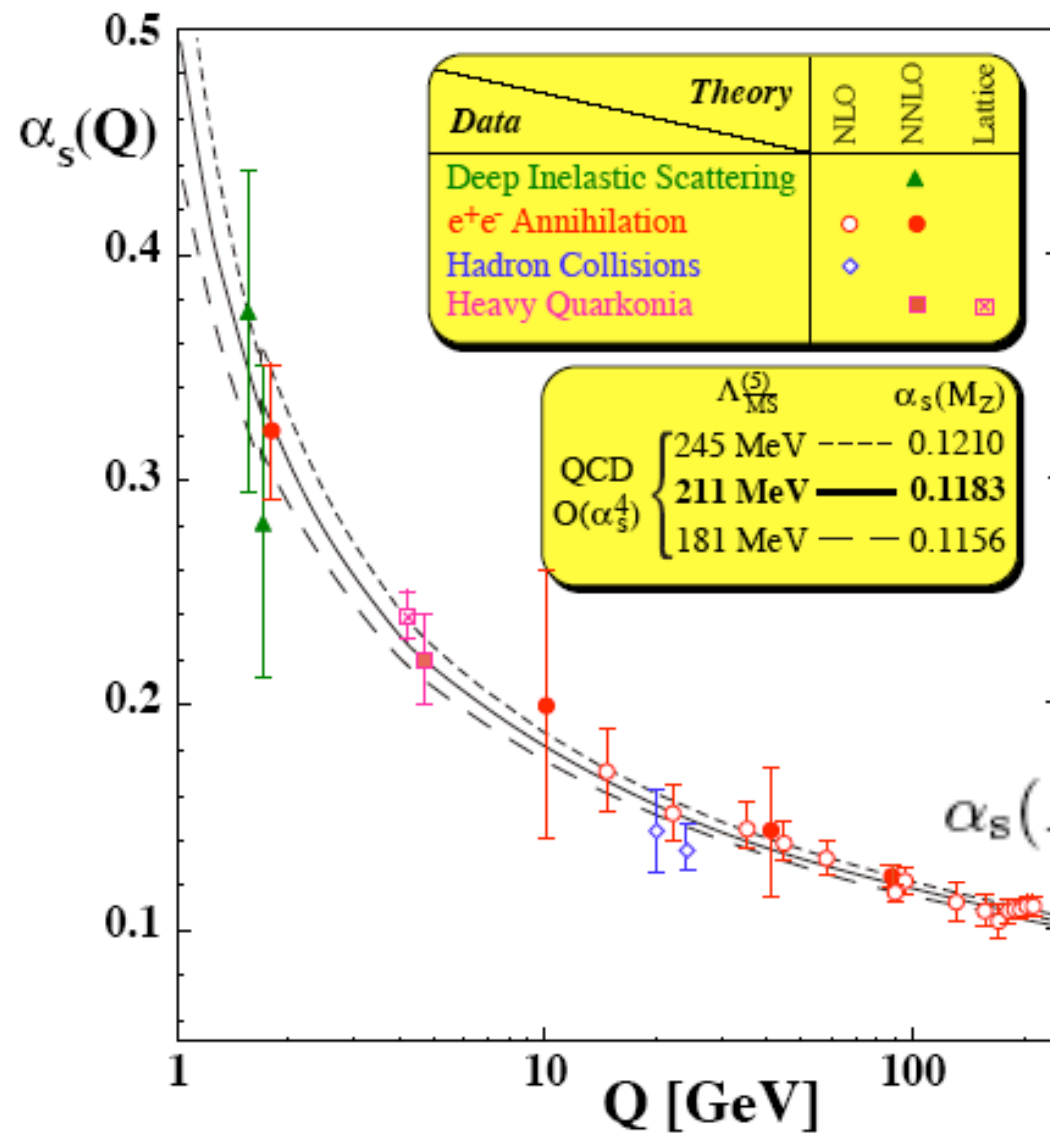


(from F. Karsch and M. Kitazawa, arXiv: 0906.3941)

Is the coupling constant large ?

Not really !

# The QCD running coupling constant

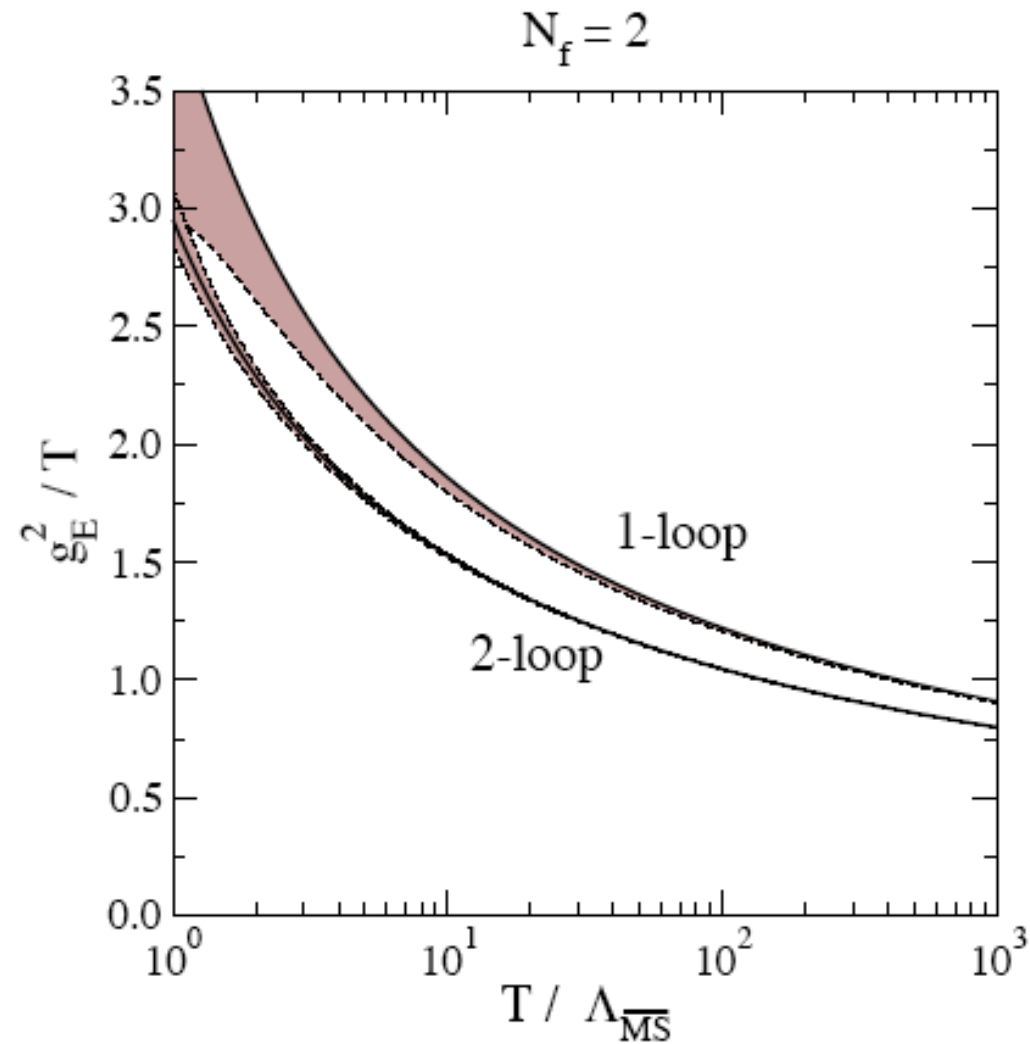


$$\alpha_s = \frac{g^2}{4\pi}$$

$$\alpha_s(M_Z) = 0.1183 \pm 0.0027$$

(S. Bethke, hep-ex/0211012)

The effective coupling is not huge, even close to  $T_c$



Strict perturbation theory breaks down

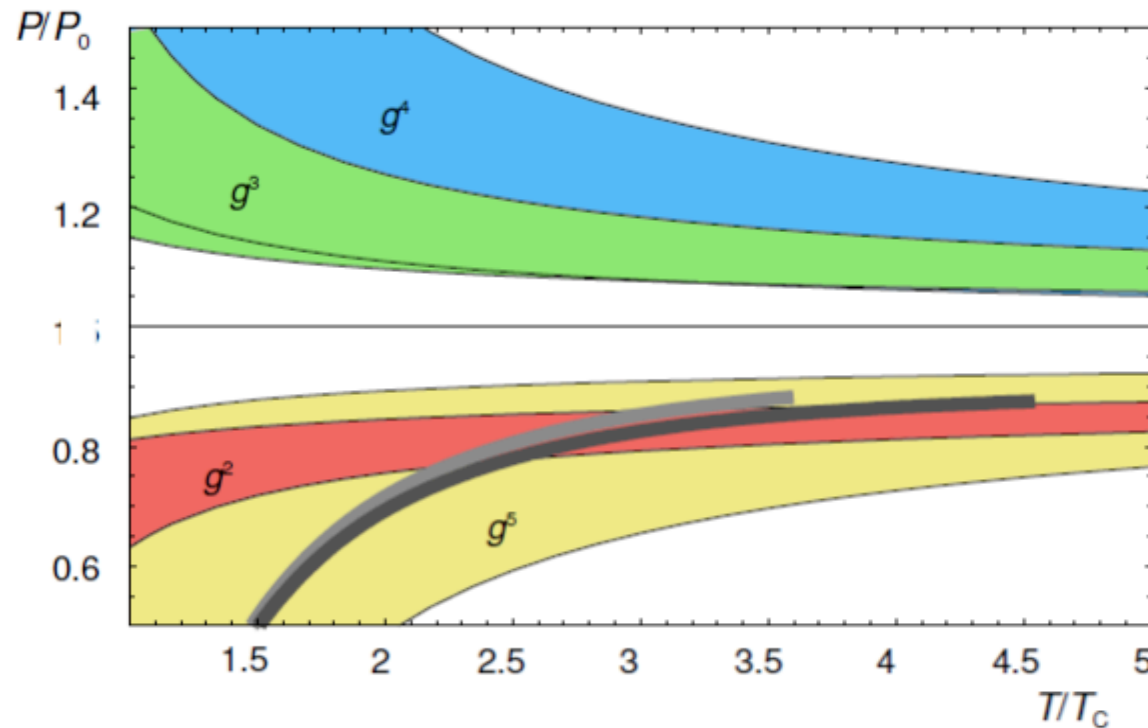
Often used as an argument in favor of strong coupling

But

- this has (almost) nothing to do with QCD
- pb can be handled with a variety of techniques (resummations, exact RG, etc)



# Perturbation theory is ill behaved at finite temperature



Perturbation theory:

$g^2$ : Shuryak; Chin (1978)

$g^3$ : Kapusta (1979)

$g^4$  In  $g$ : Toimela (1983)

$g^4$ : Arnold, Zhai (1994)

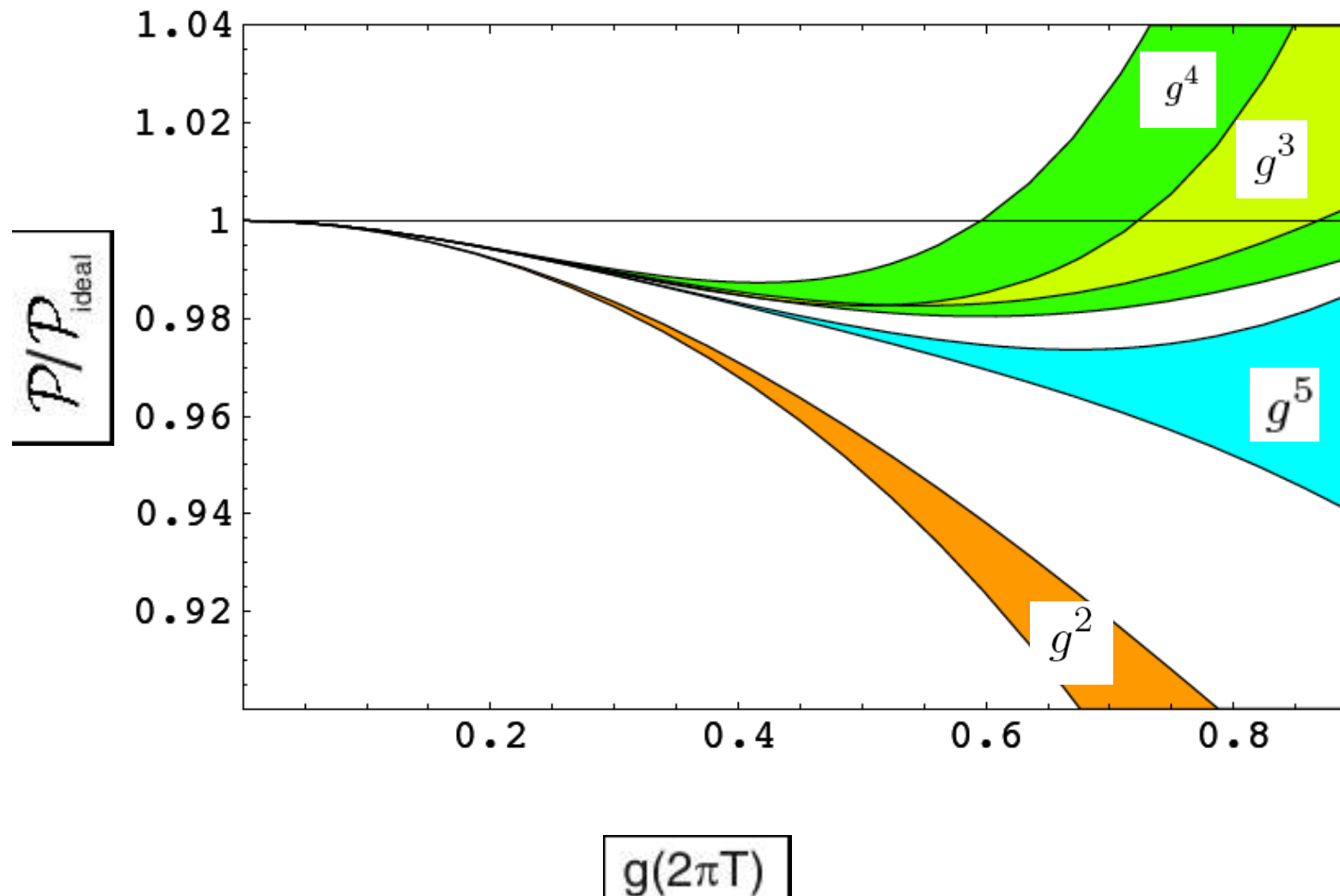
$g^5$ : Zhai, Kastening (1995),  
Braaten, Nieto (1996)

$g^6$  In  $g$ : Kajantie, Laine,  
Rummukainen, Schröder  
(2002)

$g^6$  (partly): Di Renzo, Laine,  
Miccio,  
Schröder, Torrero (2006)

Lattice data: G. Boyd et al. (1996); M. Okamoto et al. (1999).

Generic feature in (most) field theories,  
e.g in scalar field theory



## Weakly AND strongly coupled ...

Degrees of freedom with different wavelengths are differently coupled.

Expansion parameter

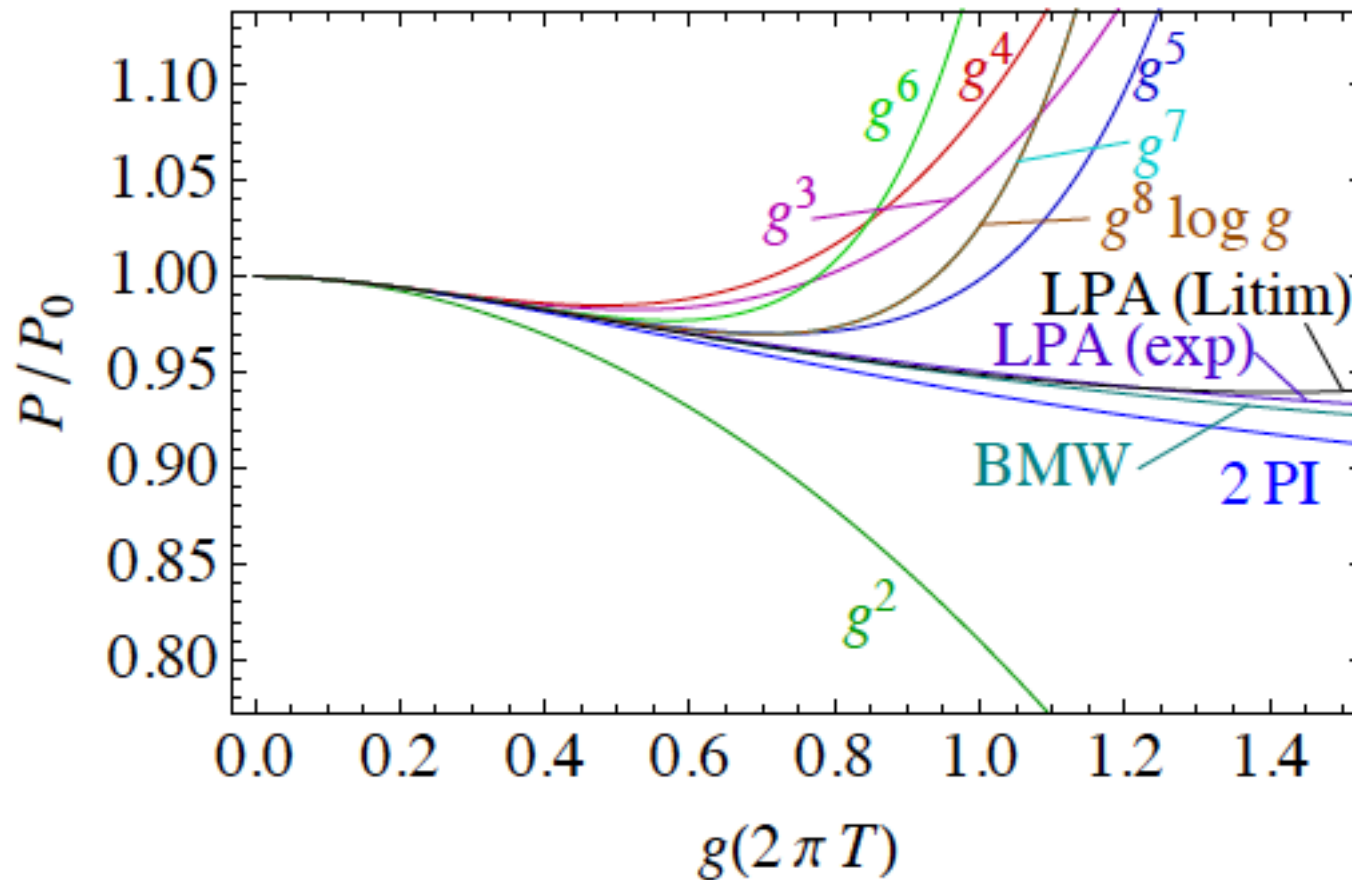
$$\gamma_K = \frac{g^2 \langle \phi^2 \rangle}{K^2} \quad \langle \phi^2 \rangle_K \sim K T \quad (K \lesssim T)$$

Dynamical scales

$$\begin{array}{ll} K \sim T & \gamma_K \sim g^2 \\ K \sim gT & \gamma_K \sim g \\ K \sim g^2 T & \gamma_K \sim 1 \end{array}$$

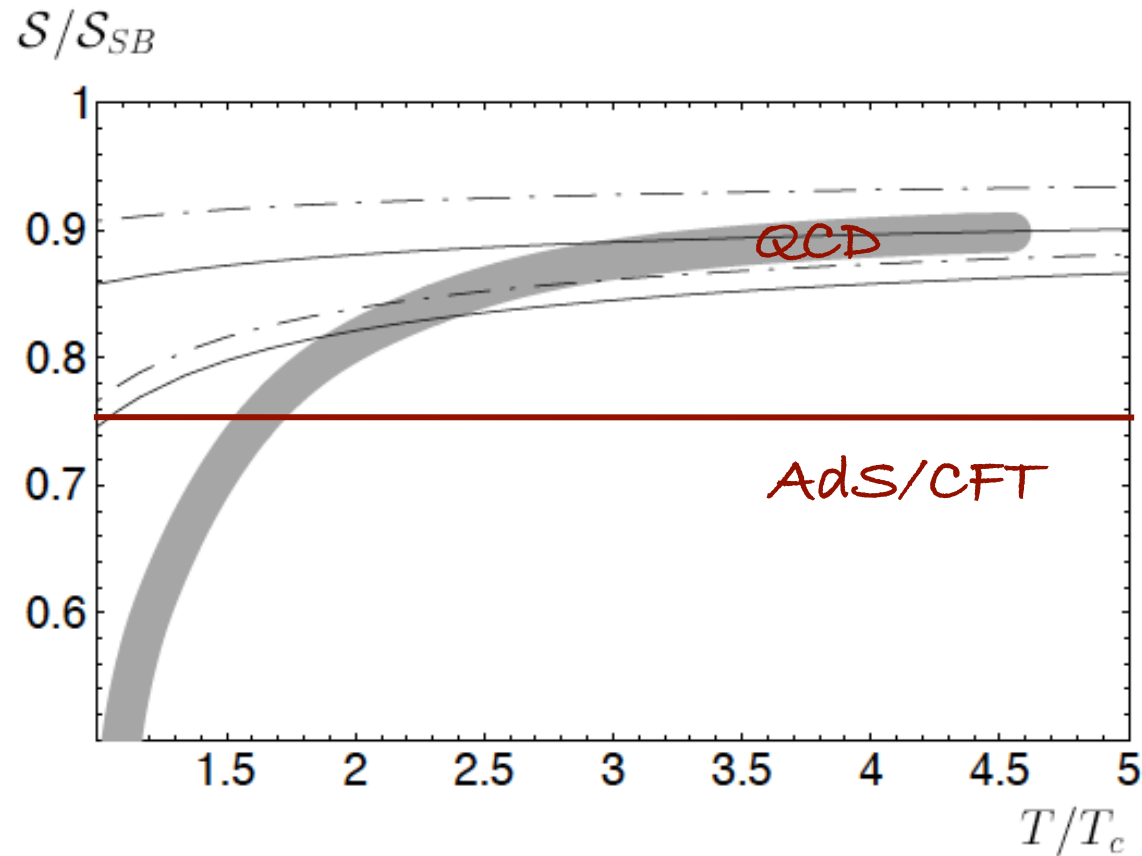
RG techniques yield smooth extrapolation to strong coupling  
(scalar field theory)

(JPB, A. Ipp, N. Wschebor, 2010)



(high orders from J. O. Andersen et al, arXiv 0903.4596)

# Entropy at weak and strong coupling

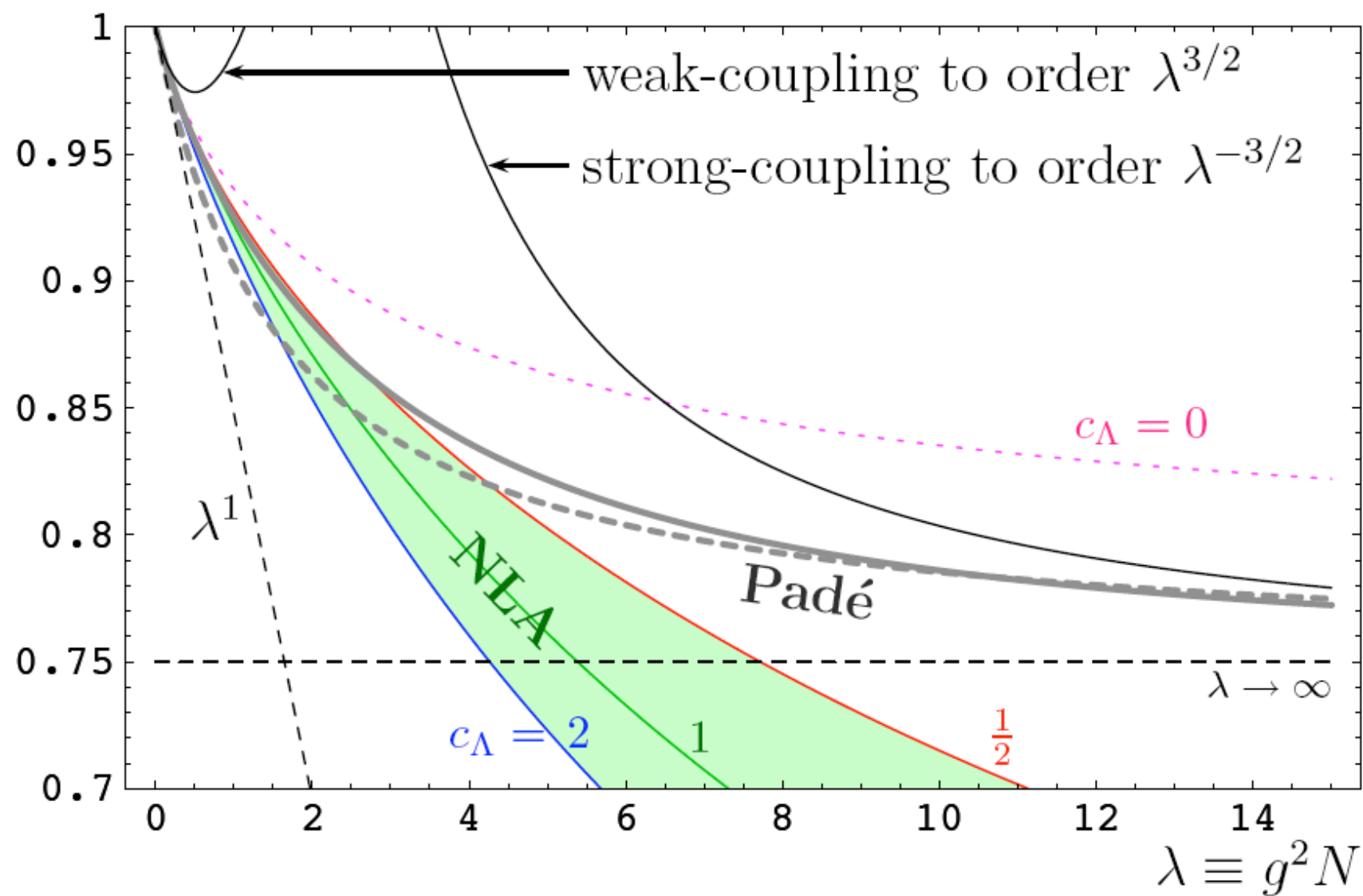


$$\frac{S}{S_{SB}} = \frac{3}{4}$$

J.-P. B., E. Iancu, A. Rebhan: Phys.Rev.D63:065003,2001  
G. Boyd *et al.*, Nucl. Phys. B469, 419 (1996).



# $\mathcal{S}/\mathcal{S}_0$ $\mathcal{N} = 4$ super-Yang-Mills



(JPB, Iancu, Kraemmer, Rebhan, hep-ph/0611393)

Is production of matter in  
heavy ion collisions  
compatible with strong coupling?

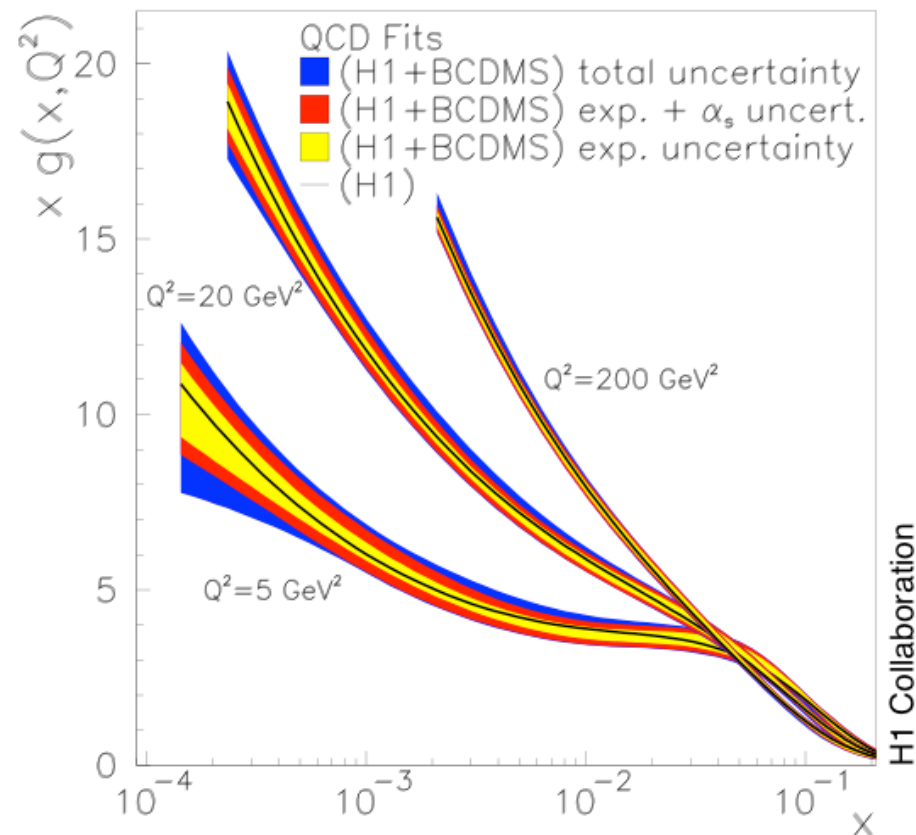
Not really (?)

# Nuclear wave-functions in high energy collisions

Bulk of particle production (  $p_T \lesssim 2 \text{ GeV}$  )

RHIC (  $\sqrt{s} = 200 \text{ GeV}$  )  $x \sim 10^{-2}$

LHC (  $\sqrt{s} = 5.5 \text{ TeV}$  )  $x \sim 4 \times 10^{-4}$



The growth of parton distribution at small  $x$  is tamed by QCD non linear effects (saturation)

Saturation momentum

$$Q_s^2 \approx \alpha_s \frac{xG(x, Q^2)}{\pi R^2} \quad \alpha_s = \alpha_s(Q_s)$$

In a nucleus

$$\frac{xG_A(x, Q^2)}{\pi R^2} \sim A^{1/3}$$

Most partons taking part in collision have  $k_T \sim Q_s$

Successful phenomenology at RHIC

weak coupling but non-perturbative

$$Q_s^2 \approx \alpha_s \frac{xG(x, Q^2)}{\pi R^2}$$

Large parton density at saturation



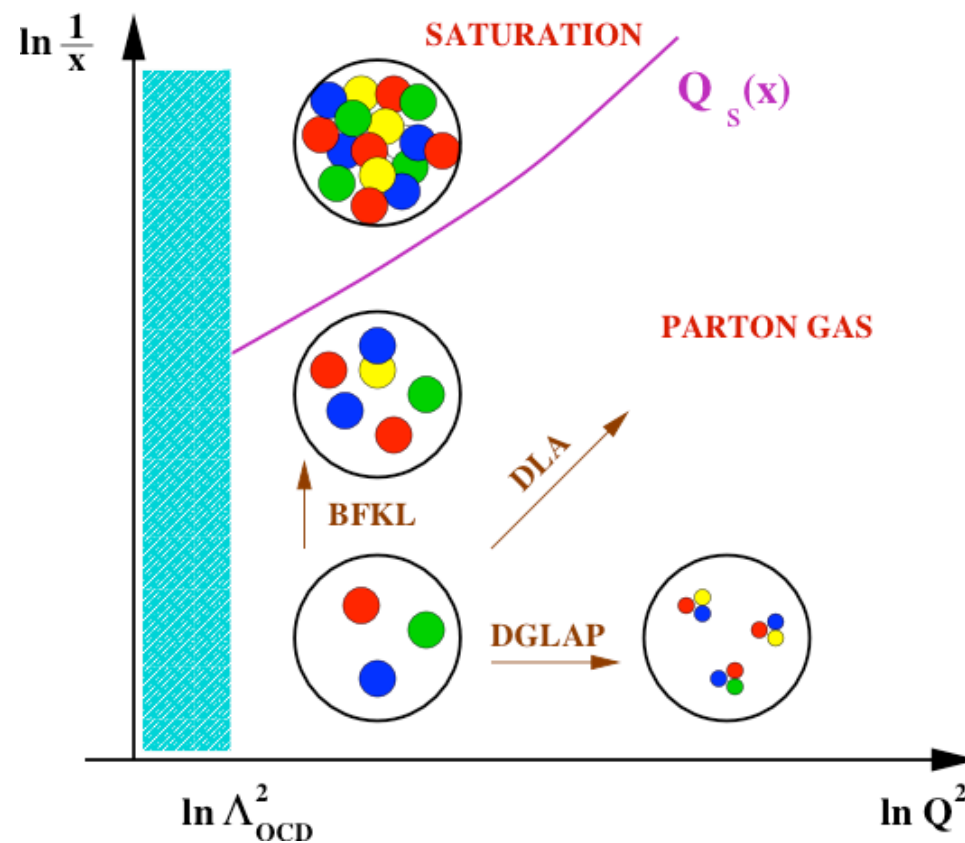
$$n \sim \frac{xG(x, Q^2)}{\pi R^2}$$

$$\frac{\pi}{Q_s^2} n \sim \frac{\pi}{\alpha_s}$$

Suggest a description in terms of classical fields



# Dense and dilute parton systems



# Conclusions

- LHC confirms the strongly coupled character of the QGP
- The strongly coupled character of the quark-gluon plasma does not seem related in any obvious way to a large value of the coupling constant.
- Non perturbative features may arise from the cooperation of many degrees of freedom, or strong classical fields.
- The quark-gluon plasma is a multiscale system (no ideal plasma, neither weakly nor strongly coupled)