

# What the Higgs has to say about New Physics

Veronica Sanz

Sussex

Christmas Workshop, IFT



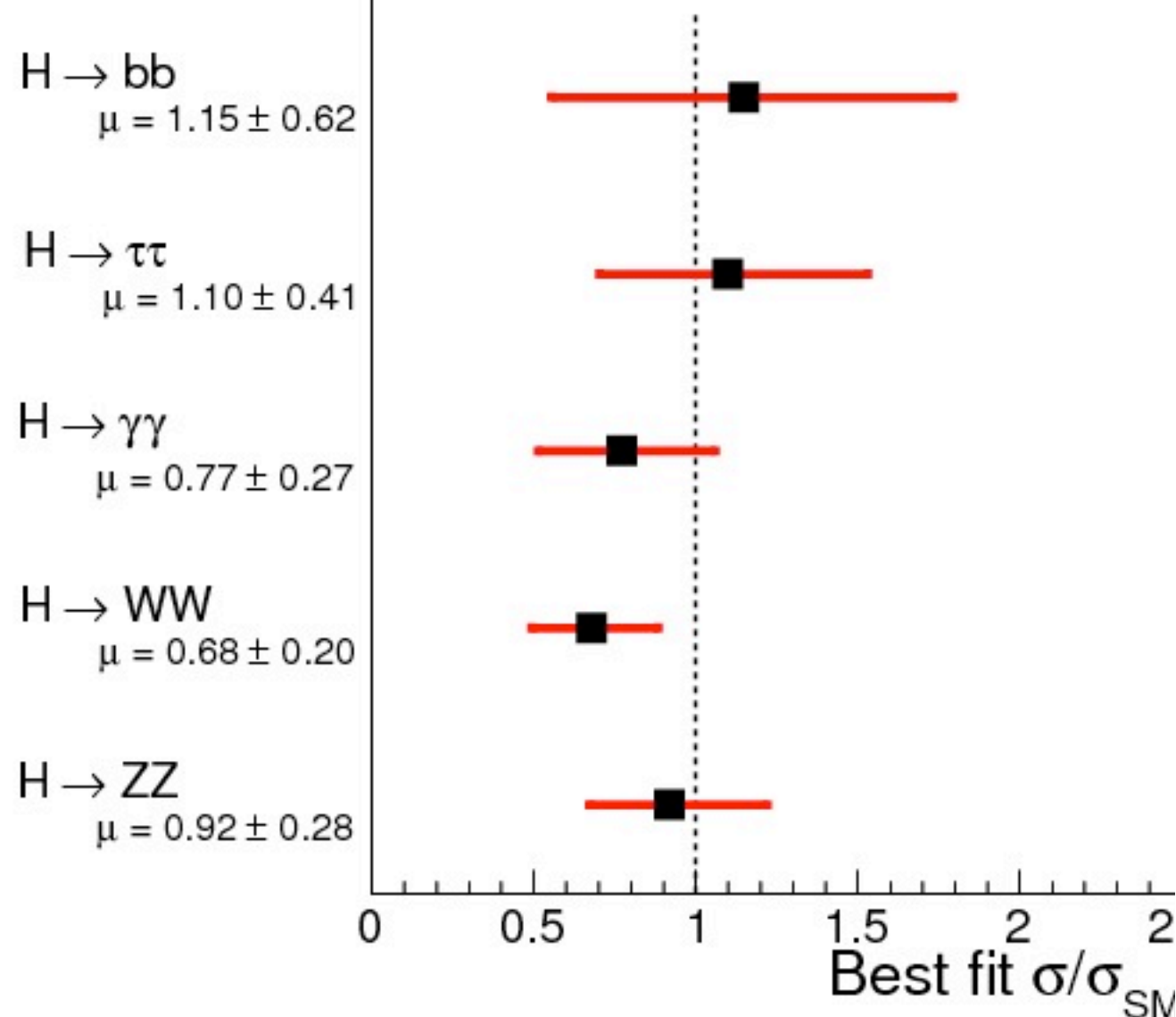
Power to the data\*

\*Quoting Concha and friends



$\sqrt{s} = 7 \text{ TeV}, L \leq 5.1 \text{ fb}^{-1}$   $\sqrt{s} = 8 \text{ TeV}, L \leq 19.6 \text{ fb}^{-1}$

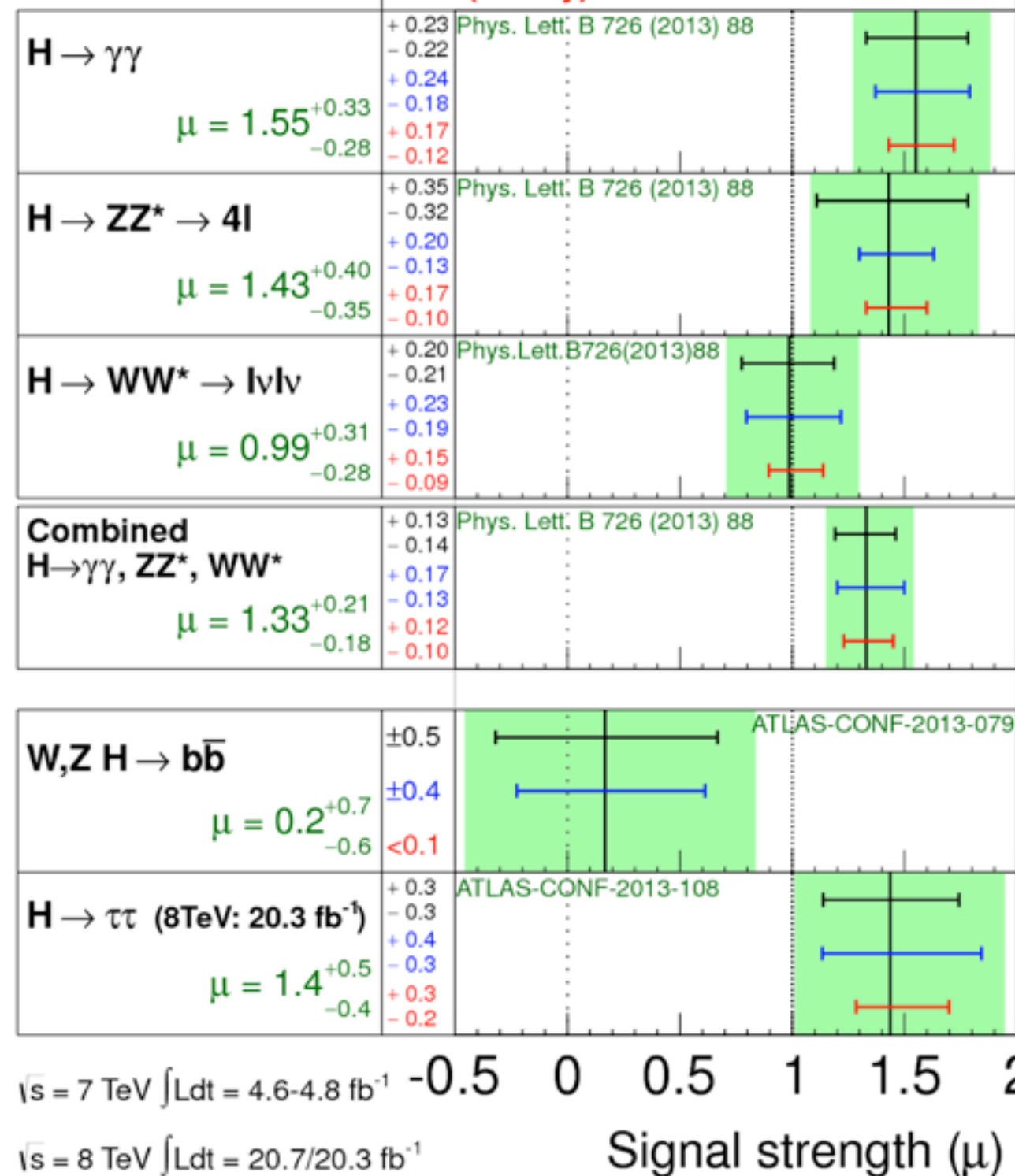
CMS Preliminary  $m_H = 125.7 \text{ GeV}$   
 $\rho_{SM} = 0.65$



**ATLAS Prelim.**

$m_H = 125.5 \text{ GeV}$

—  $\sigma(\text{statistical})$  Total uncertainty  
 —  $\sigma(\text{syst.incl.theo.})$   $\pm 1\sigma$  on  $\mu$   
 —  $\sigma(\text{theory})$



Couplings



Each channel isn't that well measured

BUT

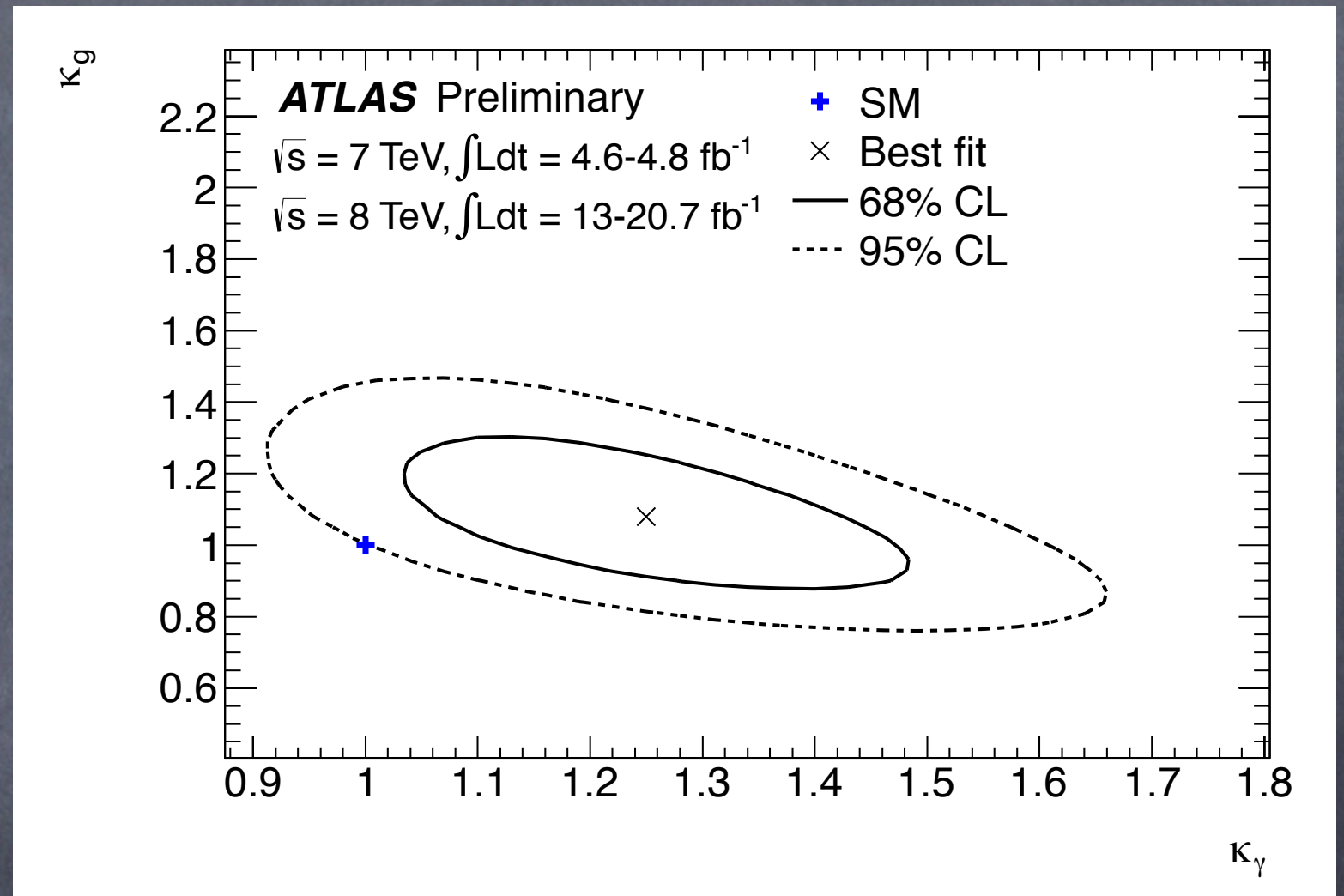
powerful information  
assume a scenario and fit



# Looking for particles in the loops

$$\kappa_g^2 = \frac{\sigma_{prod}(gg \rightarrow h)}{\sigma_{prod}(gg \rightarrow h)_{SM}}$$

$$\kappa_\gamma^2 = \frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma(h \rightarrow \gamma\gamma)_{SM}}$$





From data to ideas



Light SM-like Higgs is a reality  
stabilization of the electroweak scale

use of symmetries  
SUSY and Goldstone (shift)

examples of  
weak and strong coupling

Higgs data  
and its stabilization mechanism



# Supersymmetry



# SUSY Higgs

chiral multiplet with a fermion, Higgsino

$$\Phi = (H, \tilde{H})$$

chiral symmetry  
inherited by the Higgs

e.g. Giudice,  
Rattazzi 97

$$m_h^2 = m_Z^2 \cos^2(2\beta) + (\text{loop}) m_{\tilde{t}}^2 \log(m_{\tilde{t}_2}/m_{\tilde{t}_1})$$



# SUSY Higgs and the third generation

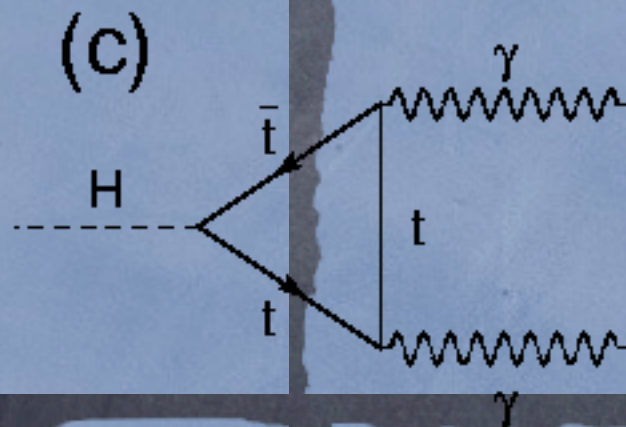
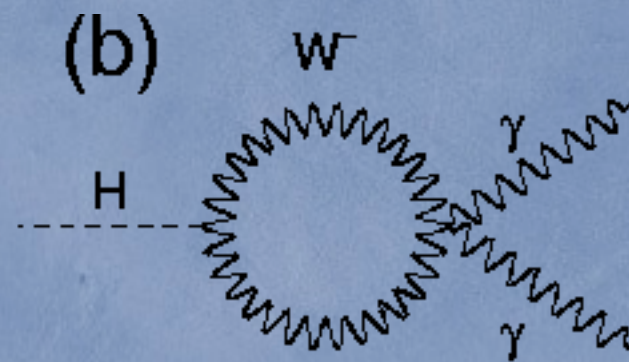
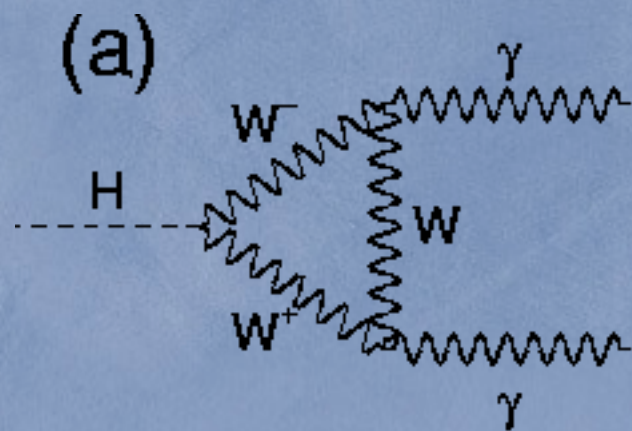
Light stops couple to the Higgs and affect its  
production and decay modes



# SUSY Higgs and the third generation

Light stops couple to the Higgs and affect its production and decay modes

## SM contributions



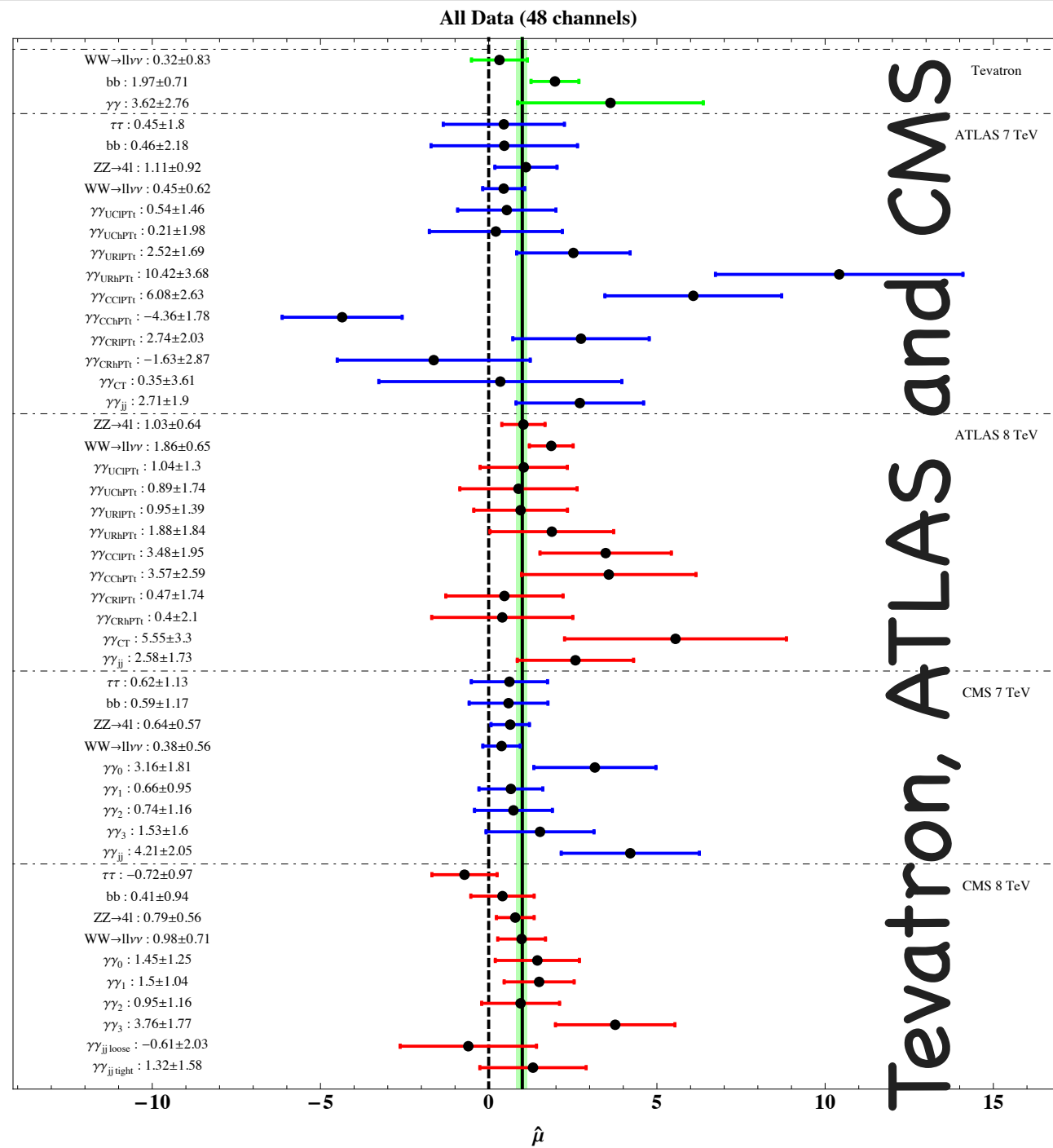
+ stops running in the loop

$$F_{\gamma}^{SM}(m_t, W, m_b \dots) = F_1(\tau_W) + \sum_{i=t,b,\dots} N_c Q_i^2 F_{1/2}(\tau_i) \left(1 - \frac{\alpha_s}{\pi}\right) \approx 1/(0.155 - 0.002 i)$$



# Fit to Higgs properties in BSM

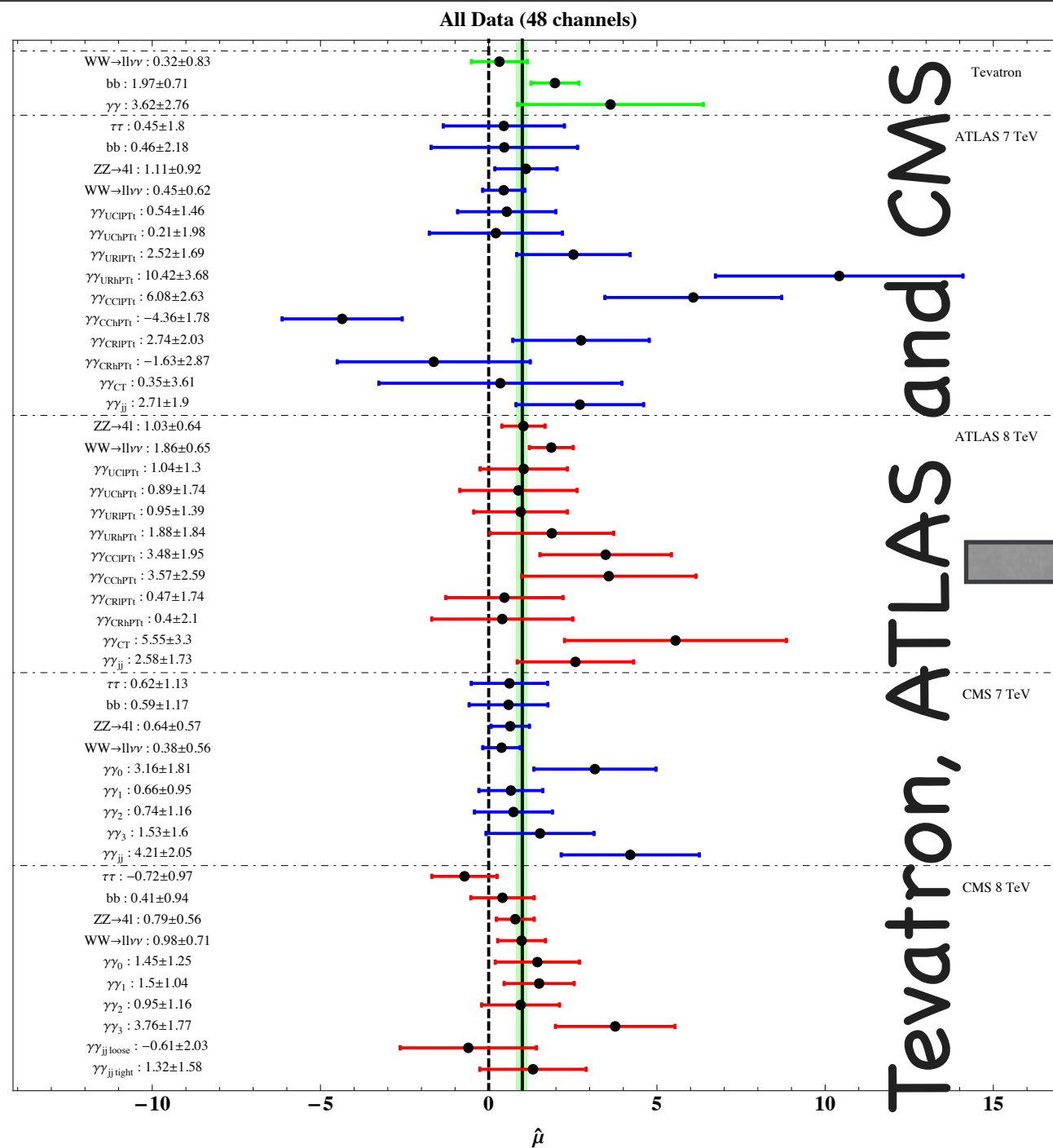
## Espinosa, Grojean, VS, Trott, 2012





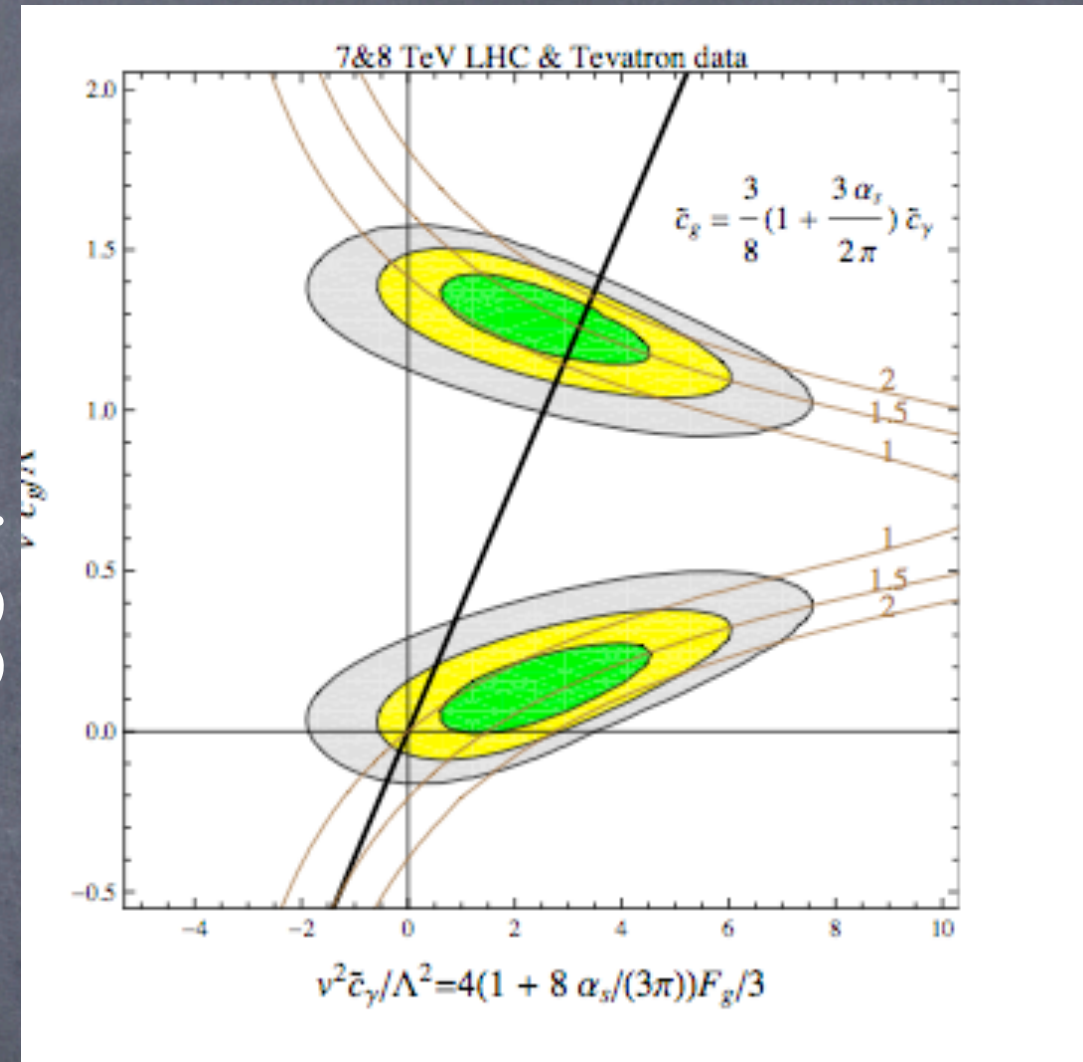
# Fit to Higgs properties in BSM

Espinosa, Grojean, VS, Trott, 2012



Tevatron, ATLAS and CMS

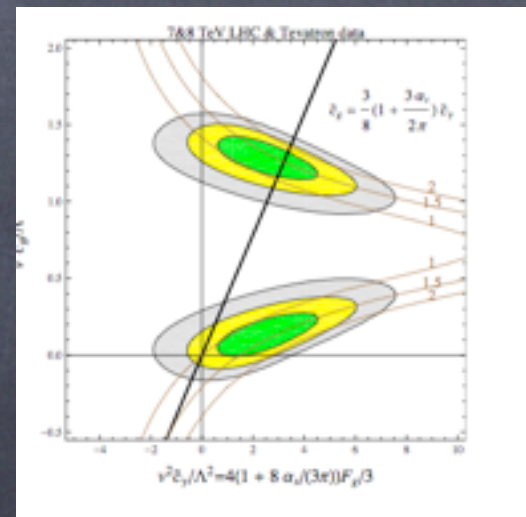
effect on gg production



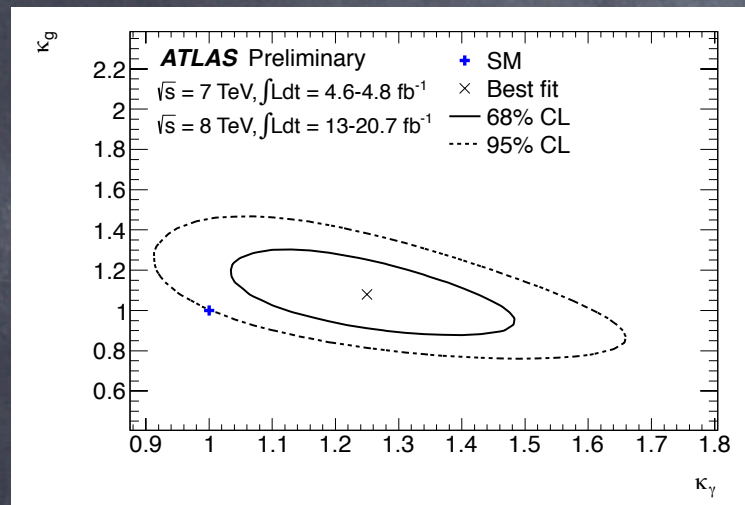
effect on diphoton

green, yellow, gray= 1,2 and 3 sigma

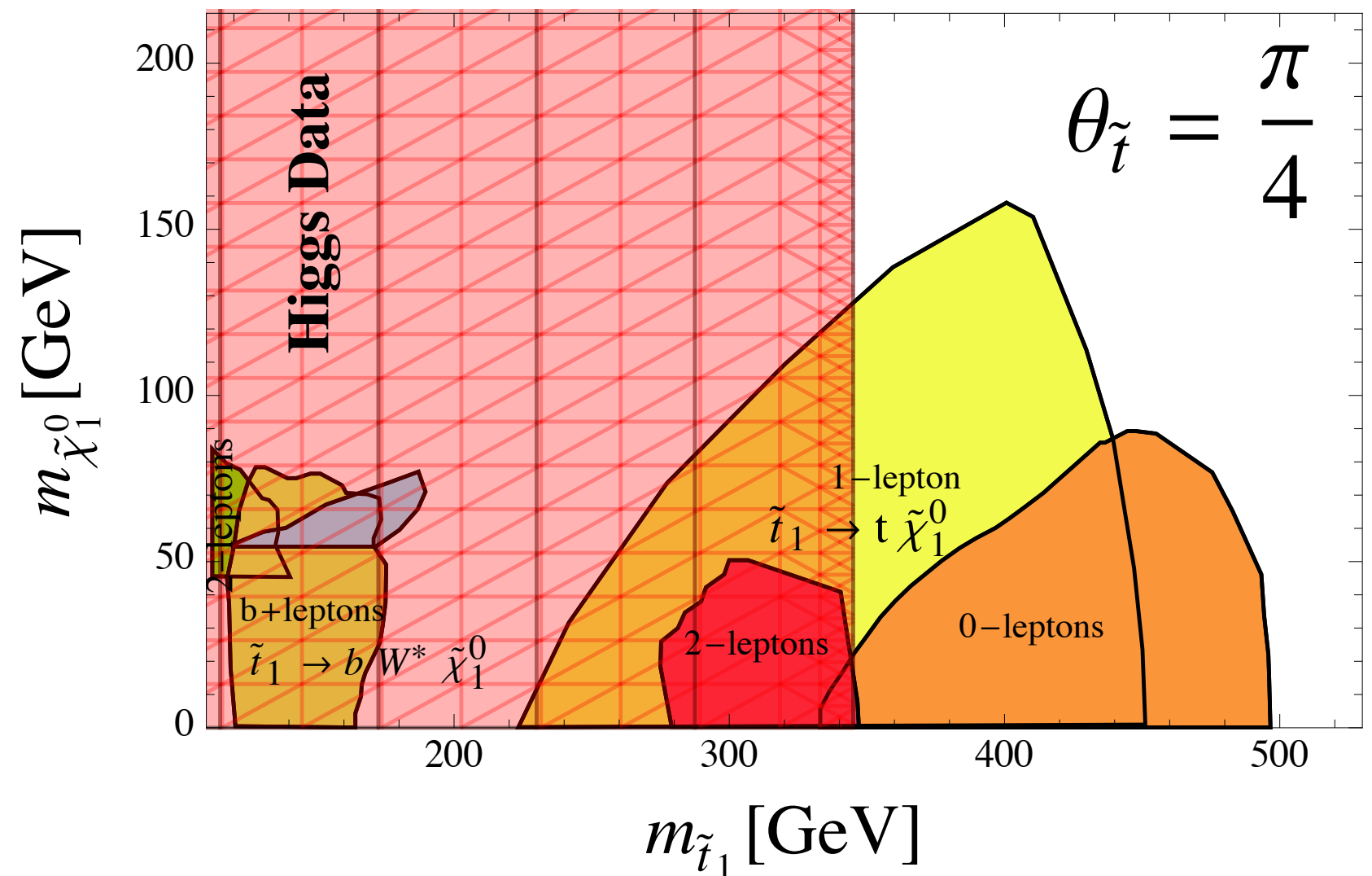




+ mW + b to s gamma  $\longrightarrow$  Stops

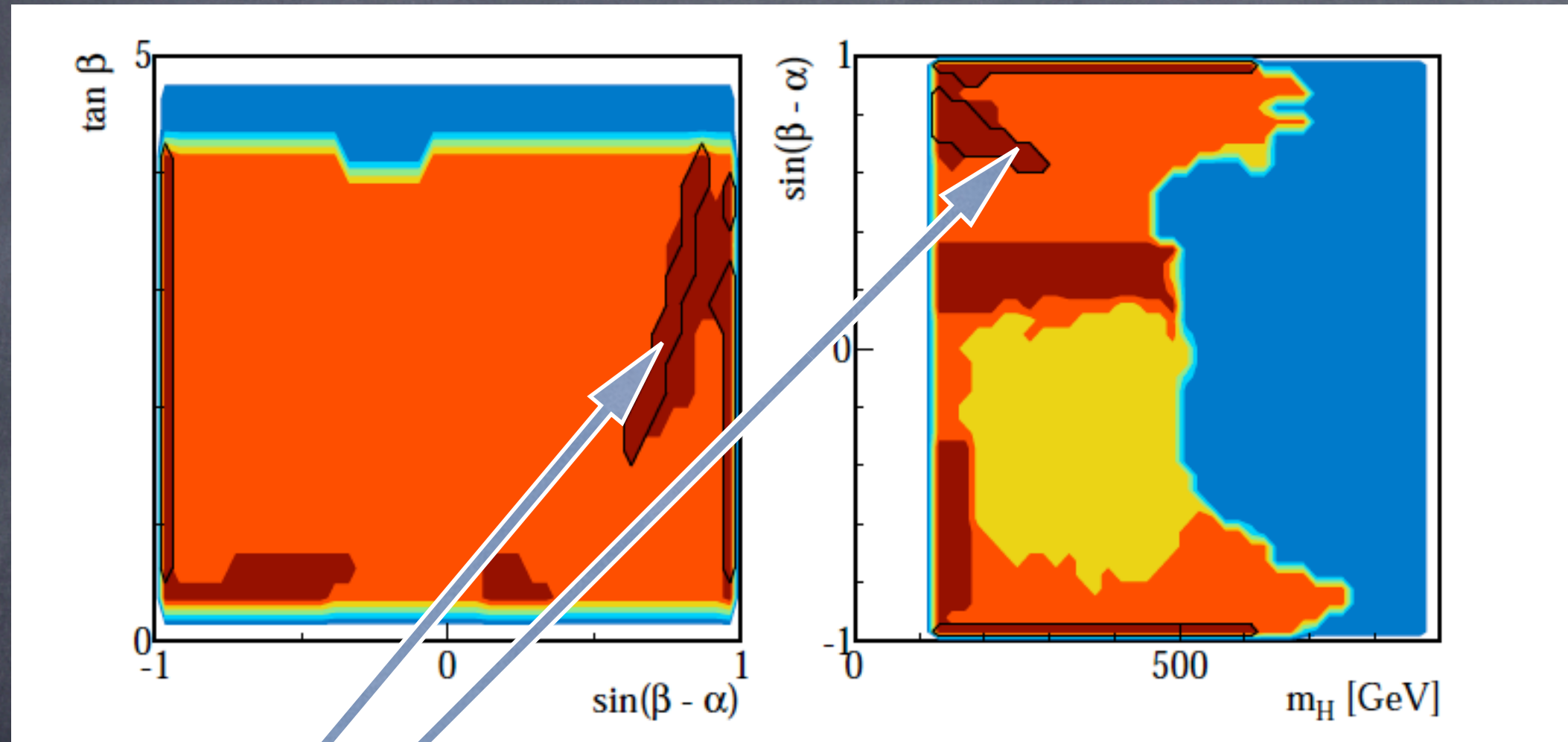


95% CL  
exclusion





In SUSY, parameter space of 2HDM



Allowed  
regions

Coleppa 1305.0002





# Composite Higgs



Another option

Higgs as a pseudo-Goldstone boson

Composite Higgs

scalar doublet, CP even



Another option

Higgs as a pseudo-Goldstone boson

Composite Higgs

scalar doublet, CP even



rho, a1, f2...  
scale symmetry  
breaking: f

$$\xi = v^2 / f^2$$

EW scale: v  
pions

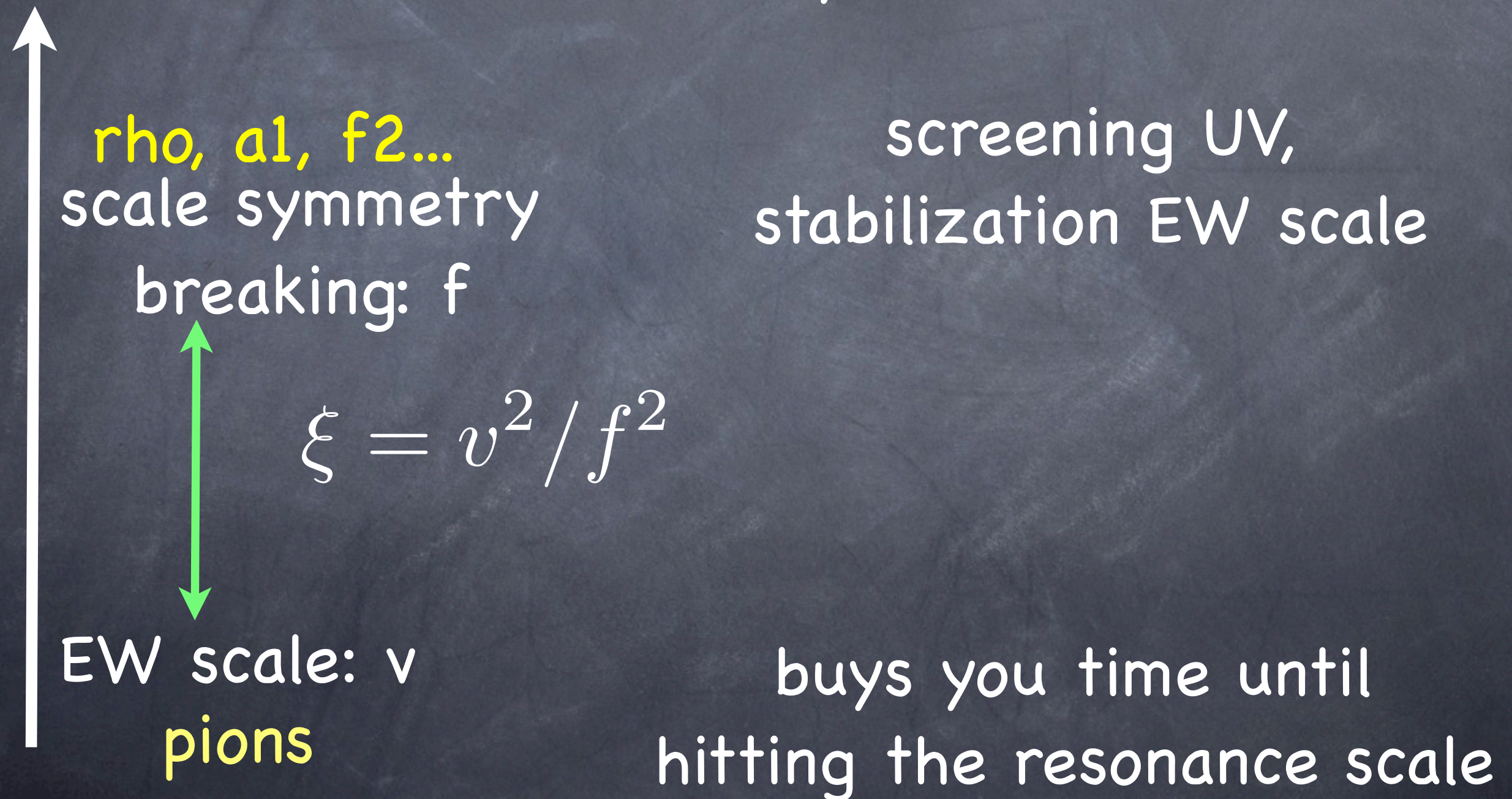


Another option

Higgs as a pseudo-Goldstone boson

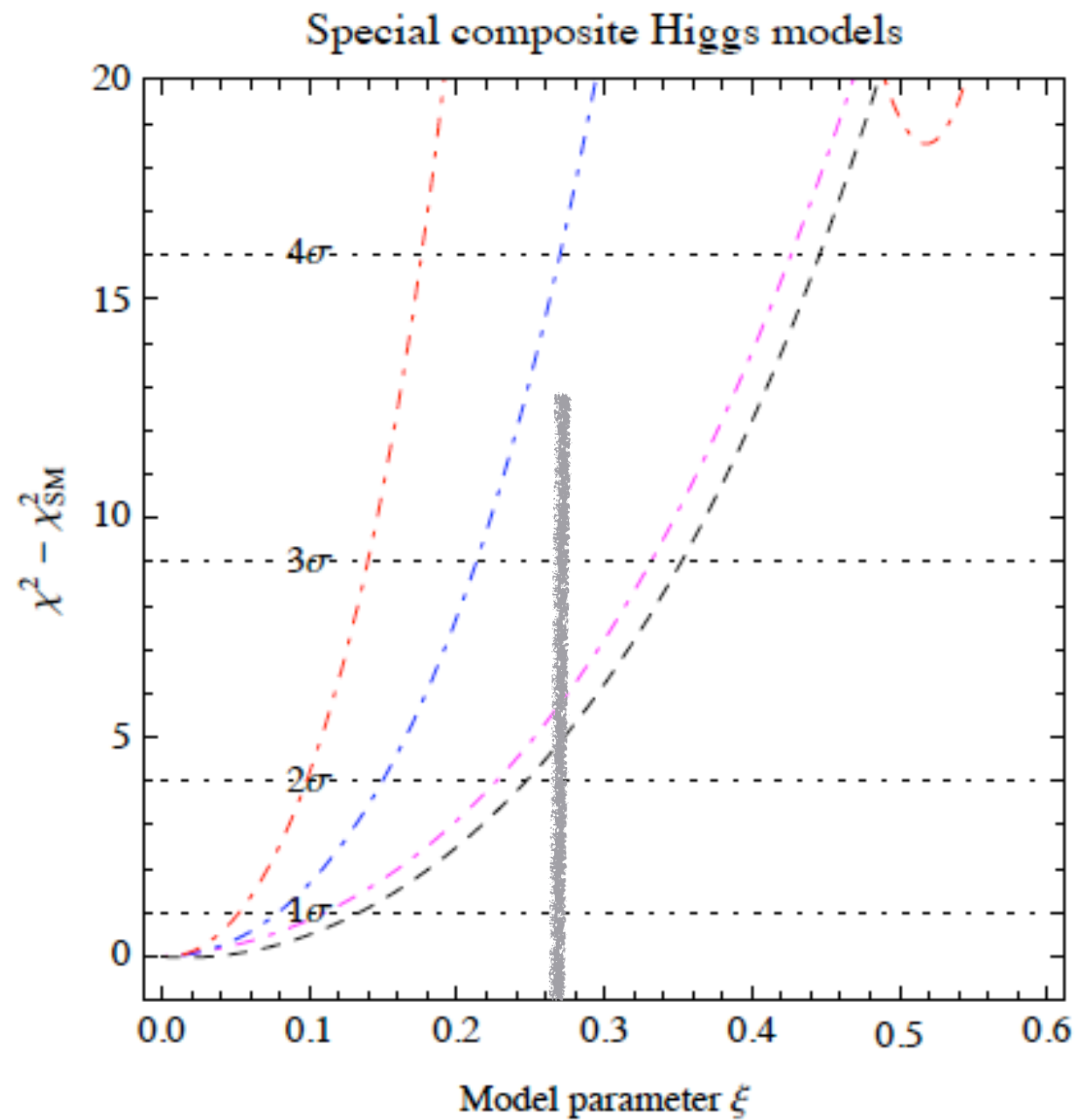
Composite Higgs

scalar doublet, CP even





The composite Higgs doesn't have the same couplings as the SM higgs



Higgs data  
no unnaturalness

$$f \gtrsim 500 \text{ GeV}$$

Strumia et al. 2013



This is (was) not **mainstream**, why?

original idea  
Georgi et al, 82

**strong** coupling  
but

the main hurdle is generating a potential  
(mass, quartic...)



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Georgi et al, 82

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but

the main hurdle is generating a potential  
(mass, quartic...)

$$\underline{m^2} \propto \underline{\lambda v^2}$$

protected  
same suppression

expect

$$v \sim f$$



But this is a symmetry argument  
in explicit models one can search for mechanisms

explicit models dealing strong coupling?

Holographic pseudoGB  
(extra-dimensions)  
and partial compositeness

Pomarol et al. 04  
Contino et al. 07

A lot of model building after that...



Higgs mass  
top-partner contribution

$$m_h^2 \sim \frac{N_c y_t^2}{16\pi^2} \frac{v^2}{f^2} m_T^2$$

One could still push

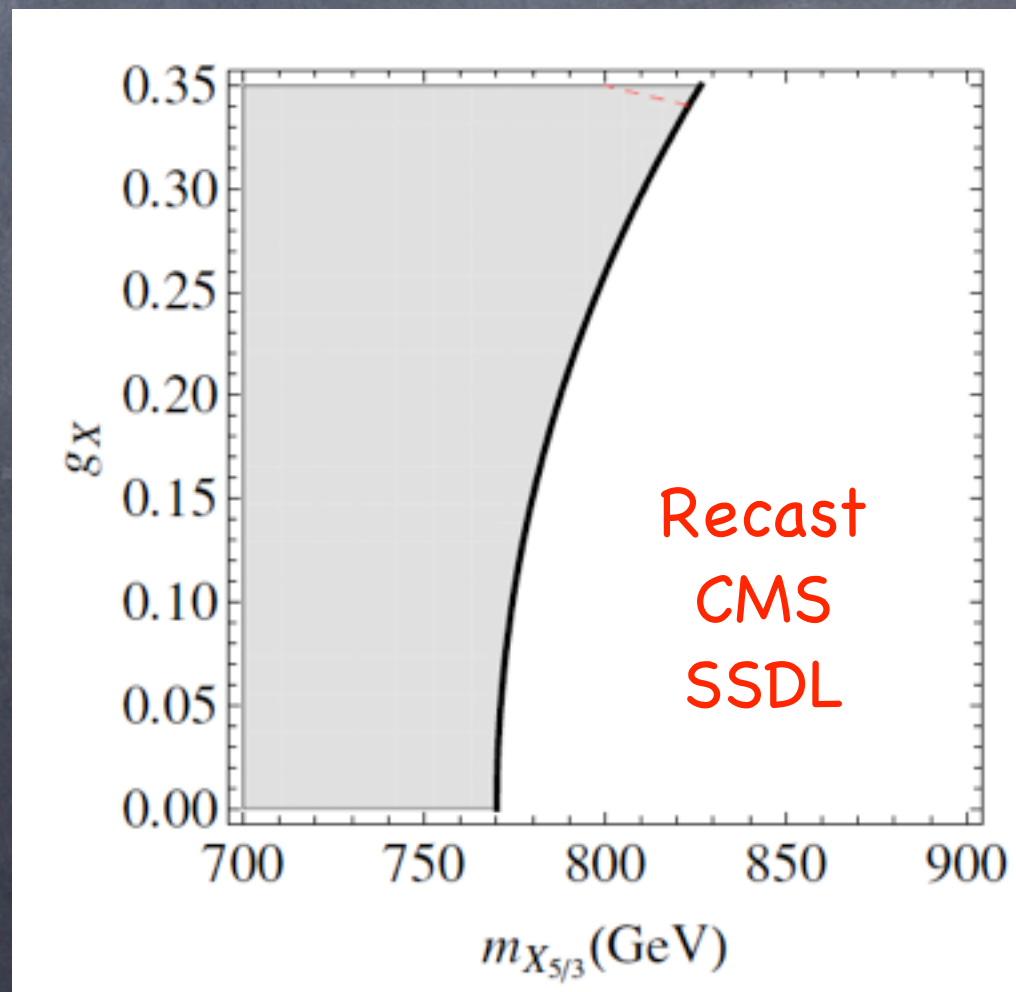
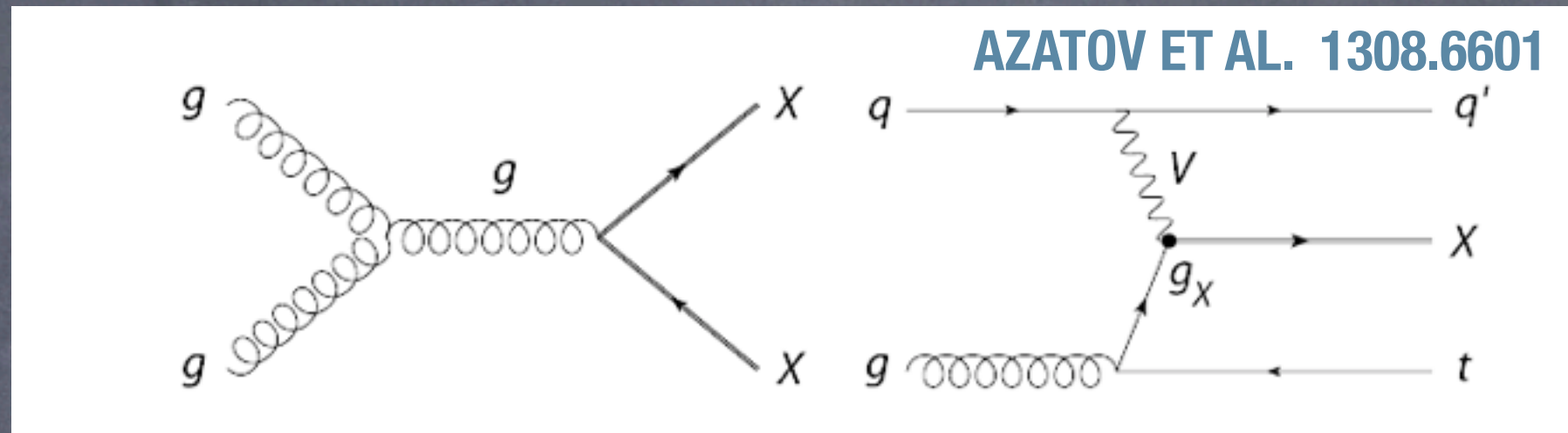
$$f \sim TeV$$

but

$$m_T \lesssim TeV$$



$$m_T \lesssim TeV$$





# Higgs Effective Lagrangian

## The LEP mentality



Now, let's be agnostic

how can we use Higgs data to constrain  
NP in a "model-independent" way?

Assume NP is heavy  
Higgs anomalous couplings



Effective theory approach:  
ex. operators involving the HVV vertices

$$\mathcal{L}_{eff} = \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i$$

$$\epsilon_i = f_i \frac{v^2}{\Lambda^2}$$

a set of  
operators

$$\mathcal{O}_W = (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_B = (D_\mu \Phi)^\dagger (D_\nu \Phi) \widehat{B}^{\mu\nu}$$

$$\mathcal{O}_{WW} = \Phi^\dagger \widehat{W}^{\mu\nu} \widehat{W}_{\mu\nu} \Phi$$

$$\mathcal{O}_{BB} = (\Phi^\dagger \Phi) \widehat{B}^{\mu\nu} \widehat{B}_{\mu\nu}$$

operators contributing to tree S,T:  
set them to zero? running

Grojean et al. 2013  
Pomarol et al. 2013  
most recent  
1312.2928 & 1312.2014



$$\mathcal{O}_W = (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi)$$

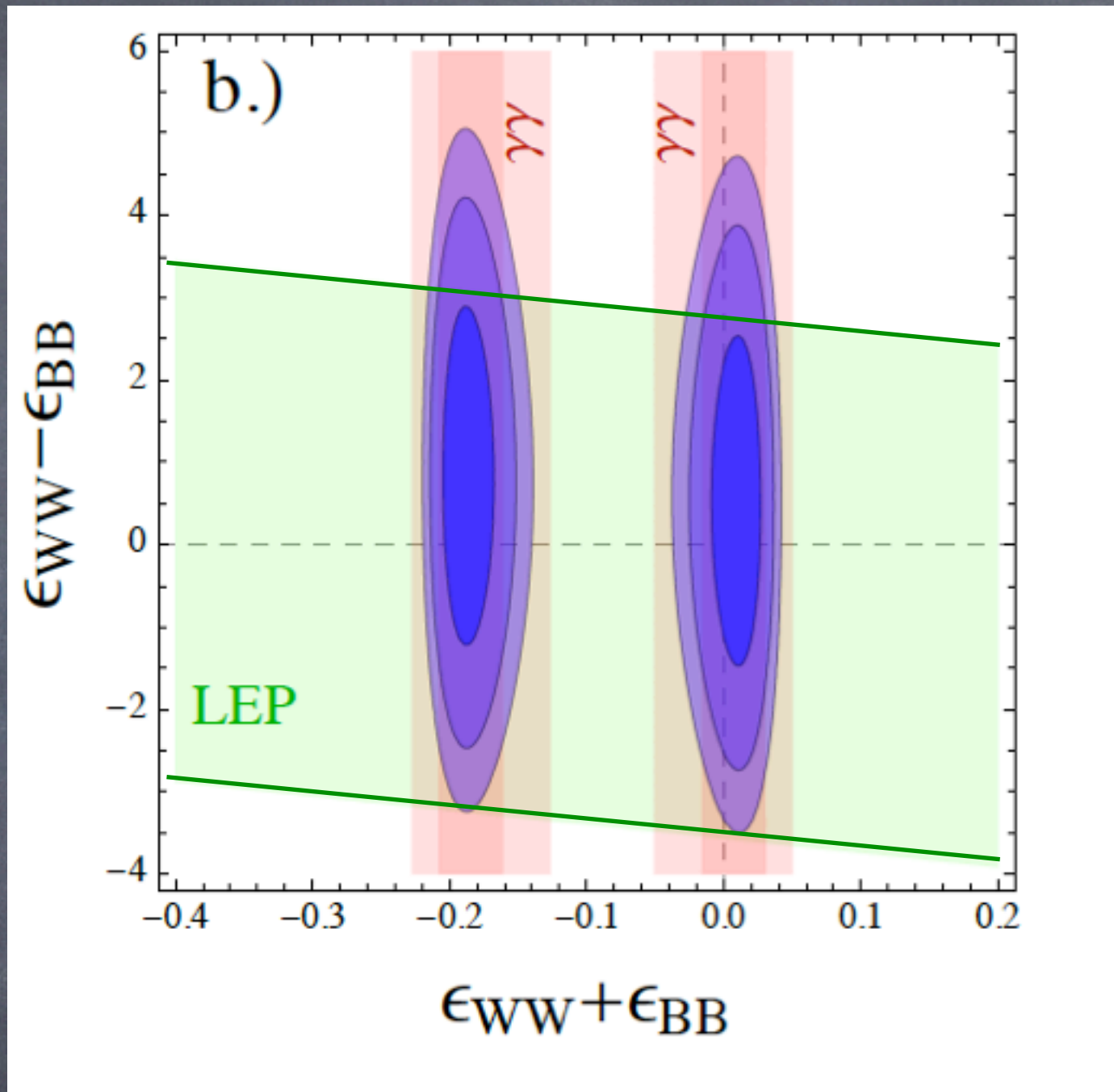
$$\mathcal{O}_B = (D_\mu \Phi)^\dagger (D_\nu \Phi) \widehat{B}^{\mu\nu}$$

$$\mathcal{O}_{WW} = \Phi^\dagger \widehat{W}^{\mu\nu} \widehat{W}_{\mu\nu} \Phi$$

$$\mathcal{O}_{BB} = (\Phi^\dagger \Phi) \widehat{B}^{\mu\nu} \widehat{B}_{\mu\nu}$$

eWW and eBB uses mostly  $h \rightarrow aa$

LHC data  
surpasses  
PEWT



Masso and VS  
PRD 2012



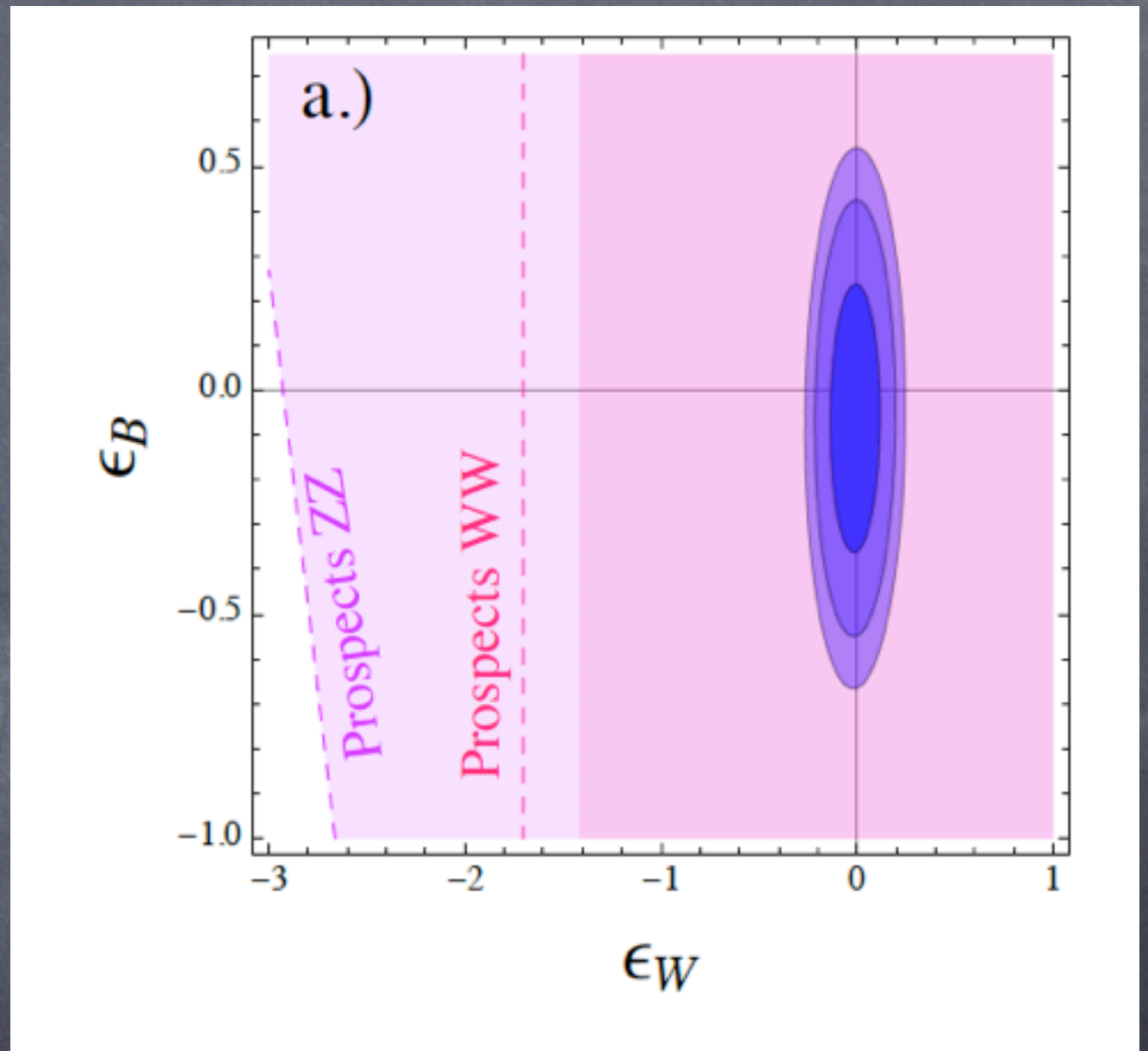
$$\mathcal{O}_W = (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi)$$

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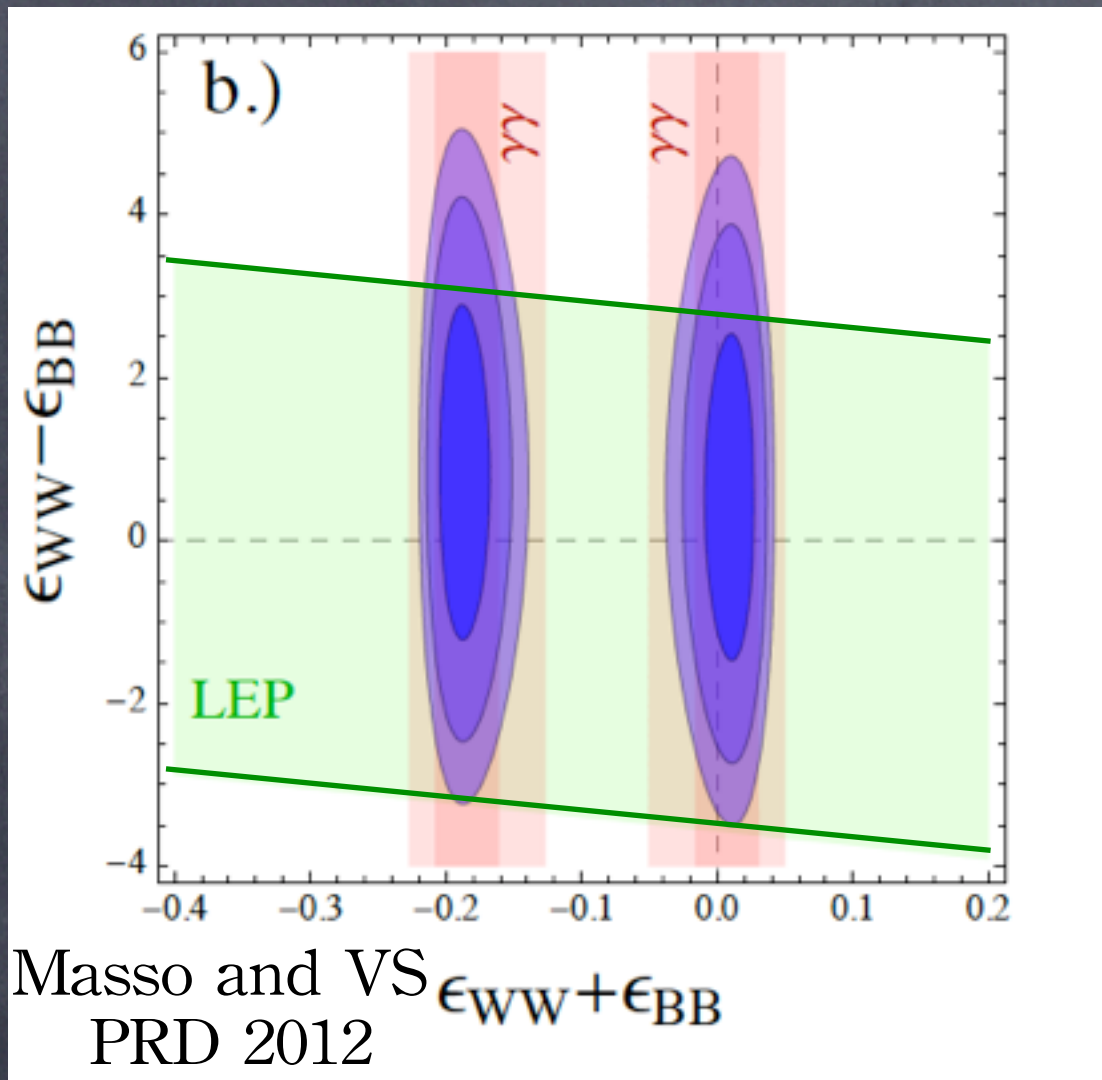
But when it comes to  $eW$  and  $eB$ ,  
data  $h \rightarrow WW, ZZ$



More ideas? differential  
information?



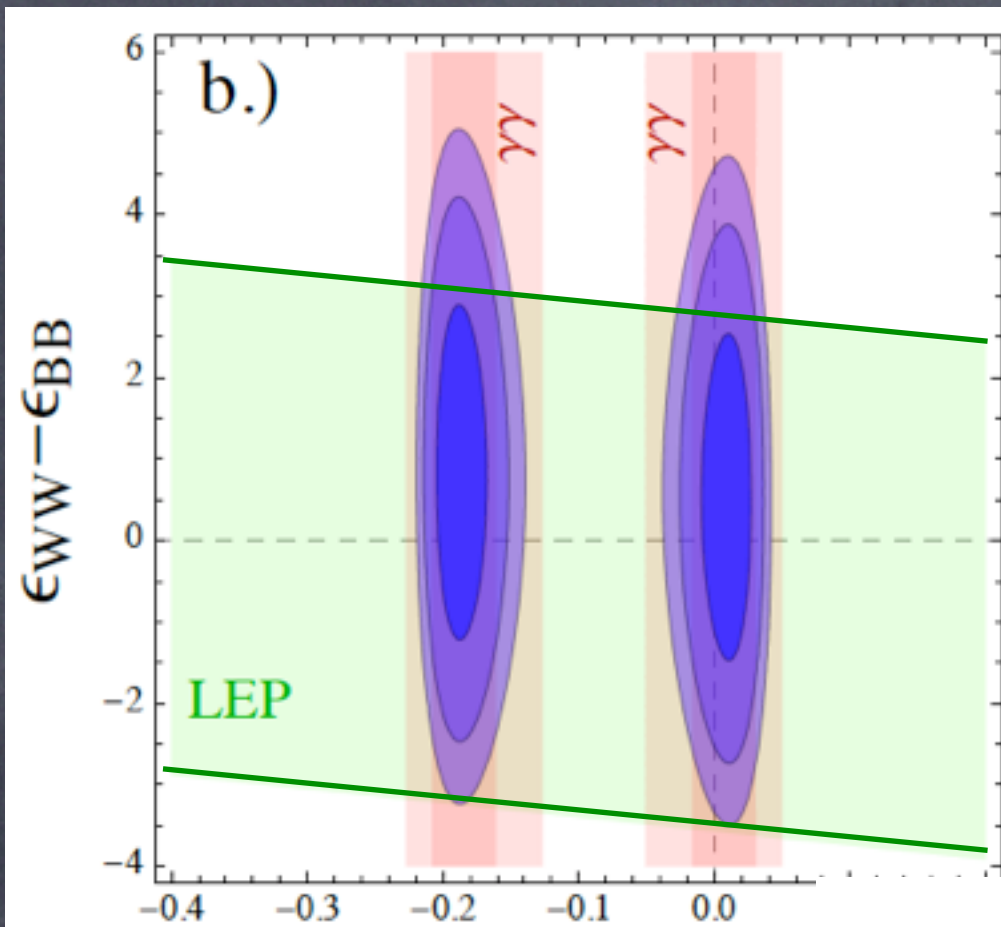
# Is the LEP/LHC comparison fair?



Not clear  
Dawson et al 1311.3107



# Is the LEP/LHC comparison fair?



Masso and VS  
PRD 2012

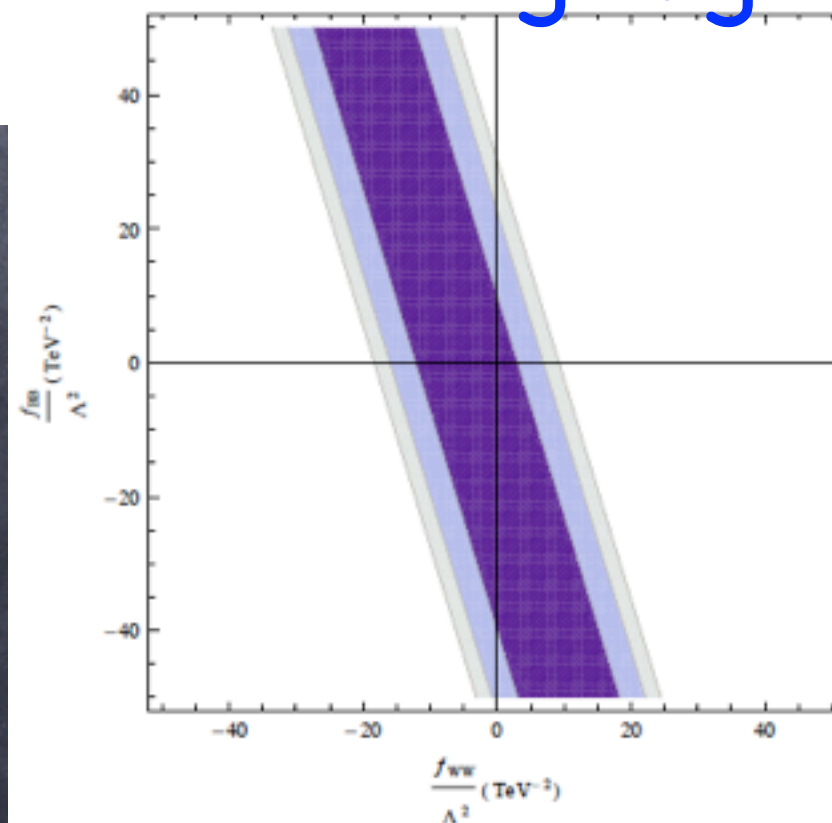
oblique  
parameter  
constraints  
MUCH weaker  
than thought  
before

Not clear

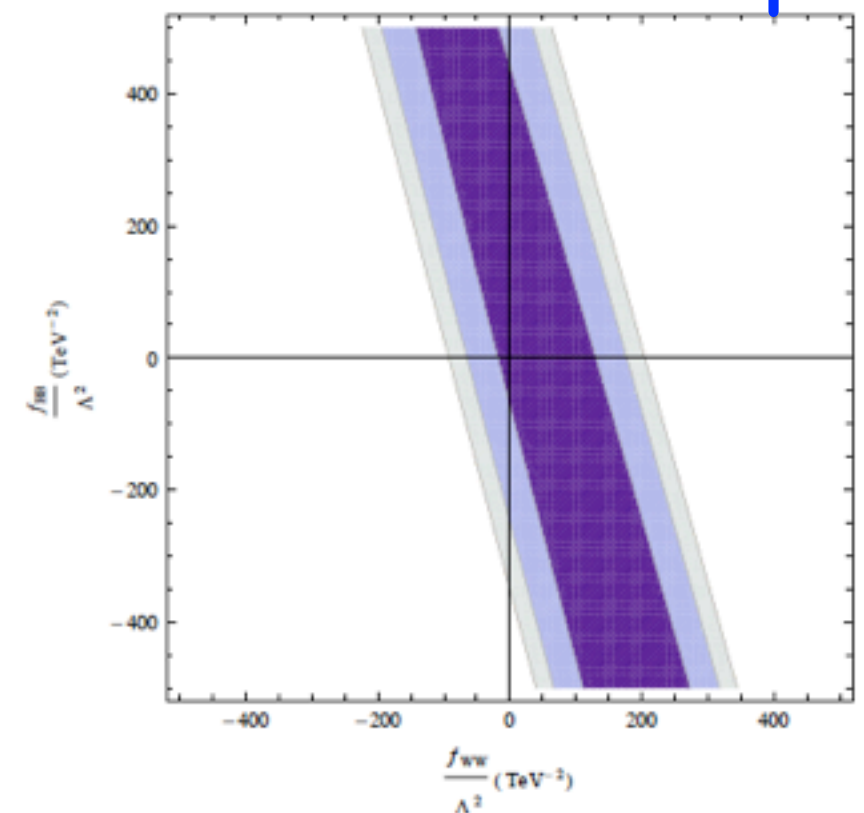
Dawson et al 1311.3107

dim6 operators at loop  
level to S,T,U  
renormalization scheme  
dependence, gauge invariance

leading logL



Ren Sch dep

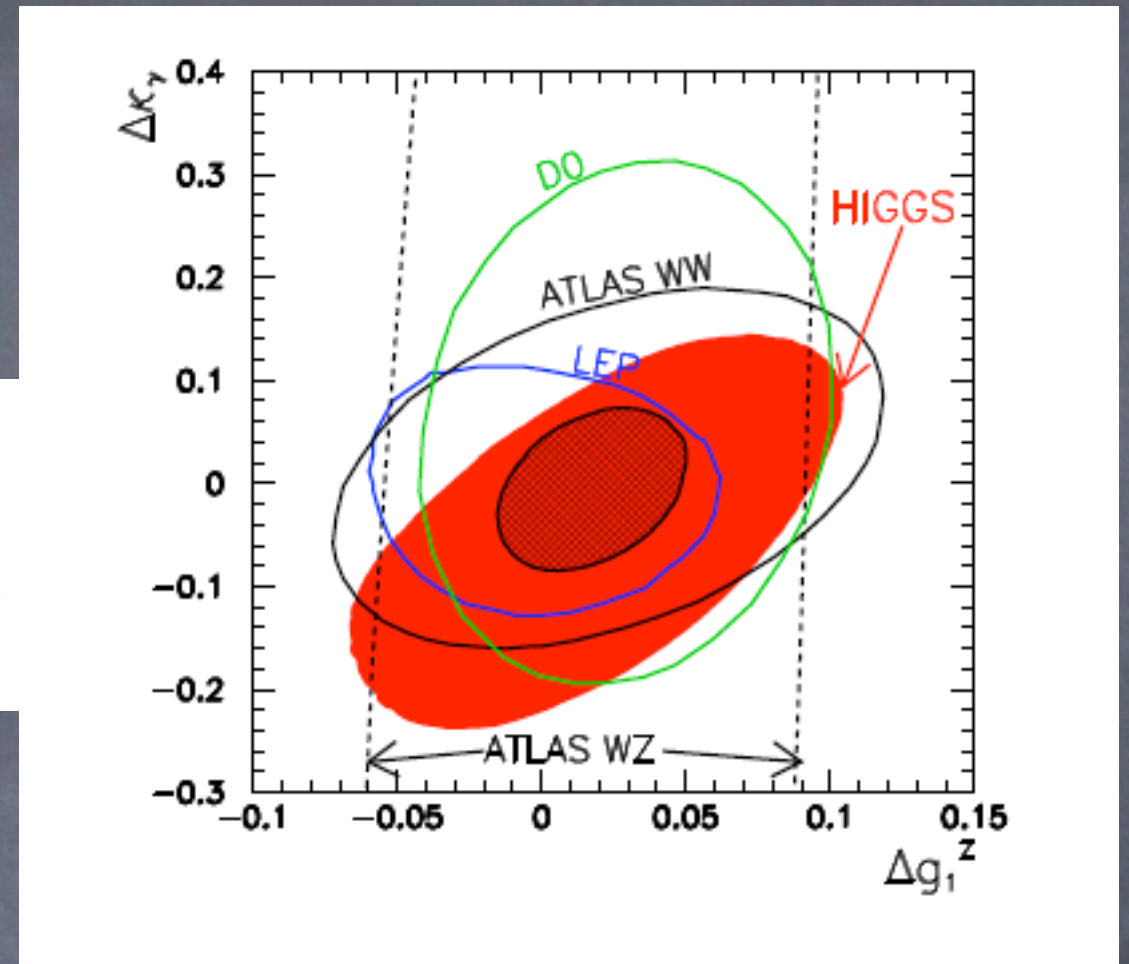




## Correlate HVV with VVV

$$\mathcal{L}_{WWV} = -ig_{WWV} \left\{ g_1^V \left( W_{\mu\nu}^+ W^{-\mu} V^\nu - W_\mu^+ V_\nu W^{-\mu\nu} \right) + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda_V}{m_W^2} W_{\mu\nu}^+ W^{-\nu\rho} V_\rho^\mu \right\}, \quad (4)$$

Corbett et al PRL 111 (2013)



## Linear vs non-linear realization

Correlations depend on realization

dim8 in linear realization may be sizeable in the non-linear expansion

Brivio et al 1311.1823



From total rates to  
differential information

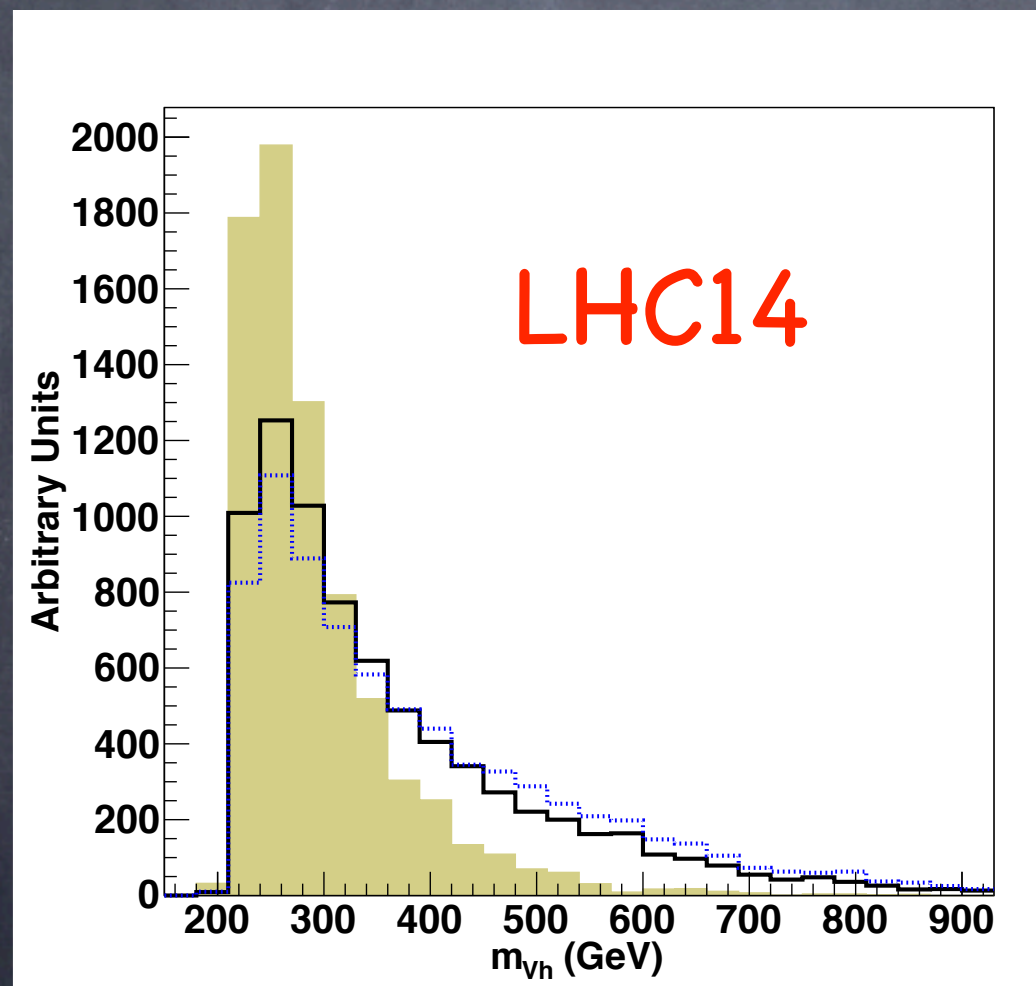


New Physics  
different Lorentz structure  
boosted topologies & COM dependence



# New Physics

different Lorentz structure  
boosted topologies & COM dependence



Alloul, Fuks, VS. 1310.5150

Invariant mass in  
associated  
production w/ HEFT

not enough data yet for  
this, but instead

ratios of COM



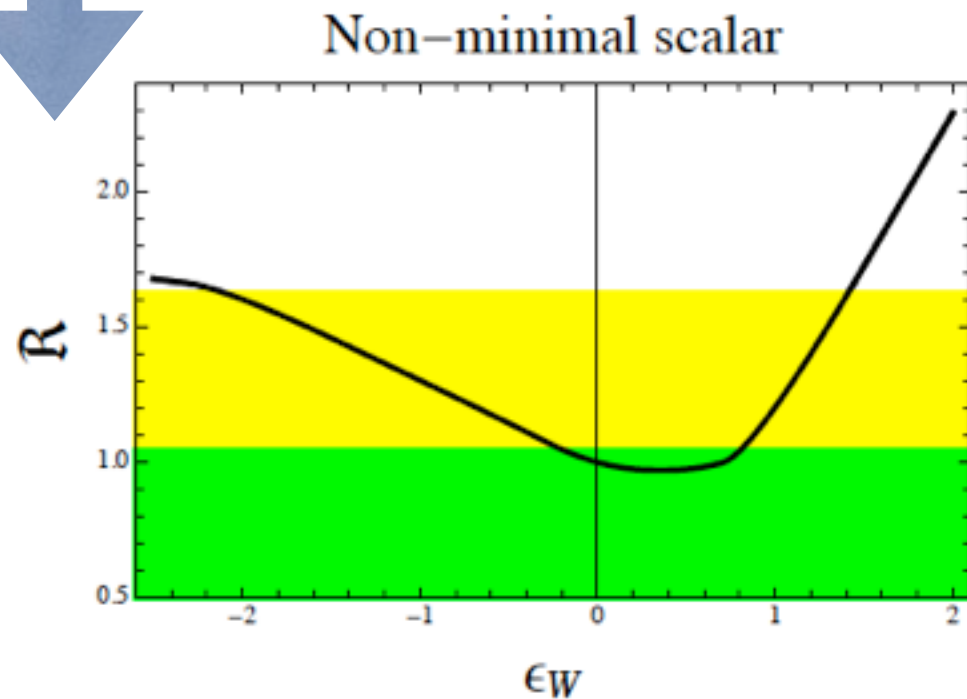
# Energy dependence in associated production

Ellis, VS, You. 2013

$$pp \rightarrow V^* \rightarrow V h$$

double ratio

$$\mathcal{R}_{\text{data}} \equiv \left( \frac{\sigma_{\text{CMS LHC 8}}^{\text{data}}}{\sigma_{\text{TeVatron}}^{\text{data}}} \right) / \left( \frac{\sigma_{\text{LHC 8}}^{0+}}{\sigma_{\text{TeVatron}}^{0+}} \right) = \frac{\sigma_{\text{TeVatron}}^{0+} \sigma_{\text{CMS LHC 8}}^{\text{data}}}{\sigma_{\text{TeVatron}}^{\text{data}} \sigma_{\text{LHC 8}}^{0+}} = \frac{\mu_{\text{LHC 8}}}{\mu_{\text{TeVatron}}} = 0.47 \pm 0.58 . \quad (16)$$



This 'differential' information is more stringent than a global fit to LHC8 data

Effective operator



With more data  
one can start asking more questions



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one can start asking more questions

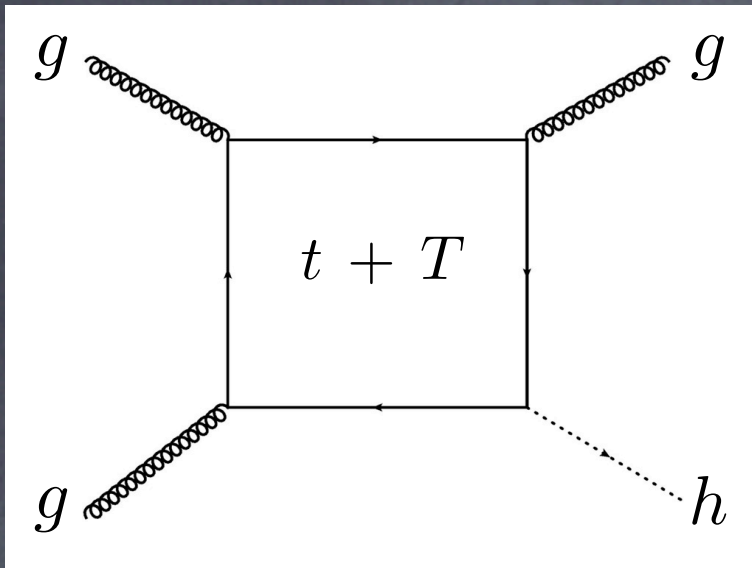
Composite Higgs and Little Higgs  
top-partners affect the loops  
BUT  
low-energy cancellation

No information on top-partners  
from Higgs couplings



# Higher order corrections evade the cancellation

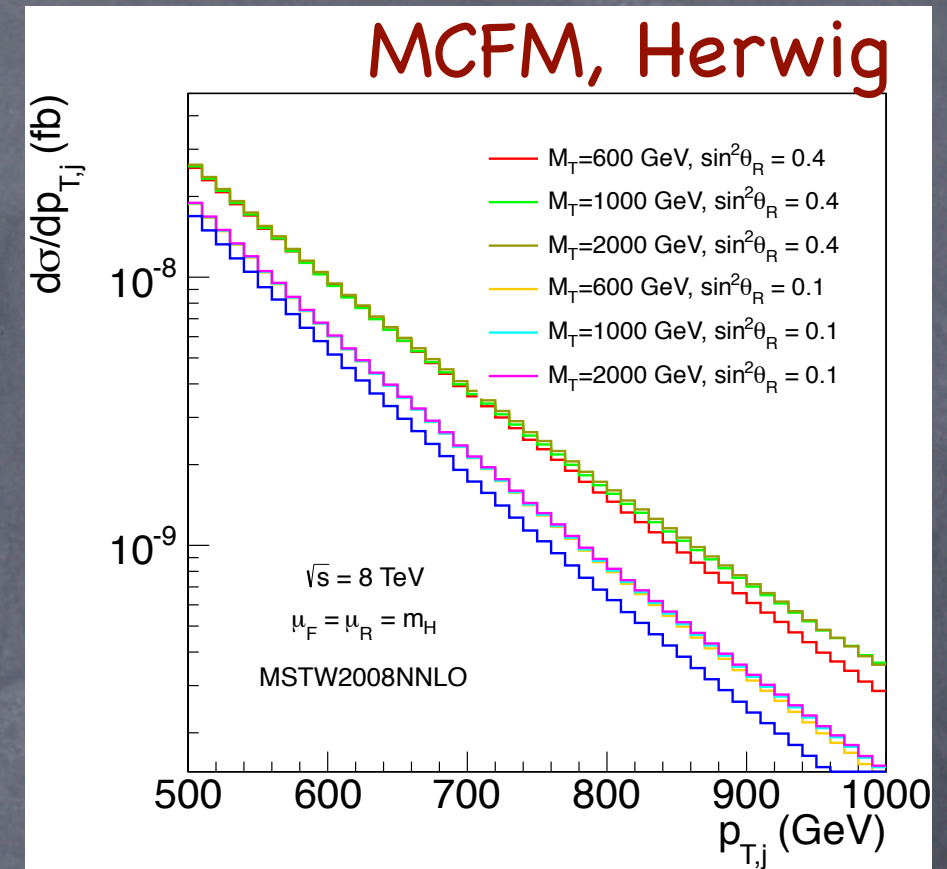
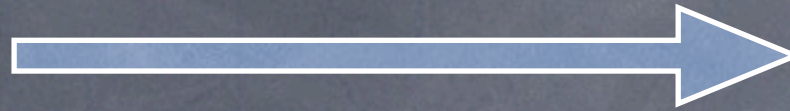
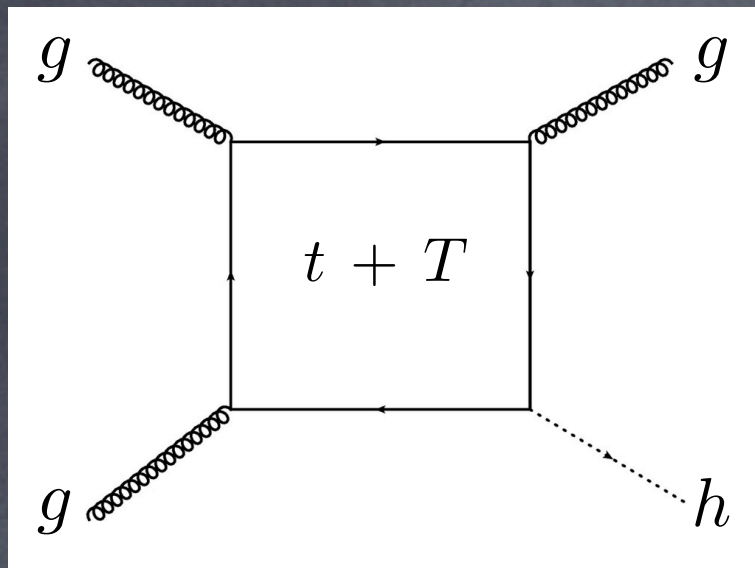
Banfi, Martin, VS. 2013





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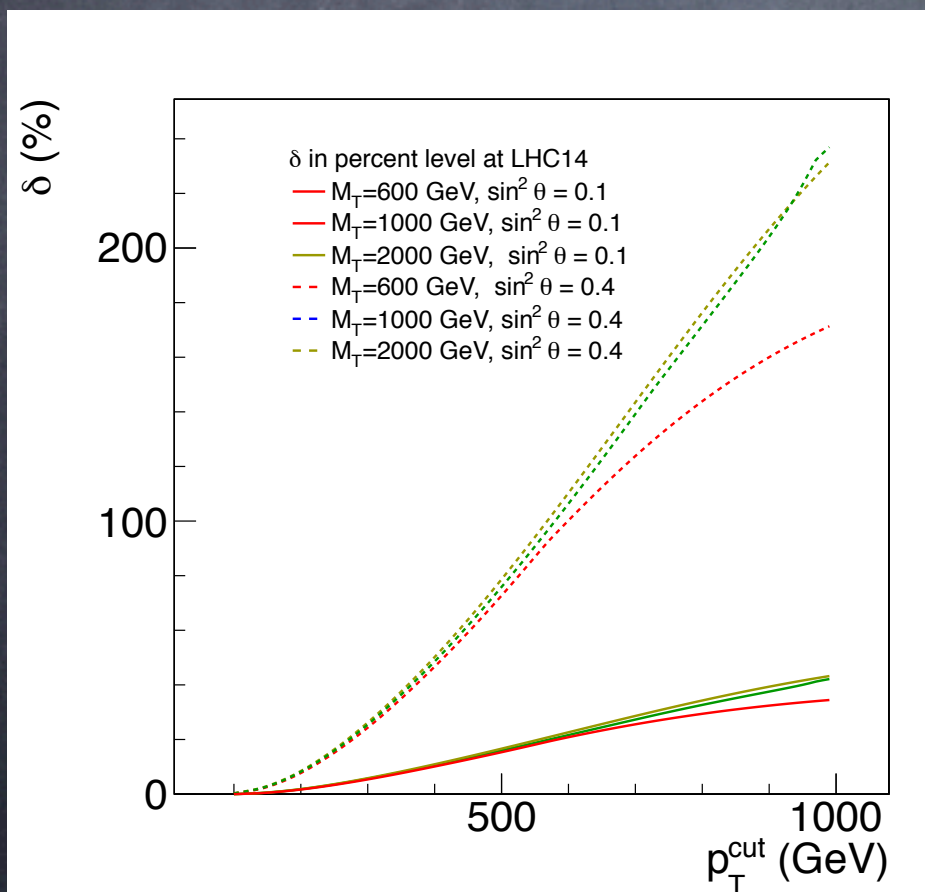
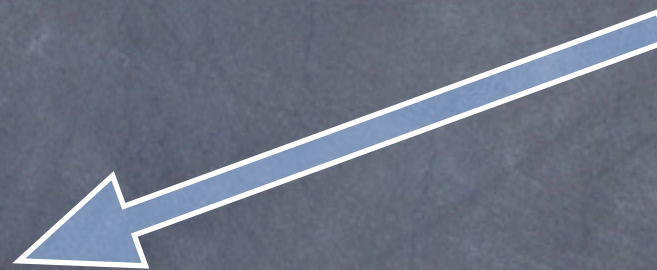
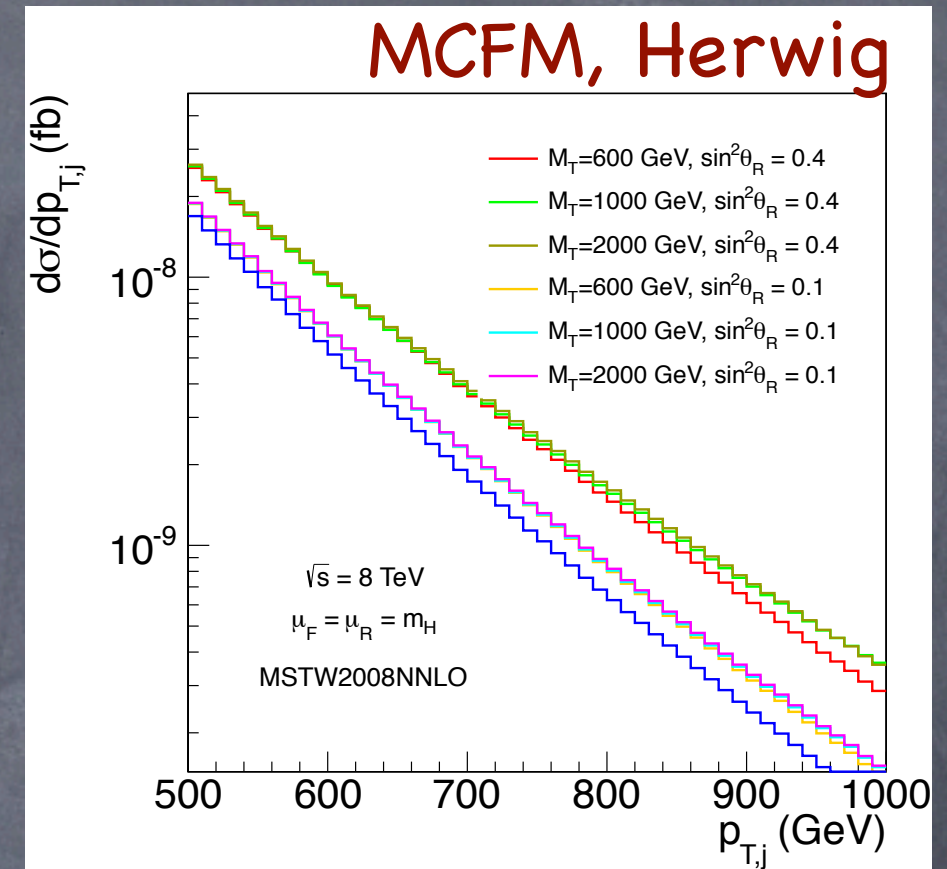
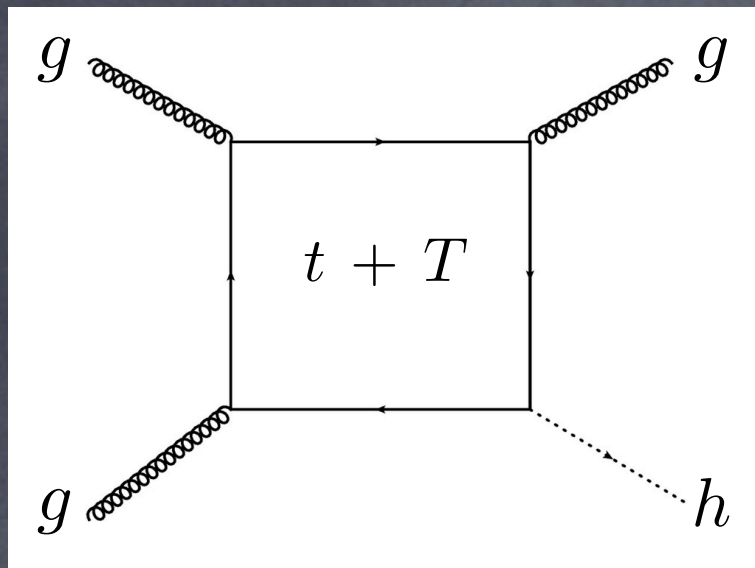
Banfi, Martin, VS. 2013





# Higher order corrections evade the cancellation

Banfi, Martin, VS. 2013

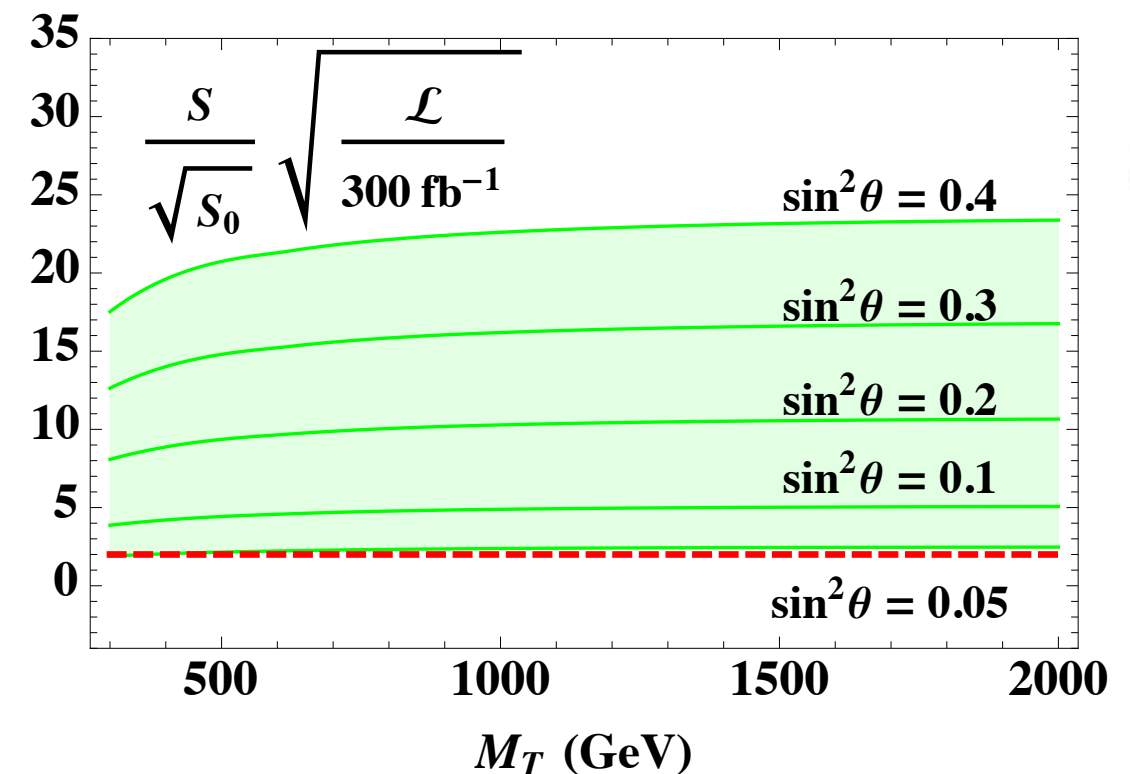
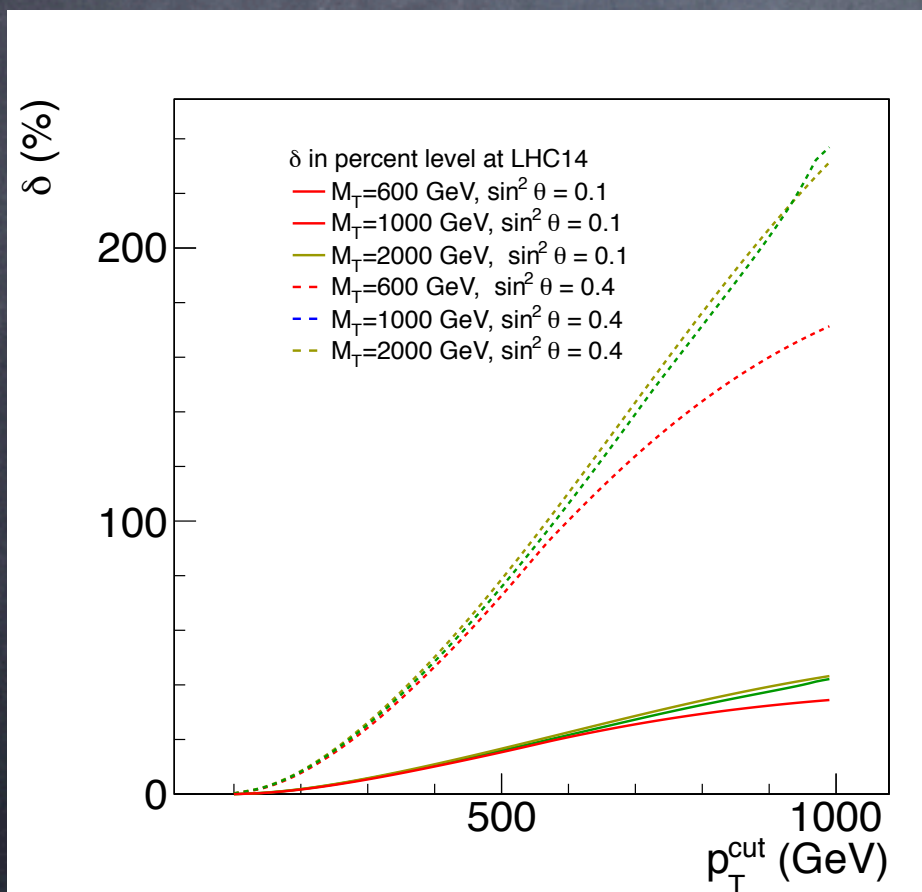
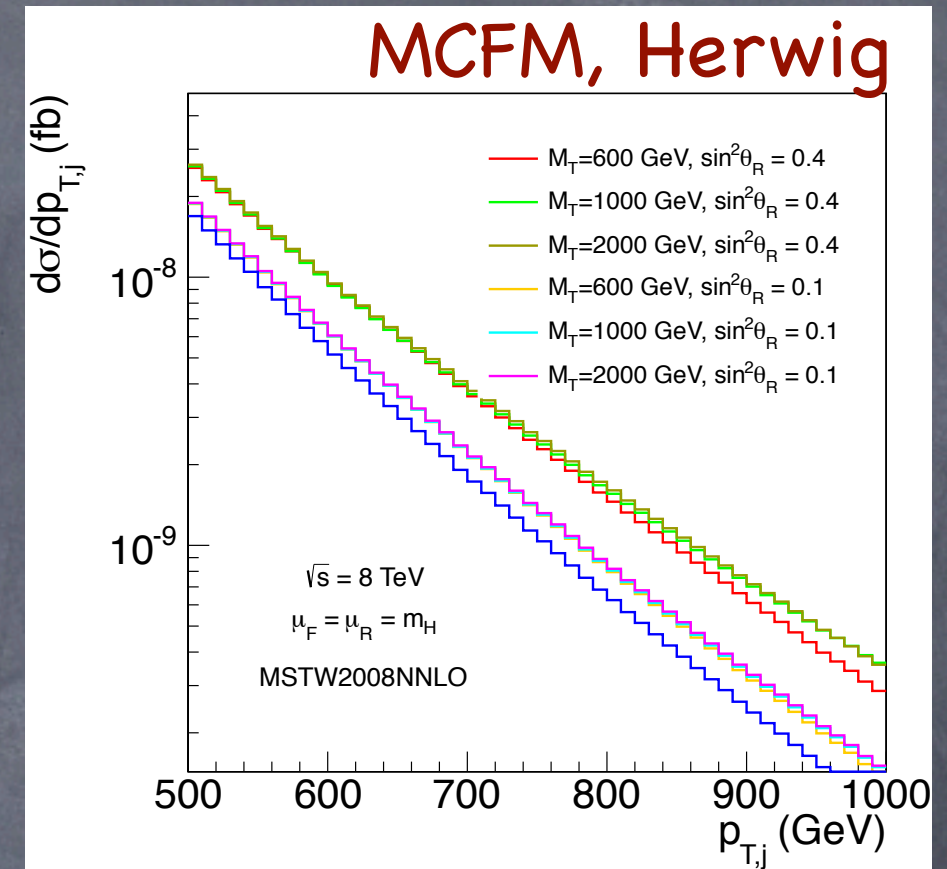
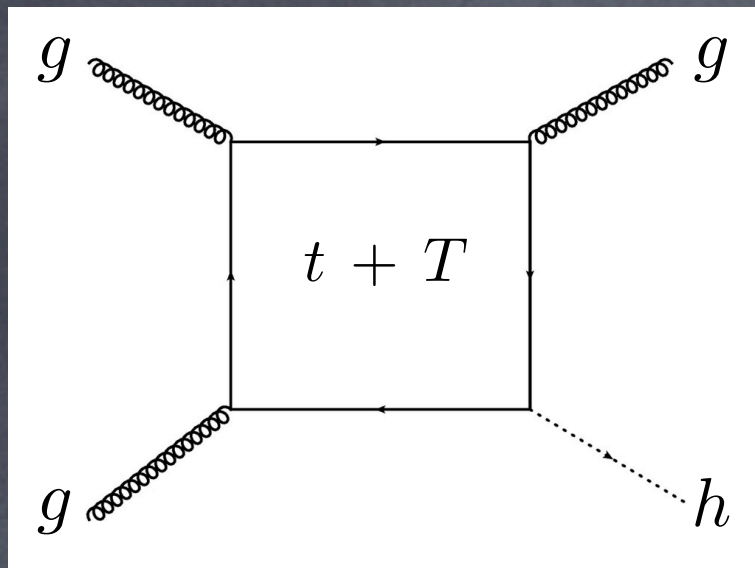


limited by lack of  
NLO computations



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Banfi, Martin, VS. 2013



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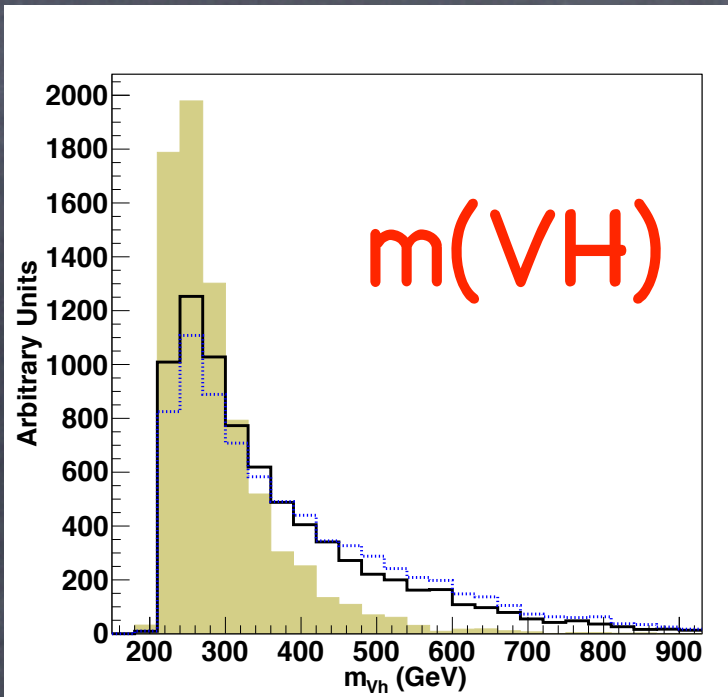


Pushing to QCD NLO  
w/Campbell, Williams. In prep



So far, Higgs anomalous couplings is LO, but

Anomalous couplings  $\rightarrow$  tails of distributions  
need higher precision to beat QCD/  
EW corrs

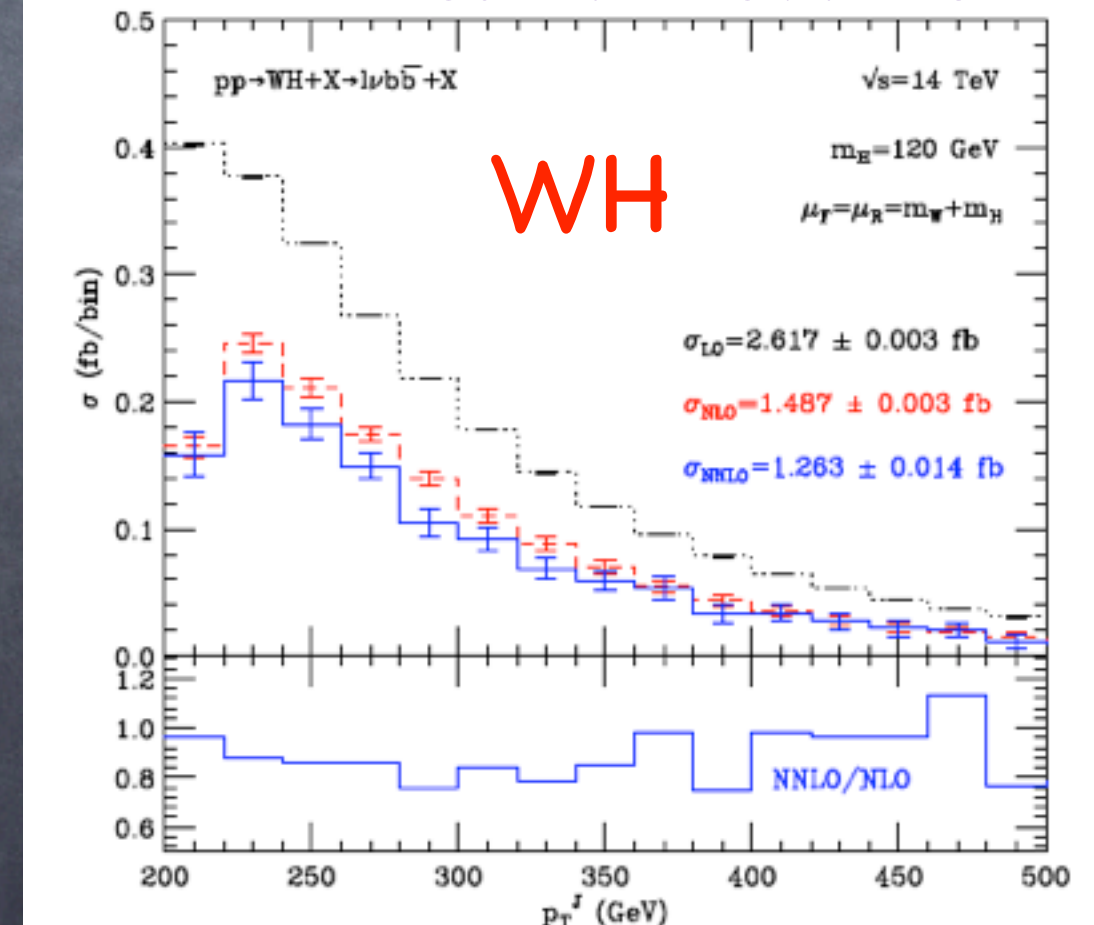


Alloul, Fuks, VS. 1310.5150

SM state of the art

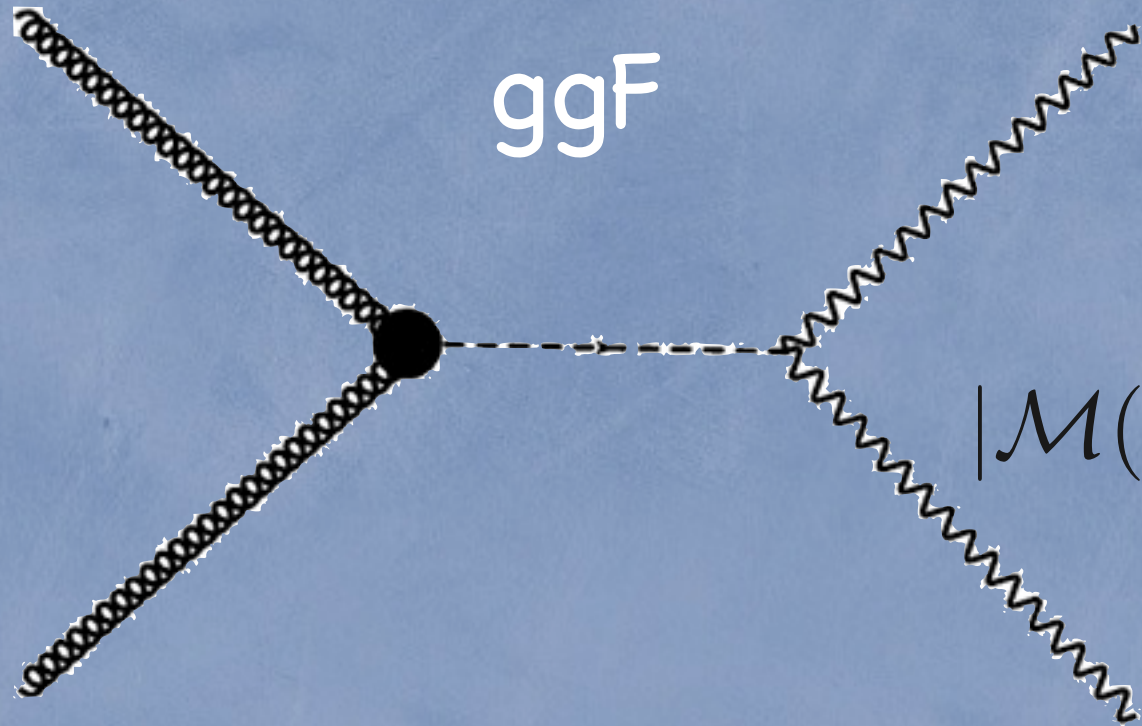
Need NLO for  
anomalous couplings:  
different Lorentz  
structure, not just a  
K factor

Grazzini et al. 1107.1164





# QCD NLO corrections in ggF, VBF and AP can be handled as follows



$$|\mathcal{M}(p_1 + p_2 \rightarrow H(\rightarrow V^*V^{(*)}) + X)|^2$$

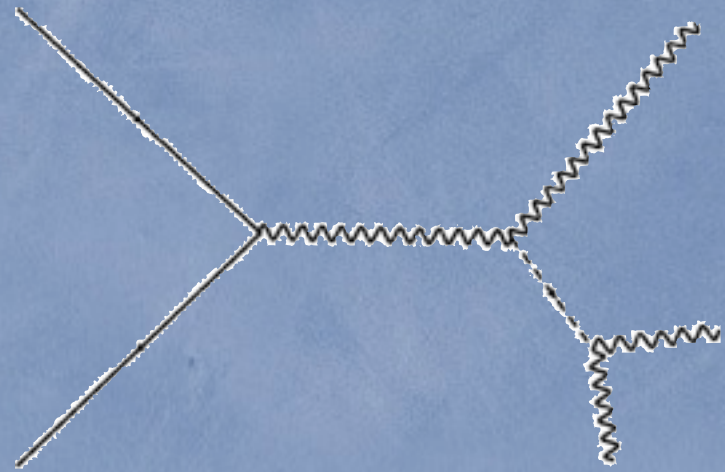
$$= |\mathcal{M}(p_1 + p_2 \rightarrow H + X)|^2 \mathcal{P}_H(s_{VV})^2 |\mathcal{M}(H \rightarrow V^*V^{(*)})|^2$$

NLO is obtained by re-weighting H+jets results by

$$\mathcal{R}(\Lambda) = \frac{|\mathcal{M}^\Lambda(H \rightarrow f)|^2}{|\mathcal{M}^{SM}(H \rightarrow f)|^2}$$



same goes for associated production



$$= \mathcal{J}^\mu(p_1, p_2, X) \mathcal{P}_{\mu\nu}(V) J^\nu(V, H)$$
$$\mathcal{A}_\Lambda^{NLO}(p_1 + p_2 \rightarrow V(\rightarrow H + V) + X) =$$
$$\mathcal{J}_{NLO}^\mu(p_1, p_2, X) \mathcal{P}_{\mu\nu}(V) J_\Lambda^\nu(V, H)$$

and vector boson fusion

soon w/ Campbell and Williams  
public MCFM+POWHEG implementation  
of Higgs anomalous couplings at QCD NLO



Into the future



# Higgs factories

1. HL-LHC

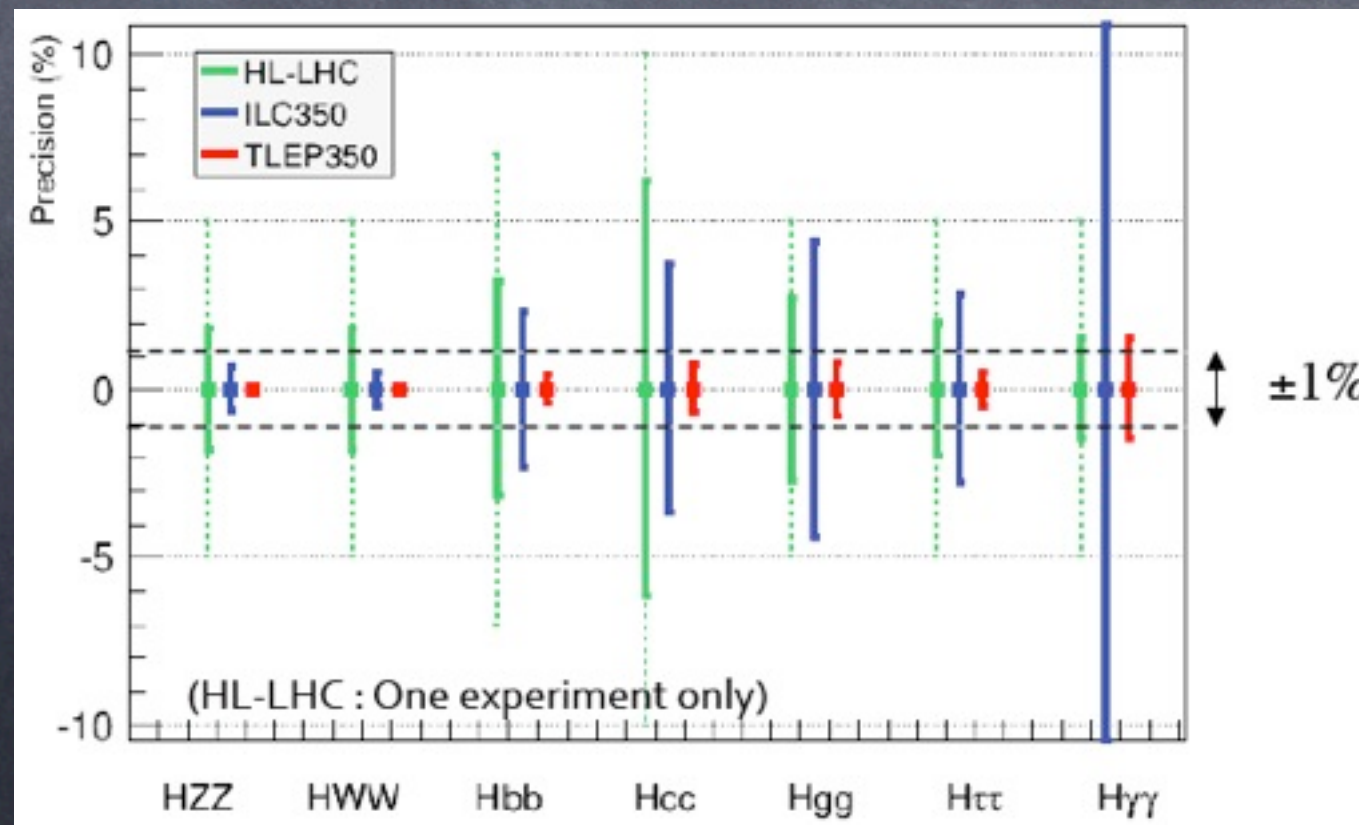
2. A linear collider

ILC up to 500 GeV

CLIC up to 3 TeV

4. A circular  $e^+e^-$  collider: LEP3, TLEP  
(host of 100 TeV of VHE collider)

5. A muon collider





# Conclusions

Higgs, scalar particle and quite SM-like  
this is just the beginning



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most powerful constraints by  
fit rates and  
differential distributions → precision



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SUSY, composite and little Higgs and Eff L  
direct/indirect (Higgs)

Higgs data at the level of PEWT for some variables



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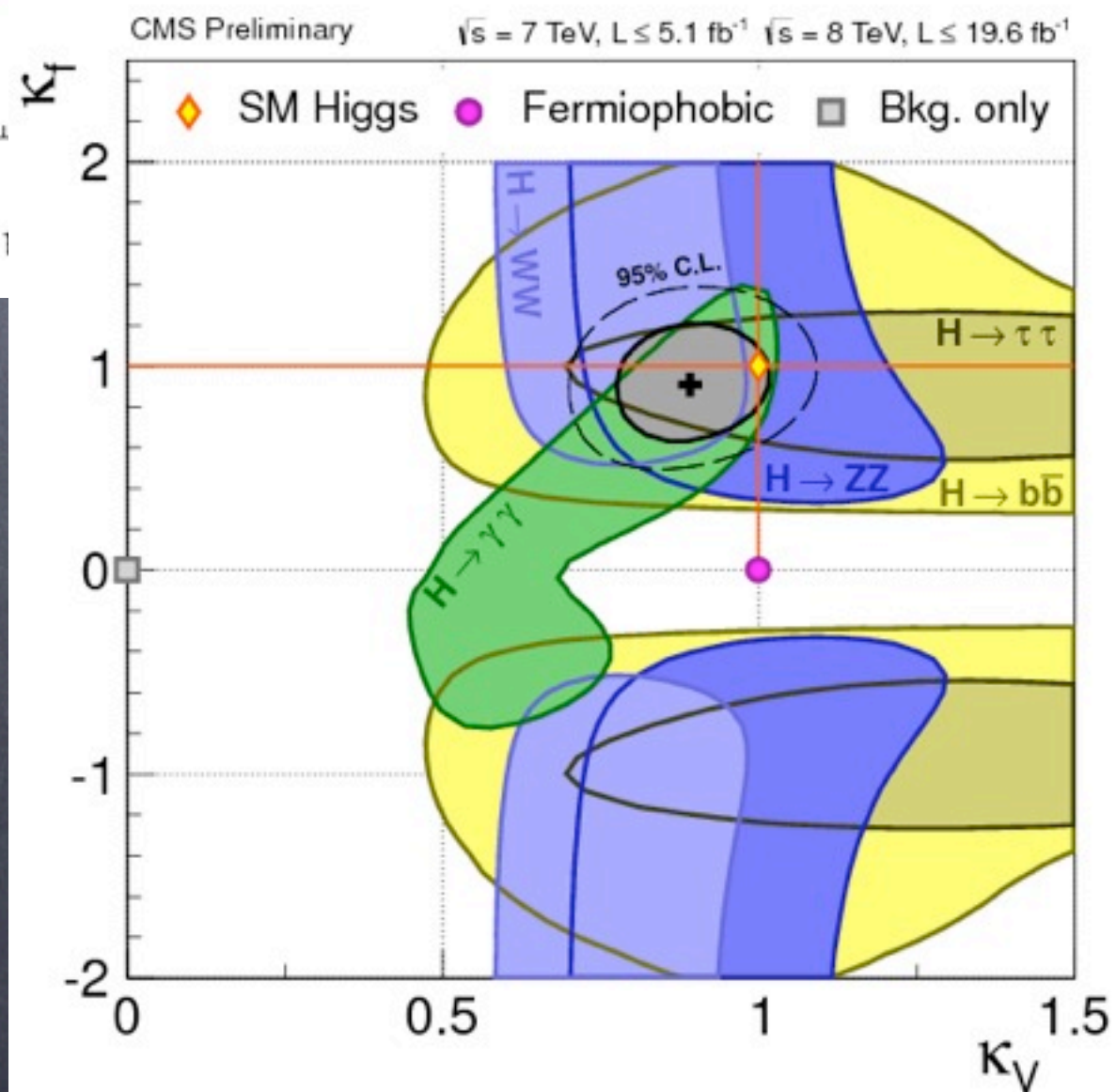
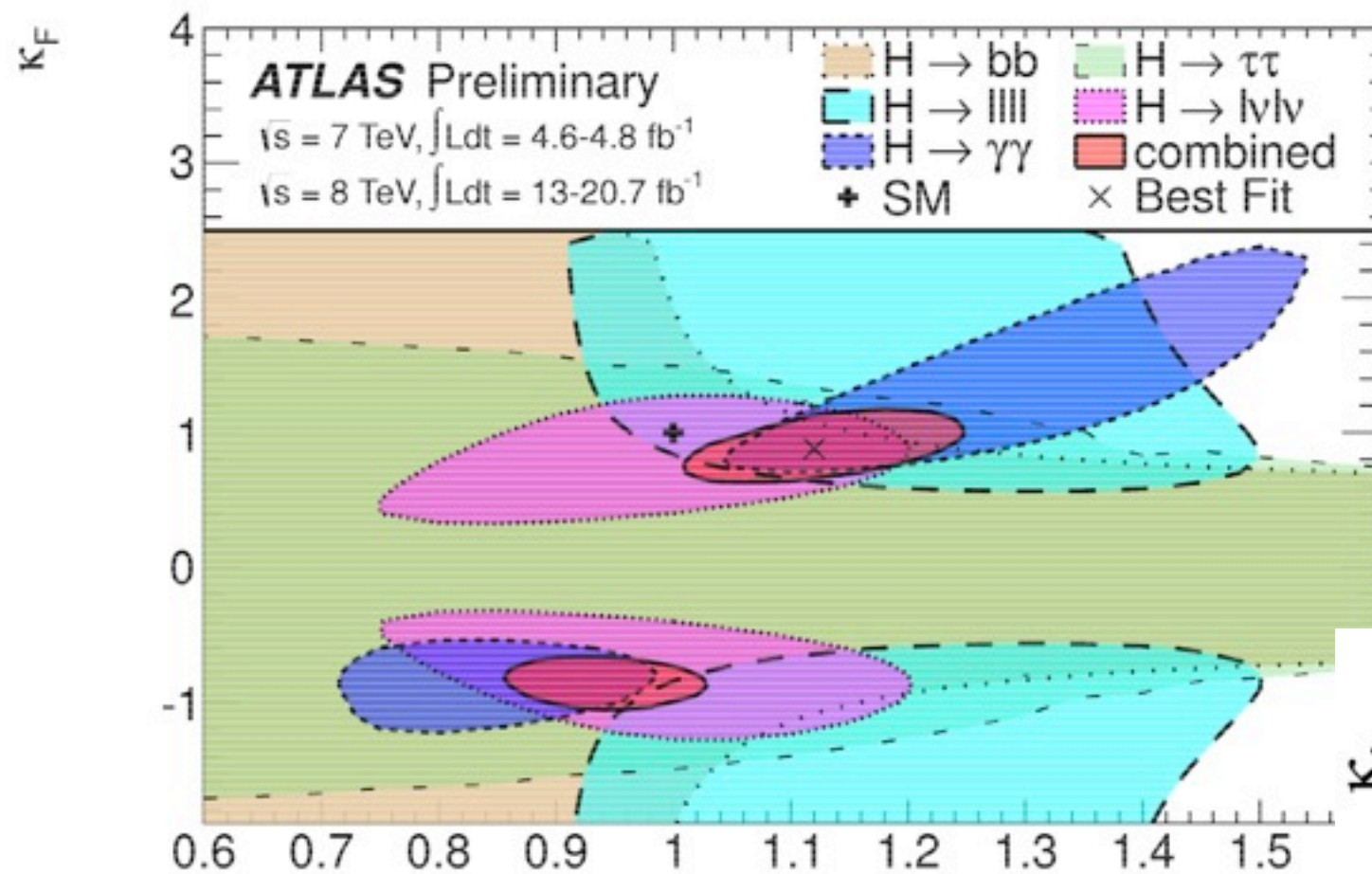
SUSY, composite and little Higgs and Eff L  
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Higgs data at the level of PEWT for some variables

What now?

a lot of the theory space cleared out  
data  $\rightarrow$  new ideas in model building







# Example

$$r_s \frac{2m_s^2}{v} S^2 h$$

Strumia et al. 2013

