

Brane Inflation : String Theory viewed from the Cosmos

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updated

A New View of the Cosmic Landscape

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How the universe can have such a small cosmological constant ?

- String theory has many meta-stable vacua, 10^{500} or more.
- They span a wide range of CC.
- Some of them have very small CC.
- This is the cosmic landscape.
- It is easy to convince ourselves that one of the string vacuum site in the landscape describes our universe.

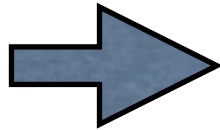
Puzzle : of all sites, why we end up in one with such a small CC ?

If we start at any meta-stable site with a long lifetime, we'll still be there now.

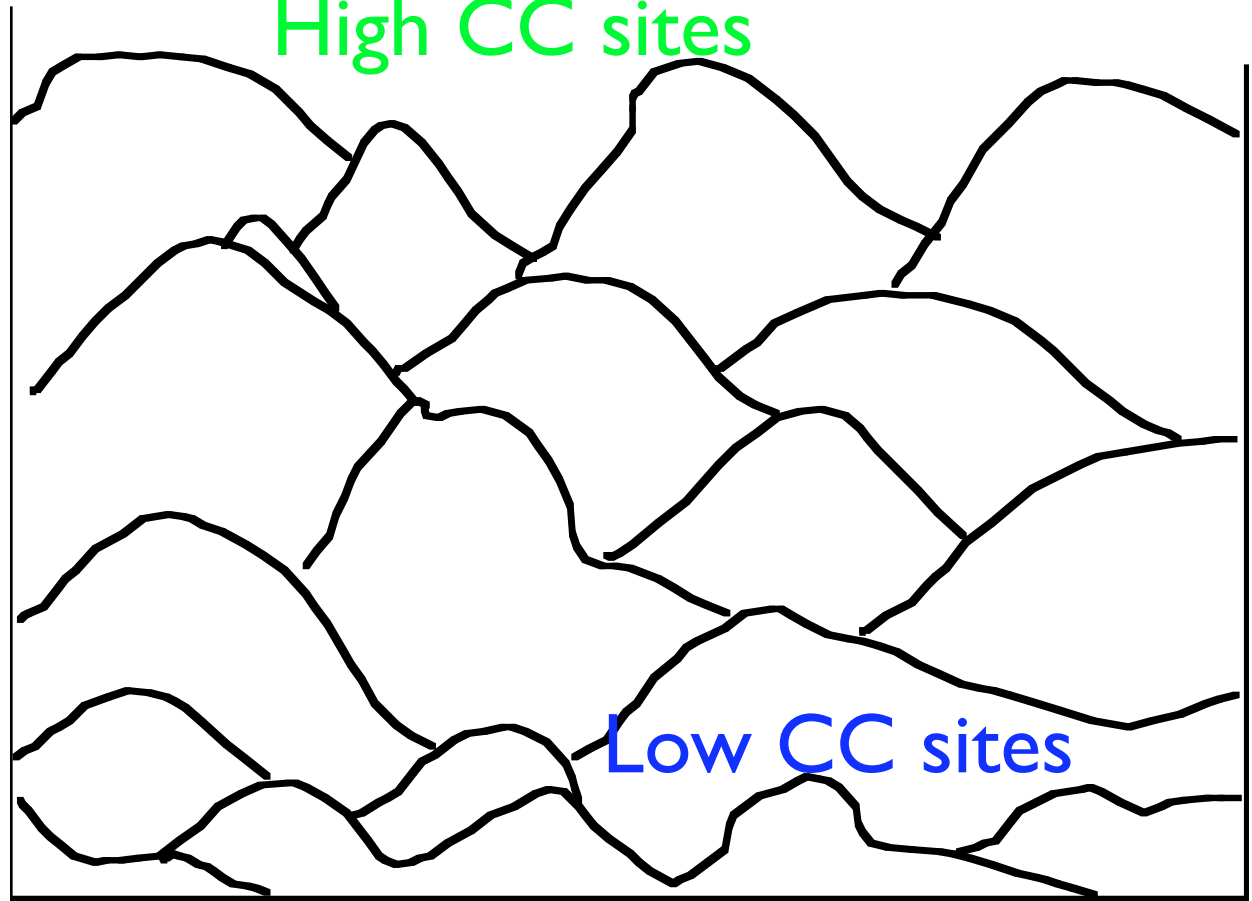
Eternal inflation !

What are the conditions under which a low CC universe will emerge naturally ?

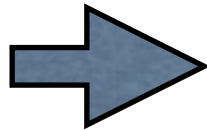
very short
lifetimes



High CC sites



exponentially
long lifetimes



Low CC sites

Suppose generic meta-stable dS sites have very short lifetimes while sites with small CC have long lifetimes.

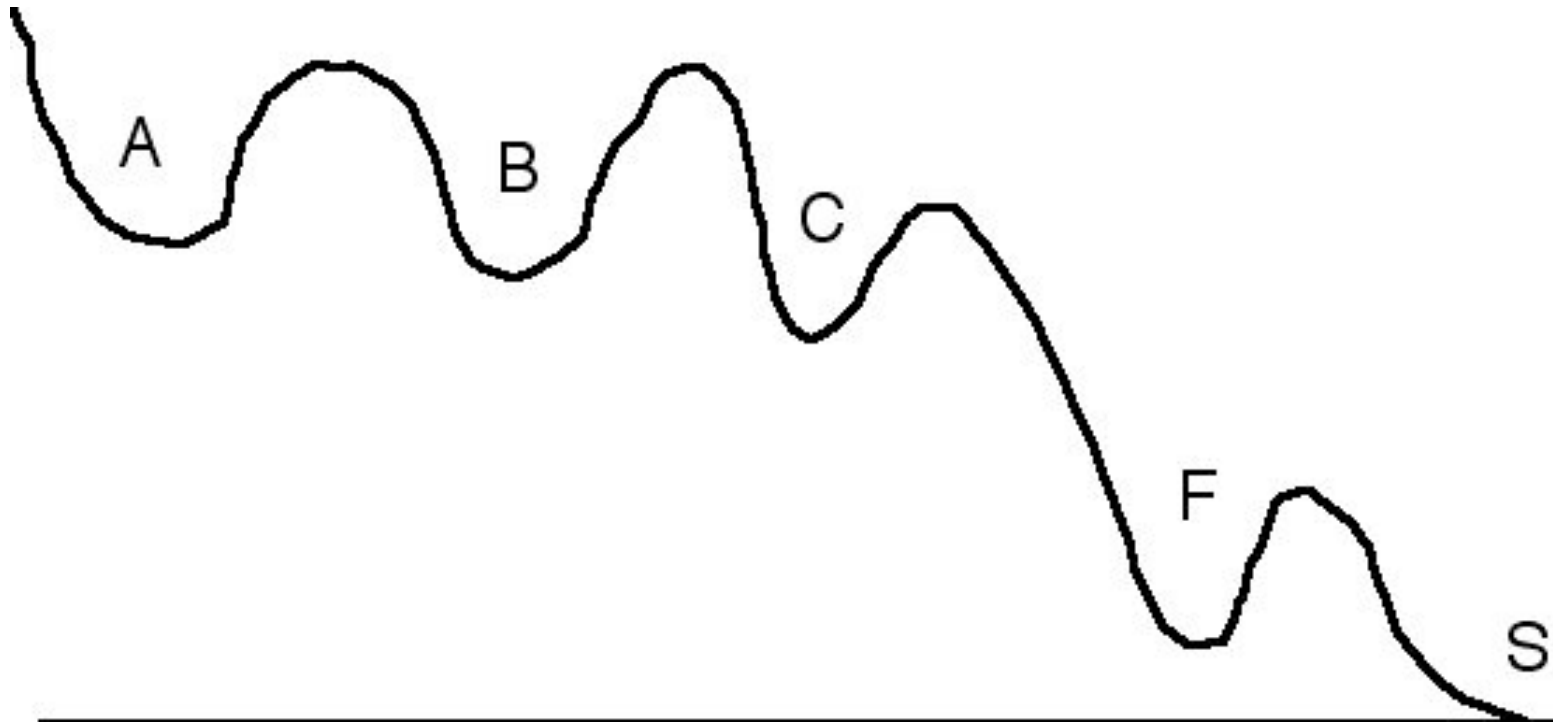
If we start at a generic site, it will quickly decay (and this decay process can repeat any number of times) until the universe reaches a site with a low CC and stays there.

EVEN BETTER :

Generic meta-stable dS sites have very short lifetimes (less than 1 second) while sites with CC below a certain value have exponentially long lifetimes (larger than the age of the universe).

Such a scenario is entirely possible.

- Vastness of the Landscape
- Resonance tunneling
- Topography of the Landscape



Resonance Tunneling :

Tunneling probabilities $T_{A \rightarrow B} \sim T_{B \rightarrow C}$ are exponentially small

When the condition is right : $T_{A \rightarrow C} \sim 1$

Resonance tunneling

Merzbacher, Quantum Mechanics, Chapter 7

$$\begin{pmatrix} \alpha_R \\ \alpha_L \end{pmatrix} = \frac{1}{2} \begin{pmatrix} \Theta + \Theta^{-1} & i(\Theta - \Theta^{-1}) \\ -i(\Theta - \Theta^{-1}) & \Theta + \Theta^{-1} \end{pmatrix} \begin{pmatrix} \beta_R \\ \beta_L \end{pmatrix}$$

where, in the WKB approximation,

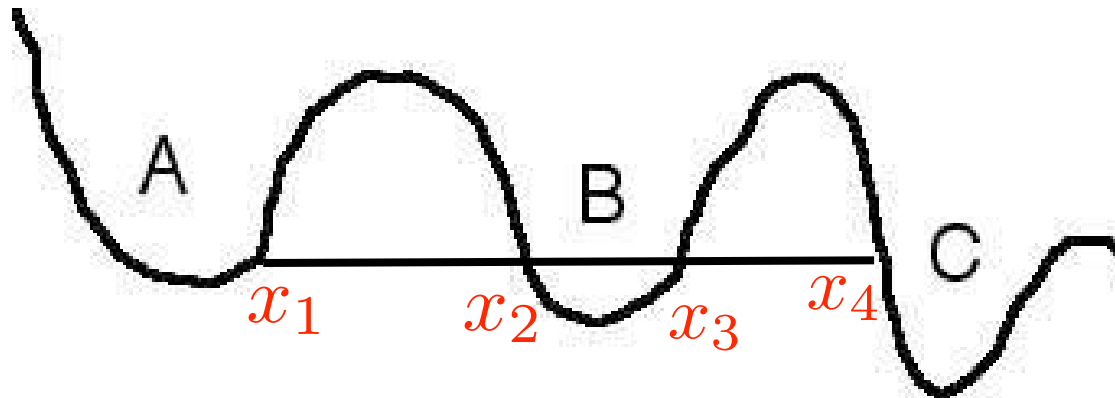
$$\Theta \simeq 2 \exp \left(\int_{x_1}^{x_2} dx \sqrt{2(V(x) - E)} \right) ,$$

$$T_{A \rightarrow B} = \left| \frac{\beta_R}{\alpha_R} \right|^2 = 4 \left(\Theta + \frac{1}{\Theta} \right)^{-2} \simeq \frac{4}{\Theta^2}$$

$$\Theta \simeq 2 \exp \left(\int_{x_1}^{x_2} dx \sqrt{2(V(x) - E)} \right) , \quad \Phi = 2 \exp \left(\int_{x_3}^{x_4} dx \sqrt{2(V(x) - E)} \right)$$

$$\frac{1}{4} \begin{pmatrix} \Theta + \Theta^{-1} & i(\Theta - \Theta^{-1}) \\ -i(\Theta - \Theta^{-1}) & \Theta + \Theta^{-1} \end{pmatrix} \begin{pmatrix} e^{-iW} & 0 \\ 0 & e^{iW} \end{pmatrix} \begin{pmatrix} \Phi + \Phi^{-1} & i(\Phi - \Phi^{-1}) \\ -i(\Phi - \Phi^{-1}) & \Phi + \Phi^{-1} \end{pmatrix}$$

$$W = \int_{x_2}^{x_3} \sqrt{2(E - V(x))} dx$$



WKB connection formula

$$\Gamma_{A \rightarrow C} = 4 \left((\Theta \Phi + \frac{1}{\Theta \Phi})^2 \cos^2 W + (\Theta / \Phi + \Phi / \Theta)^2 \sin^2 W \right)^{-1}$$

In the absence of B , $W = 0$ so $T_{A \rightarrow C}$ is very small,

$$T_{A \rightarrow C} \simeq 4\Theta^{-2}\Phi^{-2} = T_{A \rightarrow B}T_{B \rightarrow C}/4$$

However, if W satisfies the quantum condition for bound states in B ,

$$W = (n_B + 1/2)\pi$$

$$T_{A \rightarrow C} = \frac{4}{(\Theta / \Phi + \Phi / \Theta)^2}$$

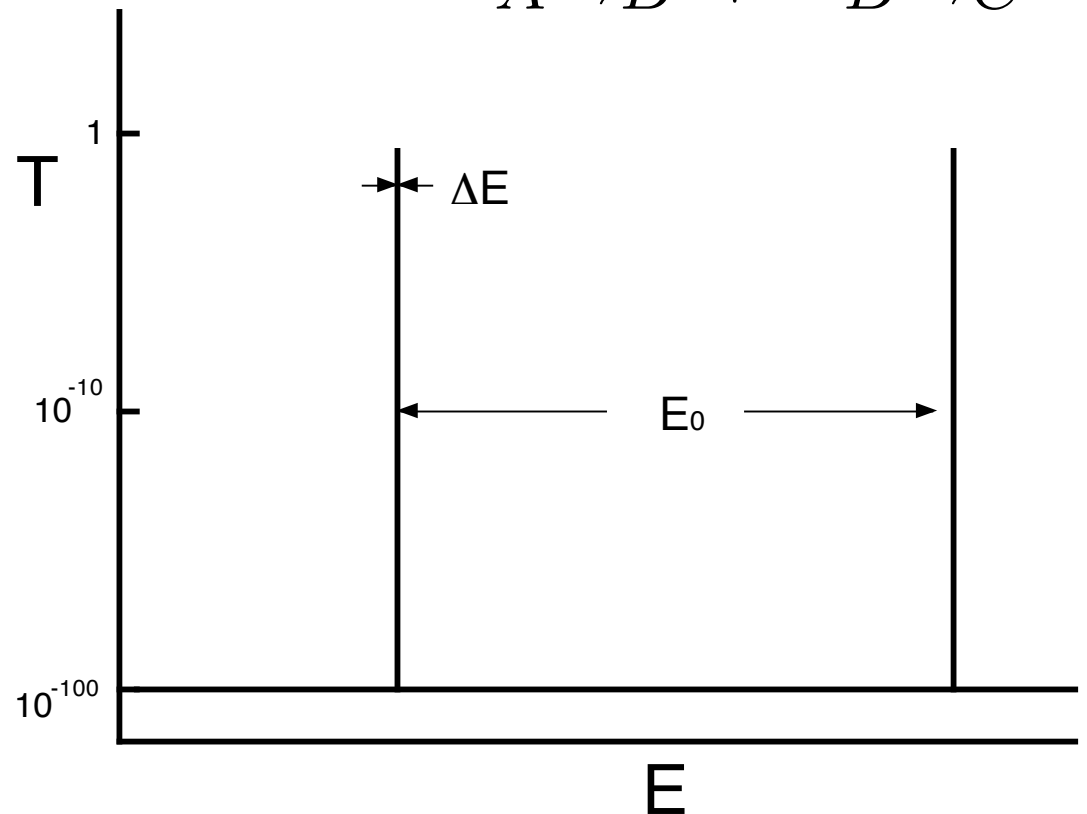
$$T_{A \rightarrow B} \rightarrow T_{B \rightarrow C} \quad \Rightarrow \quad T_{A \rightarrow C} \rightarrow 1$$

How likely to hit a resonance ?

$$P(A \rightarrow C) = \frac{\Delta E}{E_0} \simeq \frac{2}{\pi \Theta \Phi} \left(\frac{\Theta}{\Phi} + \frac{\Phi}{\Theta} \right) = \frac{1}{2\pi} (T_{A \rightarrow B} + T_{B \rightarrow C})$$

$$\langle T_{A \rightarrow C} \rangle = P(A \rightarrow C) T_{A \rightarrow C} \sim \frac{T_{A \rightarrow B} T_{B \rightarrow C}}{T_{A \rightarrow B} + T_{B \rightarrow C}}$$

Average tunneling probability is exponentially larger than the naive estimate.

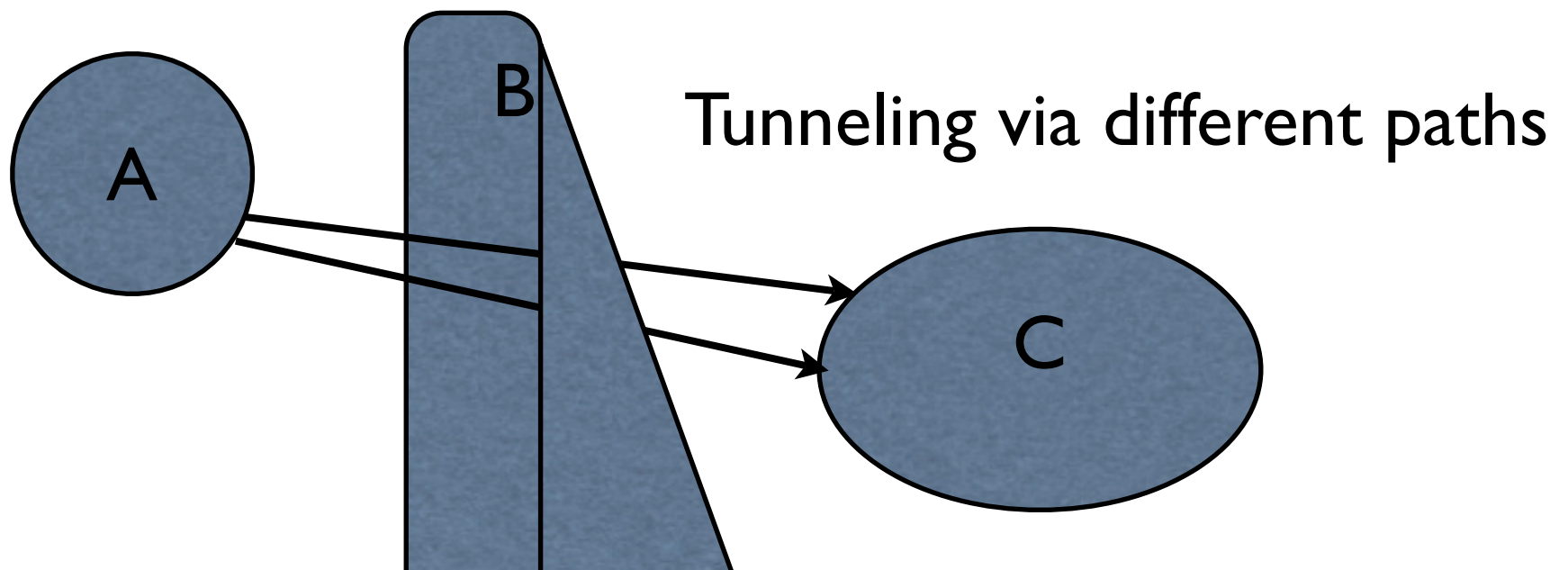


QM to QFT/Gravity

$$\phi_i(\mathbf{r}, t) \rightarrow \phi_i(t) = x_i(t)$$

$$U(\phi_i) \rightarrow V(x_i)$$

Resonance tunneling happens in QFT
as well as in Gravity ?



Application to The Landscape

- Tunneling from a dS site to an AdS site is ignored.
- Tunneling from a dS site to another dS site with a larger CC is ignored.
- Evolution in a classically allowed region should be treated quantum mechanically.
- Sum over tunneling paths.

n-step tunneling channel :

$$A \rightarrow B_1 \rightarrow B_2 \rightarrow \dots \rightarrow B_{n-1} \rightarrow C$$

Naive :

$$\hat{T}_{A \rightarrow C}(n) \simeq T_{A \rightarrow B_1} T_{B_1 \rightarrow B_2} \dots T_{B_{n-1} \rightarrow C}$$

$$\langle T_{A \rightarrow C}(n) \rangle = P(A \rightarrow C) T_{A \rightarrow C} \sim \text{Min} \left(T_{\min}, \sqrt{\hat{T}_{A \rightarrow C}(n)} \right)$$

Tunneling probability is decreasing rapidly but still exponentially larger than the naive estimate.

Probability of hitting an efficient tunneling channel in n steps :

$$P(A \rightarrow C) \sim \left(T_{A \rightarrow B_1} \dots T_{B_{n-1} \rightarrow C} \right)^{1/2} \quad n \geq 2$$

The number of efficient tunneling channels available to A :

$$N(A) = \sum_C P(A \rightarrow C) \quad \Lambda_C < \Lambda_A$$

The number of sites C available to A is exponentially large.

Suppose the distribution of vacua is independent of CC.
A site with GUT scale CC has about 10^{100} more channels
to tunnel to than our universe has.

$$N(A) = \sum_C P(A \rightarrow C) \quad \Lambda_C < \Lambda_A$$

Both the number of decay sites C and the number of tunneling channels to C open to a site A decreases as its CC decreases.

If $N(A) > 1$, then $T_A \sim 1$

If $N(A) < 1$

T_A would be exponentially small.

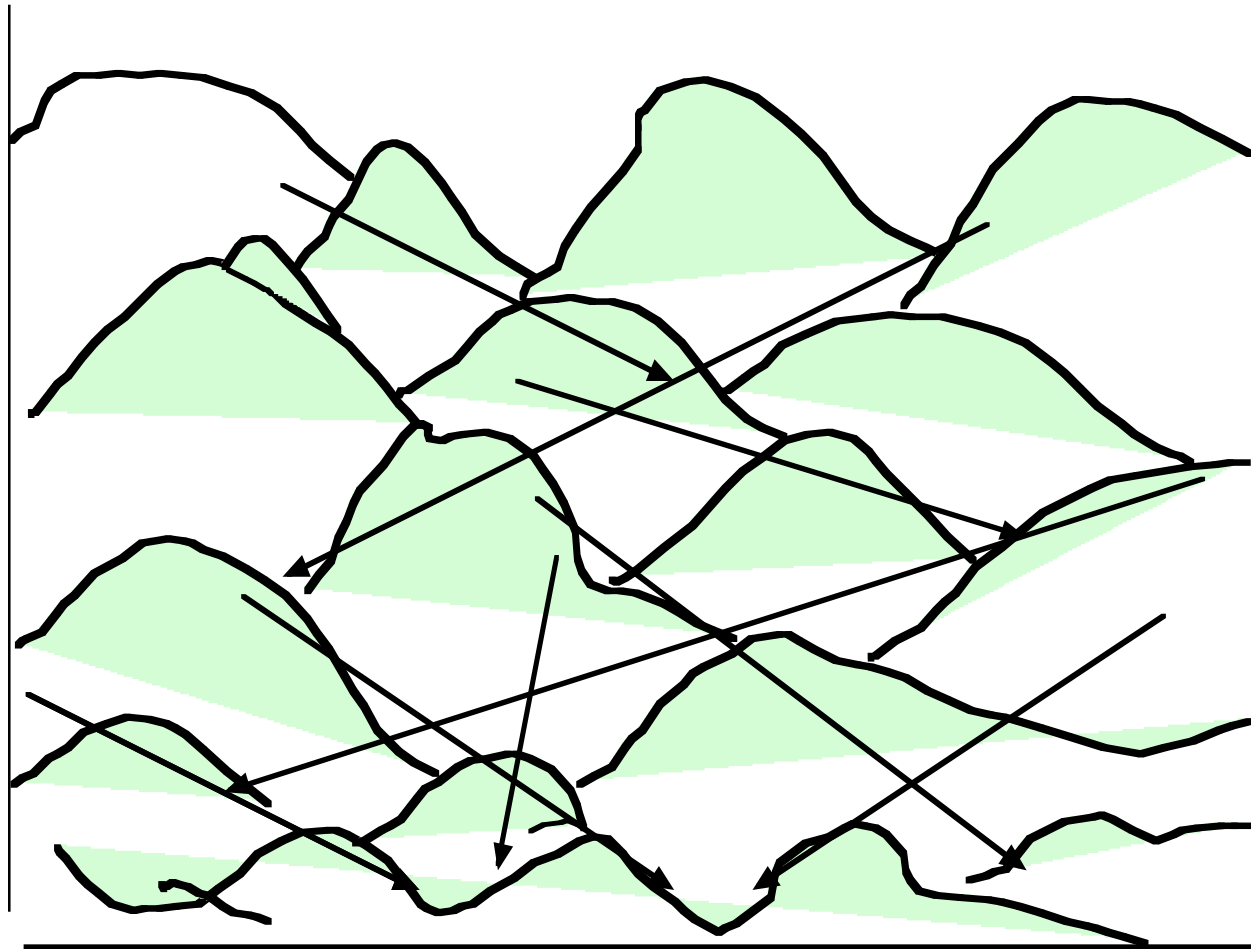
The Quantum Landscape

There is a critical CC :

$N(A) > 1$ for $\Lambda_A > \Lambda_c$, and $N(A) < 1$ for $\Lambda_A < \Lambda_c$.

↓
Decays quickly

↓
very long lifetime



Eternal Inflation ?

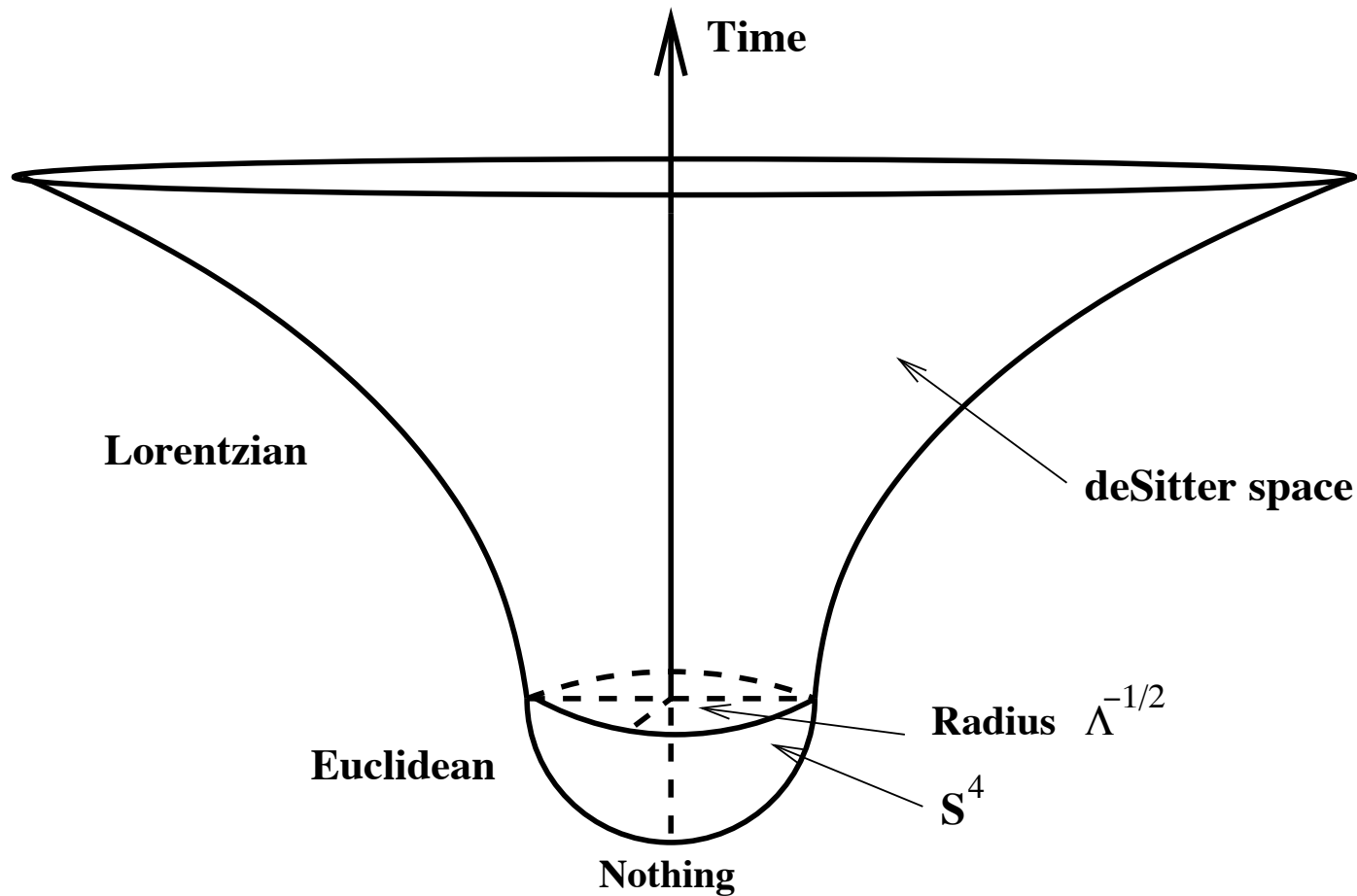
$$T \sim 1$$

$$\Gamma_A \sim M_s$$

$$H \sim M_s^2/M_P < \Gamma_A$$

No eternal inflation !

Tunneling from Nothing

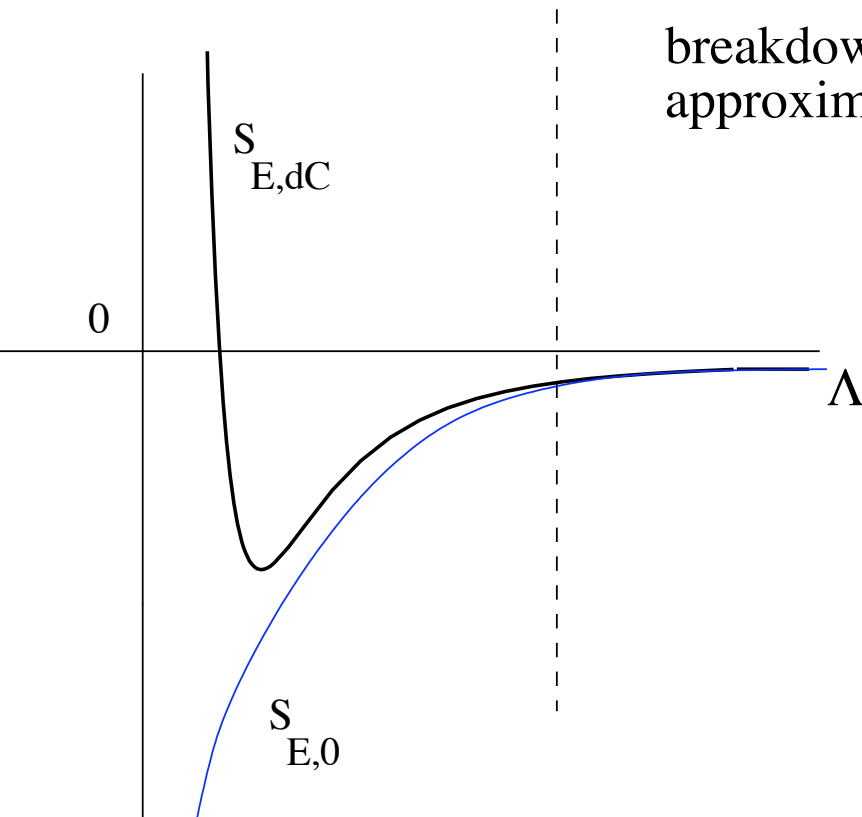


Vilenkin, 1982, 1983

Hartle and Hawking, 1983

The improved Euclidean action

$S_{E,dC}$ vs $S_{E,0}$



breakdown of semiclassical
approximation for large Λ

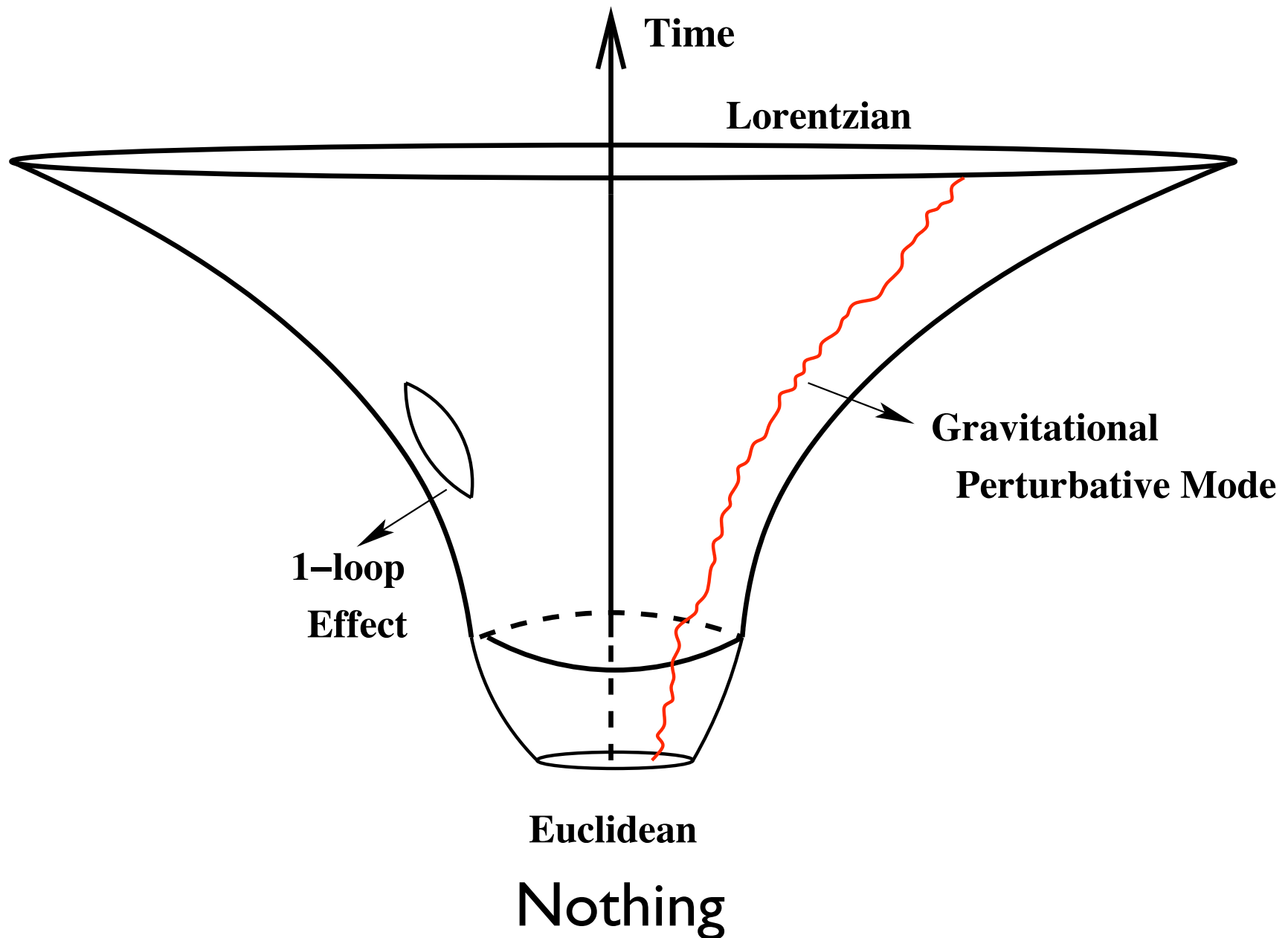
Tunneling probability

$$P \simeq e^{-S_{E,dC}} = e^F$$

$$F = \frac{3\pi}{G\Lambda} - \frac{12n_d}{\Lambda^2 l_s^4}$$

$S_{E,0}$ is unbounded from below, but the interaction with the environment has made $S_{E,dC} = -F$ bounded from below.

Tunneling from Nothing



History

- The Universe is a spontaneous creation from NOTHING.
- It has a CC somewhere below the string scale.
- It then rapidly decays to a vacuum site in the landscape with an exponentially long lifetime.
- This is where we live.
- Resonance tunneling may happen both before and after inflation.

Conclusion

If the scenario is correct, we can appreciate string theory in this new light : it provides a vast landscape so that a small CC vacuum is among its solutions, and the same vastness destabilizes all vacua except ones with a very small CC, thus allowing our universe to emerge, survive and grow.

We still need to fold this scenario into the cosmological evolution of the universe. A proper treatment requires us to start with the wavefunction of the universe. It is a challenging but tractable exercise to check the validity of this scenario.